

THE WAVELENGTH SPECTRUM SHIFT OF A  
CAVITY-DUMPED ARGON LASER-PUMPED  
RHODAMINE 6-G ORGANIC-DYE LASER

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THESIS

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by

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September 1973

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Argon Laser-Pumped Rhodamine 6-G Organic-Dye Laser

by

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ABSTRACT

The wavelength spectrum of a CW pumped Rhodamine 6-G organic dye laser is observed to shift approximately 150 Angstroms toward shorter wavelengths when pumped by 30 nsec pulses at a one megahertz repetition rate from a cavity-dumped Argon laser. Experimental evidence of the shift is presented, gain and rate equations are developed for a simplified dye laser model, and theoretical results are obtained for a computer simulation of the experiment. A comparison is made of the theoretical and experimental results and satisfactory agreement is obtained within the limits of the values of the parameters used and the assumptions made in formulating the model.





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## I. INTRODUCTION

Man's quest for knowledge, inherent curiosity, and need for a tunable coherent source of light have rapidly advanced the development of the organic dye laser since stimulated emission was observed from Trivalent Uranium (a four level laser system) by Sorokin and Stevenson [1] in 1960.

Brock, et al. [2] in 1961 proposed that organic compounds could be used as a laser medium, but stimulated emission from an organic material was not observed until 1962 when Morantz, White and Wright [3] observed it from benzophenone and naphthalene imbedded in a glass matrix. This was followed in 1963 by the observation of stimulated processes from Eu-benzoylacetate by Lempicki and Samuelson [4].

Stimulated emission from an organic dye, however, was not accomplished until 1966 when it was observed by Sorokin and Lankard [5] from Chloroaluminum phthalocyanine dissolved in ethyl alcohol when pumped by a giant pulse ruby laser. Sorokin and Lankard [6] also reported in 1967 the first observed stimulated emission from flashlamp pumped Acridine Red, Rhodamine 6-G (the most common dye used in present day organic dye lasers), and Fluorescein. Schafer, Schmidt, and Volze [7] reported stimulated emission from several other organic dyes during this same time period while using the same basic procedure as Sorokin and Lankard [5], and they were the first to suggest that a dye laser operated on transitions from the excited singlet state to vibrational levels of the ground singlet state.

The feasibility of a continuous wave dye laser was discussed by Snavely [8] and was demonstrated by Snavely and Schafer [9] in 1969. CW operation was achieved in 1970 by Peterson, Tuccio and Snavely [10] and



Hercher and Pike [11] using an Argon ion laser as the pump source and Rhodamine 6-G as the organic dye.

The first CW dye lasers were only about one percent efficient. This has been raised to the present day level of thirty-five percent through the efforts of several groups of people including Hercher and Pike [12, 13], Kohn, Shank, Ippen and Dienes [14, 15], and Tuccio and Strome [16, 17].

The evolution of the dye laser has seen it constructed in many ways. The tunable distributed-feedback dye laser built by Shank, Bjorkholm and Kogelnik [18], the prism-dye laser constructed by Chandra, Takeuchi, and Hartmann [19], and the evanescent-field-pumped dye laser demonstrated by Ippen and Shank [20] are a few of the many sophisticated methods that have been used to obtain dye laser emission. The most common dye laser configuration, however, has the pump and dye laser cavities aligned with the beams coincident and the dye flowing perpendicular to the pump beam through a dye cell.

Several uses have been proposed for dye lasers. Bloom [21] has suggested its use for attenuation measurements on narrow-band absorption lines in molecular gases, and Sorokin, Lankard, Moruzzi and Hammond [22] have proposed using the dye laser for optical studies of rare-earth ions, photochemistry, and double-quantum absorption spectroscopy. The dye laser characteristics which make it exceptionally well suited to the above uses are its tunable range which can be as much as 1100 Angstroms from a single dye, and its narrow linewidth which is normally less than 0.5 Angstroms.

Uses for the dye laser which have already been published include infrared difference-frequency generation [23], megawatt tunable second harmonic and sum frequency generation [24], studies of the Sodium D





resonance lines by high resolution Spectroscopy [25], and detection of OH in the atmosphere [26]. Further uses of the dye laser are limited only by the interest and ingenuity of man.

In this work it has been observed that the output wavelength spectrum of a CW pumped Rhodamine 6-G organic dye laser shifts toward shorter wavelengths when pumped by short intense pulses from a cavity-dumped Argon laser. This shift to the green region of the visible wavelength spectrum is important in that the new spectrum can be matched to the wavelength response curve of a crystal to produce a tunable dye laser optical memory system. The new spectrum also provides the opportunity to investigate laser propagation in blue and green sea water over an appreciable wavelength range using a single dye.

The purpose of this thesis is to provide experimental results of the observed spectrum shift, to obtain a computer simulation of a dye laser model, and to compare experimental and theoretical results. Section II discusses dye laser theory based on energy level considerations and gives the development of the dye laser gain and rate equations for an assumed physical model. Section III presents experimental procedures and results and Section IV discusses the solution procedures and results of the theoretical model. Section V provides a comparison of experimental and theoretical results and conclusions. The computer program used to solve the dye laser model is listed in Appendix A.



## II. DYE LASER THEORY

### A. ORGANIC DYE PROPERTIES

#### 1. Energy Level Description

A typical energy level diagram for an organic dye is shown in Figure 1. The reference level A is the electronic ground state of the molecular singlet state  $S_0$ .  $S_1$  and  $S_2$  are the first and second excited singlet states, and  $T_1$  and  $T_2$  are the first and second excited triplet states. The small letters a and b indicate molecular vibrational energy levels and the primed letters indicate molecular rotational energy levels within a state. The typical separation of vibrational energy levels is 0.1 electron volts and that of rotational energy levels is 0.001 electron volts [27] so that each state may be viewed as a continuous band of energy.

#### 2. Dye Laser Process

The dye laser process starts with the excitation of molecules from level A of  $S_0$  to an upper vibrational or rotational level of  $S_1$ . The excited molecules then decay very rapidly ( $10^{-10}$  -  $10^{-13}$  sec) [28] by nonradiative internal conversion to B.

The excited molecule in B has three options. It may decay spontaneously to a or a' (called fluorescence), it may undergo a stimulated transition to a or a', or it may travel via intersystem crossing to the lower level of  $T_1$ .

Fluorescence depends on the natural lifetime of B. The fluorescence spectrum of an organic dye is governed by the Franck-Condon principle [29] which states that preferred electron transitions are determined by the wave functions of the individual energy levels.



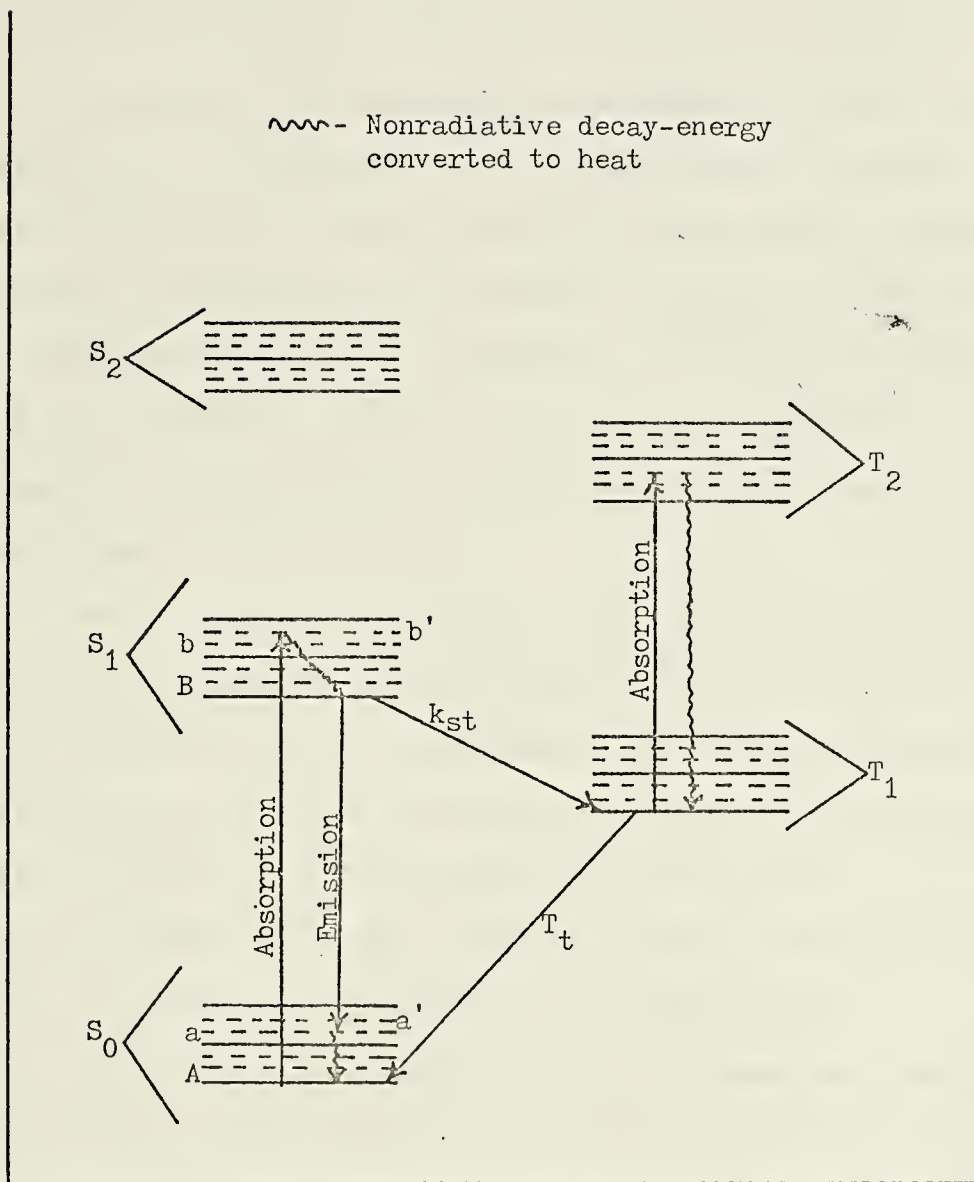


Figure 1: Organic dye energy levels.



Stimulated transition to a or a' can occur only if the excitation pulse is fast and intense enough to create a much larger than normal concentration of molecules in  $S_1$ , called critical inversion, so that coherent emission may take place from the dye. The critical inversion as well as the emission wavelength are governed by losses in the dye laser system. Once the excited molecule reaches a or a' it decays nonradiatively to A.

Intersystem crossing is the least desirable transition for an excited molecule. The singlet state is characterized by opposite electron spins and the triplet state by parallel electron spins. As a result, according to the laws of quantum mechanics, the  $S_1 - T_1$  transition is spin forbidden and is in fact approximately  $10^{-6}$  [30] less probable than the  $S_1 - S_0$  transition. However, the  $S_1 - T_1$  transition does occur and is significant enough to have an important negative effect on the dye laser process.

The intersystem crossing rate time constant,  $k_{st}^{-1}$ , is generally much smaller than the  $T_1$  state lifetime,  $T_t$ , therefore, state  $T_1$  acts as a time dependent trap for dye molecules. This not only reduces the possible dye laser efficiency by removing dye molecules from the singlet system, but also, since the absorption spectrum of the  $T_1 - T_2$  transition of an organic dye usually overlaps the dye fluorescence spectrum, another loss mechanism is added to the dye laser system.  $S_1 - S_2$  absorption is possible but it is usually neglected because the fraction of molecules in  $S_1$  needed to achieve critical inversion is very small [31].

Transitions to A are generally assumed to be from the lowest level of  $T_1$ . These transitions may be either radiative, called phosphorescence, or, as is the situation for most organic dyes, nonradiative. The longer  $T_t$  is, the faster the dye pumping pulse has to be to excite





enough molecules to reach critical inversion before the  $T_1$  trap prevents dye laser emission.

The  $T_1$  trap problem has caused a lot of research to be done to find substances which will shorten  $T_t$  via collisions with molecules in  $T_1$ . This action is called quenching and two very effective quenching agents are cyclooctatetraene (COT) and molecular oxygen.

## B. DYE LASER MATHEMATICS

### 1. Dye Laser Gain Equation

The dye laser gain equation is a mathematical statement of the production and loss rates for photons in the laser cavity. Photons are produced by stimulated emission from the excited singlet state to the ground singlet state. Photons are lost by  $S_0 - S_1$  and  $T_1 - T_2$  absorption and by loss through extrinsic means such as output mirrors, optical surface scattering, etc. Photon production is also possible by stimulated emission from the excited triplet state but nonradiative decay is so rapid from this state to the lower triplet state that this source of photons can be neglected.

A simplified schematic diagram of a dye laser is shown in Figure 2.  $L_1$  is the dye cell length,  $R_1$  and  $R_2$  are the input and output mirror reflectivities,  $T_1$  and  $T_2$  are the dye cell window transmittances, and  $n_0, n_1, n_2$ , etc. are the photon densities at any time at the indicated positions of a round trip in the laser cavity. This particular cavity is used as the model for the following derivation of the gain equation.

The rate of photon production per unit length in the dye may be written as follows:

$$\frac{dn}{dl}_{\text{total}} = \frac{dn}{dl}_{\text{emission}} - \frac{dn}{dl}_{\text{absorption}} \quad (1)$$







The first term on the right hand side of (1) is photon production by stimulated emission and the second term is photon loss by singlet and triplet state absorption.

Yariv and Gordon [32] have shown that the expression for the transition rate caused by a monochromatic beam of light of wavelength  $\lambda$  is

$$\frac{dn}{dt}_e = \frac{N_s \lambda^4 E(\lambda) n(\lambda)}{8 \pi T_s \bar{n}^3} \text{ cm}^{-3} \text{ sec}^{-1} \quad (2)$$

where  $N_s$  is the population density of the excited singlet state,  $n(\lambda)$  is the number of photons of wavelength  $\lambda$  per  $\text{cm}^3$  of active medium,  $T_s$  is the excited singlet state lifetime,  $\bar{n}$  is the index of refraction, and  $E(\lambda)$  is the fluorescence lineshape function normalized so that

$$\int_0^{\infty} E(\lambda) d\lambda = \phi = \text{quantum yield} = \frac{\text{photons emitted}}{\text{photons absorbed}}$$

An expression for the first term of (1) can now be derived. Since

$$\frac{dn}{dt}_e = \frac{dn}{dl}_e \times \frac{dl}{dt} \quad \text{and} \quad \frac{dl}{dt} = \text{velocity} = \frac{c}{\bar{n}}$$

where  $c$  is the speed of light, it follows that

$$\frac{dn}{dl}_e = \frac{N_s \lambda^4 E(\lambda) n(\lambda)}{8 \pi c T_s \bar{n}^2} \quad (3)$$

The molecular extinction coefficient, which is the absorption cross-section for a single dye molecule, can be used to write the singlet and triplet absorption loss per unit length for (1). These terms are

$$\frac{dn}{dl}_{\text{singlet}} = n(\lambda) N_o \sigma_s(\lambda) \quad (4)$$

$$\frac{dn}{dl}_{\text{triplet}} = n(\lambda) N_t \sigma_t(\lambda) \quad (5)$$

where  $N_o$  and  $N_t$  are the population densities of the ground singlet and triplet states, and  $\sigma_s$  and  $\sigma_t$  are the wavelength dependent singlet and triplet absorption cross-sections.



Let

$$\sigma_e(\lambda) = \frac{\lambda^4 E(\lambda)}{8\pi c T_s \bar{n}^2}$$

and substitute (3), (4) and (5) into (1). This yields the total production rate of photons per unit length in the dye medium which is

$$\frac{dn}{dl}_{\text{total}} = n(\lambda) \sigma_e(\lambda) N_s - n(\lambda) \sigma_s(\lambda) N_o - n(\lambda) \sigma_t(\lambda) N_t. \quad (6)$$

Separation of variables of (6) and integration over the length of the dye cell with the boundary conditions that where  $l = 0$ ,  $n = n_1$ , and where  $l = L_1$ ,  $n = n_2$ , yields the solution

$$n_2 = n_1 \exp[\sigma_e N_s - \sigma_s N_o - \sigma_t N_t] L_1. \quad (7)$$

The  $n$ 's and  $\sigma$ 's are still wavelength dependent and from Figure 2 it can be seen that since the paths traveled by photons  $n_1$  and  $n_5$  are through the same medium but in opposite directions that (7) also holds for the relation between  $n_6$  and  $n_5$ .

It is obvious from Figure 2 that  $n_1 = n_0 T_1$ ,  $n_3 = n_2 T_2$ ,  $n_4 = n_3 R_2$ ,  $n_5 = n_4 T_2$ ,  $n_7 = n_6 T_1$ , and  $n_8 = n_7 R_1$ . Substitution of these relations consecutively into (7) gives the result

$$n_8 = n_0 \exp[\sigma_e N_s - \sigma_s N_o - \sigma_t N_t + \frac{1}{2L_1} \ln(R_1 R_2 T_1^2 T_2^2)] 2L_1. \quad (8)$$

Let  $G(\lambda)$  be the equation in brackets in (8) and define it to be the gain. Then

$$n_8 = n_0 \exp[2 G(\lambda) L_1], \text{ and}$$

$$G(\lambda) = \sigma_e N_s - \sigma_s N_o - \sigma_t N_t + \frac{1}{2L_1} \ln[R_1 R_2 T_1^2 T_2^2] \quad (9)$$

which is the desired gain equation and is similar to the results obtained by Snavely [8] and McColgin, et al. [33].





## 2. Dye Laser Rate Equations

The dye laser will lase at the value of  $\lambda$  where (9) is a maximum. To find this wavelength it is necessary to solve the state population density rate equations for the dye laser. These equations are

$$\frac{dN_s}{dt} = -\frac{1}{T_s} N_s + P(t)N_o \quad (10A)$$

$$\frac{dN_t}{dt} = -\frac{1}{T_t} N_t + k_{st} N_s \quad (10B)$$

$$\frac{dN_o}{dt} = -P(t)N_o + \left(\frac{1}{T_s} - k_{st}\right)N_s + \frac{1}{T_t} N_t \quad (10C)$$

$$N = N_o + N_s + N_t \quad (10D)$$

$N$  is the total population density and (10D) is true only for an enclosed system undergoing no photochemical processes.  $N_t$ ,  $N_o$ , and  $N_s$  are as defined for the gain equation,  $T_s$  and  $T_t$  are the lifetimes of the singlet and triplet states,  $k_{st}$  is the excited singlet to ground triplet intersystem crossing rate constant, and  $P(t)$  is the optical pumping rate.

There are several assumptions made in writing the rate equations in the form of (10A), (10B), (10C) and (10D). These assumptions are:

(a) The number of molecules involved in intersystem crossing from the triplet to the singlet state is negligible.

(b) The molecules initially excited into the first excited singlet state reach thermal equilibrium in a short time compared to the pumping time.

(c) Ground triplet to excited triplet state absorption is not negligible, but the population of the excited triplet state is negligible because of rapid nonradiative decay.

(d) The dye laser has essentially only three levels. (The fine vibrational and rotational levels in each main state are neglected.)



(e) The effect on state population densities caused by dye self-absorption of the lasing wavelength is negligible. (There are approximately two orders of magnitude difference between the absorption coefficient for the pumping wavelength and the expected lasing wavelength so this assumption is not too restrictive.)

(f) The point of maximum gain is reached so fast that the effect of reduction of the excited singlet state due to stimulated emission can be neglected. (This assumption has been shown to be realistic by Atkinson and Pace [34] and is very true at the onset of lasing; however, as lasing progresses and photon buildup occurs in the laser cavity the stimulated emission term can become very significant. The same is true of the self-absorption term since it also depends on intracavity power.)

Equations (9), (10A), (10B), and (10D) are used to evaluate the dye laser model in a computer simulation in Section IV. Equation (9), the gain equation, is used in the discussion of theoretical results to explain the theoretical time dependent wavelength sweep of the dye laser.



### III. EXPERIMENTAL PROCEDURES AND RESULTS

#### A. EXPERIMENTAL SET UP

The apparatus used in generating and measuring the output of the Argon and dye lasers was set up as shown in Figure 3. A Beck Wavelength Reversion Spectroscope with a single line resolution of approximately four Angstroms was also used to monitor the dye laser output wavelength.

The beam splitter indicated in Figure 3 was actually an integral part of the front end of the Argon laser and is normally used as part of a power monitoring system. This system was modified so that the Argon laser output pulse could be monitored on an oscilloscope via a photodiode detector.

The diode detector and power meter were both tested for linearity and saturation at higher powers than could be expected at their locations in Figure 3. They both exhibited no signs of approaching saturation and their deviations from linearity over the testing range were minimal.

The risetime of the Argon laser is normally limited to seven nanoseconds because of cable transit times, etc., however, in the particular set up of Figure 3 the risetime was limited to 12 to 15 nanoseconds because that was the lower limit of the pulse generator risetime. An explanation of the cavity dump mechanism can be found in Maydan [35].

#### B. EXPERIMENTAL PROCEDURES

The first step in the procedure was the alignment of the Argon laser cavity dump system. This system consisted of a curved front half-mirror with vertical and horizontal adjustments, a rear curved mirror with vertical and horizontal plane tilt controls, and an acousto-optic cell housing



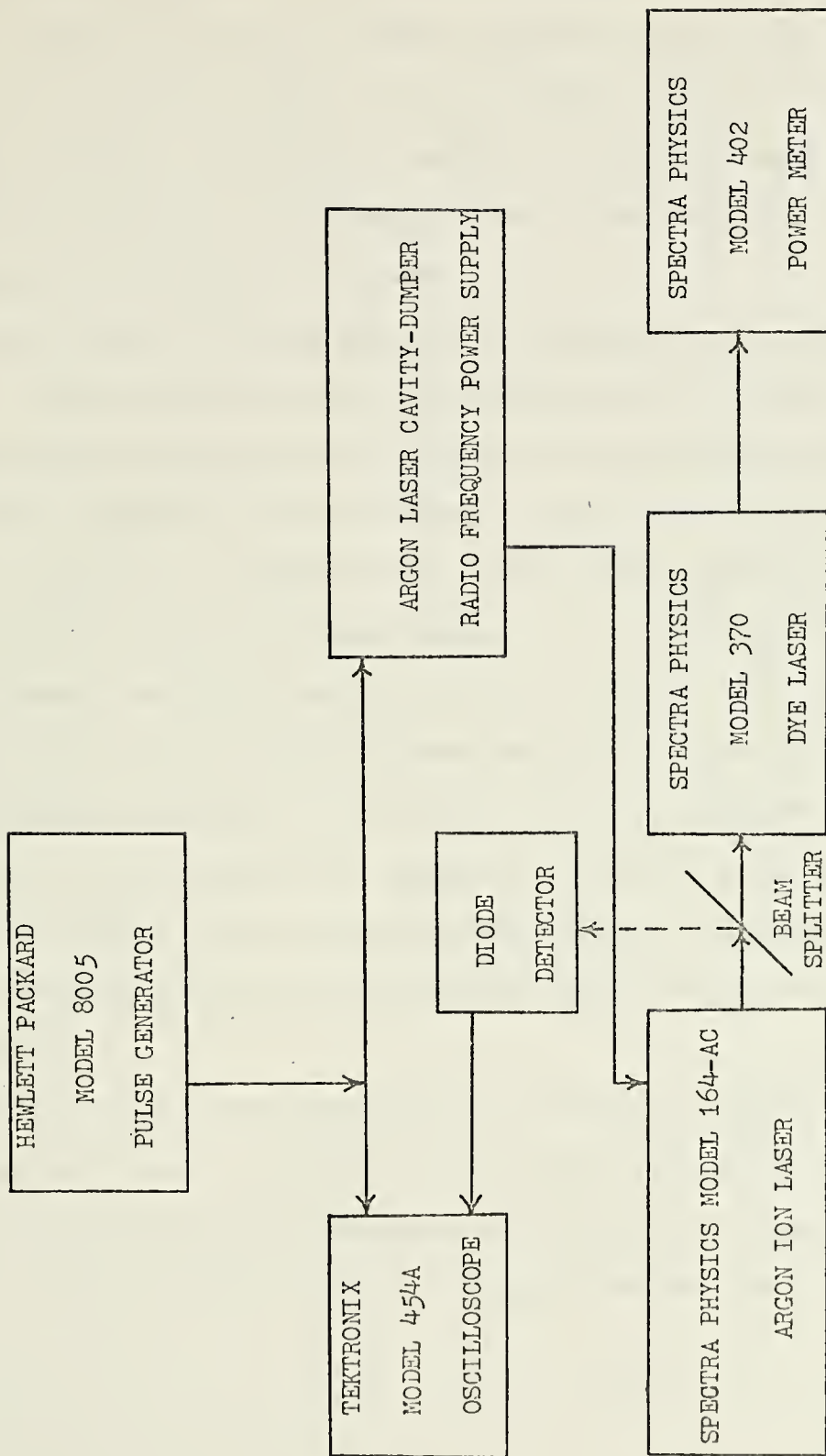


Figure 3: Experimental set up.





with controls to adjust its longitudinal and vertical positions as well as its angular alignment. All of these controls had to be precisely set for the cavity dump system to function properly, and the Argon laser power output and pulse shape had to be constantly monitored because some of the adjustments were extremely delicate. A slight movement of the rear mirror adjustments, for example, severely degraded the Argon laser output pulse and beam mode structure.

The next step in the procedure was the alignment of the Argon and dye lasers. This was done on a specially designed platform which had permanent hold down clamps for the Argon laser and kinematic mounts for the dye laser. Alignment of the two lasers was very sensitive and was a critical part of the experimentation since any nonparallelism of the dye and Argon laser cavities severely reduced the power conversion efficiency.

The Argon laser power supply was set for a current of 30 amperes after alignment was achieved. With the pulse generator delivering three volt, thirty nanosecond pulses at a repetition rate of one megahertz to the cavity-dump rf power supply, the average power output of the Argon laser was 500 milliwatts. These values and settings were then used as Argon laser operating reference points and standards for the rest of the experimentation.

The final step in the procedure was to monitor the output of the Argon and dye lasers. The Argon and dye laser powers were recorded, the Argon laser pulse shape was sketched, and the various dye laser wavelengths noted as the pulse width of the pulse generator was varied.



### C. EXPERIMENTAL RESULTS

The results yielded some new and interesting properties of the dye laser. The wavelength of peak output power and efficiency of the cavity-dumped laser-pumped dye laser was found to occur at a shorter wavelength than that of the continuous wave laser-pumped dye laser. In addition, the tuning range of the pulsed dye laser was narrowed.

The normalized output for both the continuous-wave laser pumped and pulsed-laser pumped dye laser are shown for comparison in Figure 4. The same dye laser was used to obtain both curves.

A two piece mathematical model which is an excellent fit to the actual Argon laser output pulse used to obtain the pulsed dye laser curve is shown in Figure 5. The repetition rate of this pulse was one megahertz.

It was also discovered that as the pulse width driving the cavity-dumped rf power supply was lengthened the wavelength of peak emission shifted toward longer wavelength and the laser efficiency decreased. In addition, for a given pulse width, an increase in the output power of the Argon laser (accomplished by increasing the current to the Argon laser power supply) shifted the peak emission wavelength to shorter wavelengths.

The lowest peak wavelength observed was 5650 Angstroms and the highest conversion efficiency obtained was almost thirty percent. This same dye laser pumped with a continuous wave laser had an efficiency of only seventeen percent at the peak output wavelength of 5800 Angstroms.

The fluorescence spectrum of the dye in the laser was also measured by use of a spectroscope with a high intensity white light source. This spectrum along with the wavelength-power curve of the cavity-dumped laser-pumped dye laser are compared in Figure 6. It is obvious from this figure that the peak lasing wavelength for the pulsed pump is approaching the peak of the dye fluorescence curve.



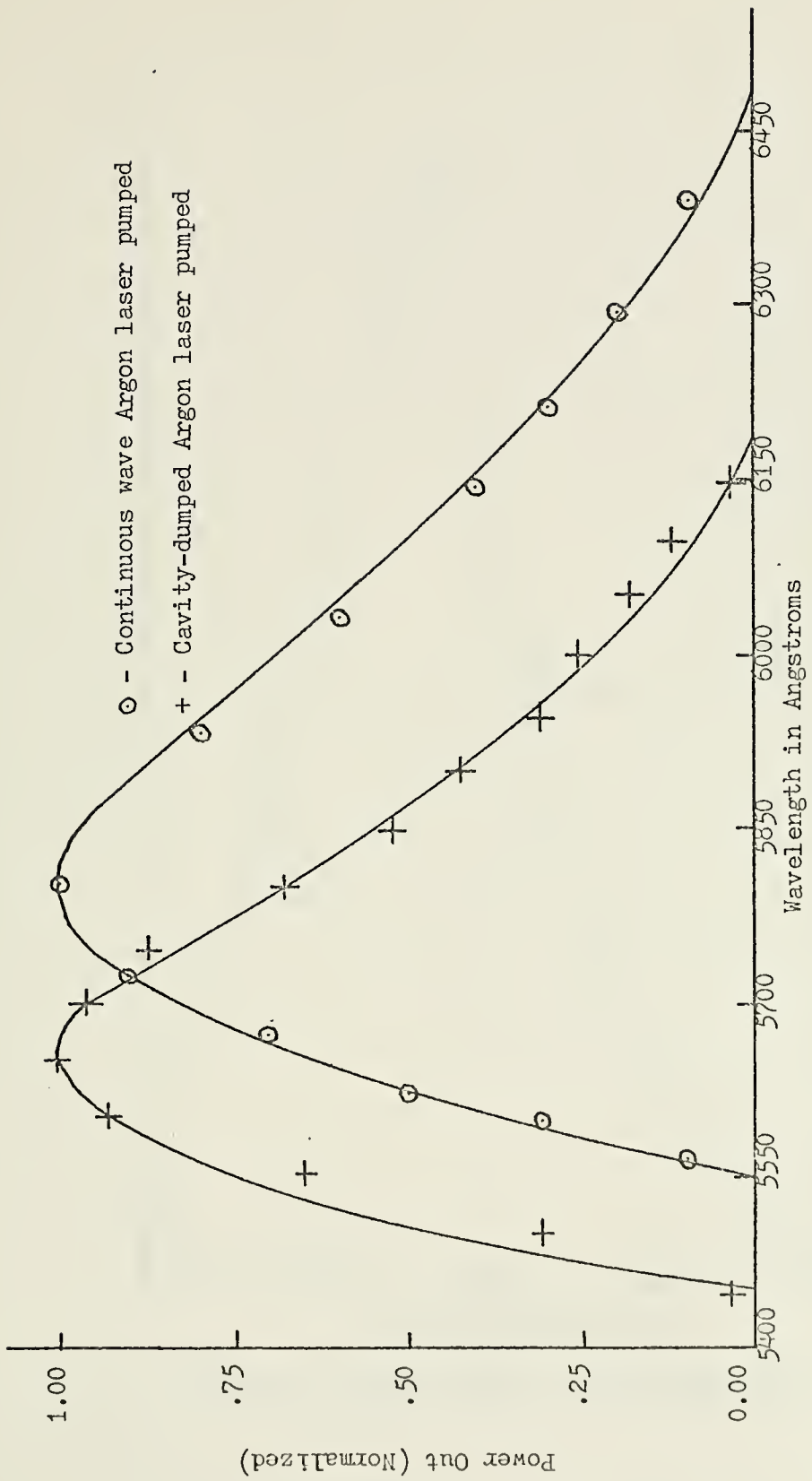


Figure 4: Power - wavelength spectrums of a Rhodamine 6-G dye laser.



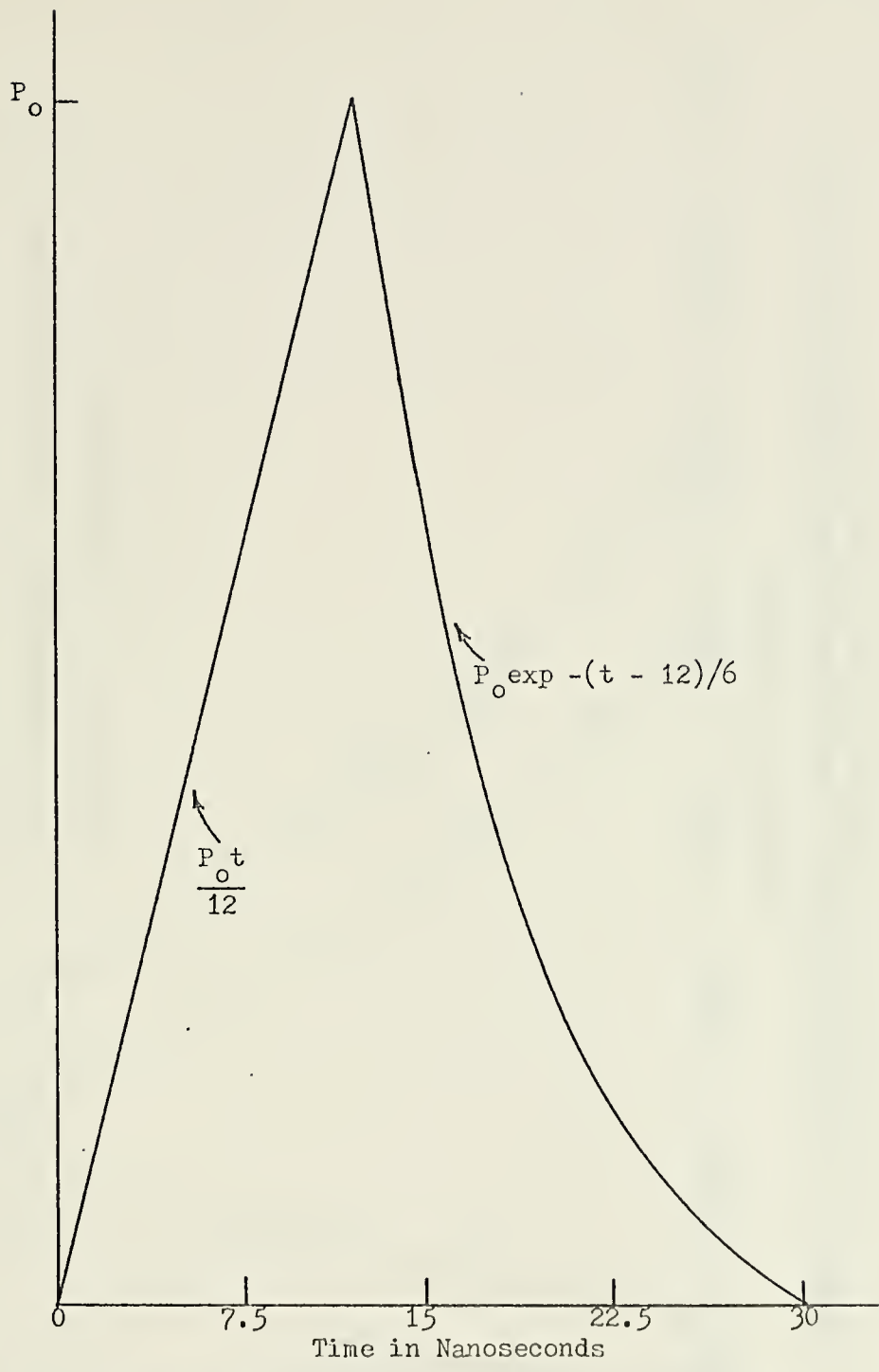


Figure 5: Typical Argon laser output pulse.





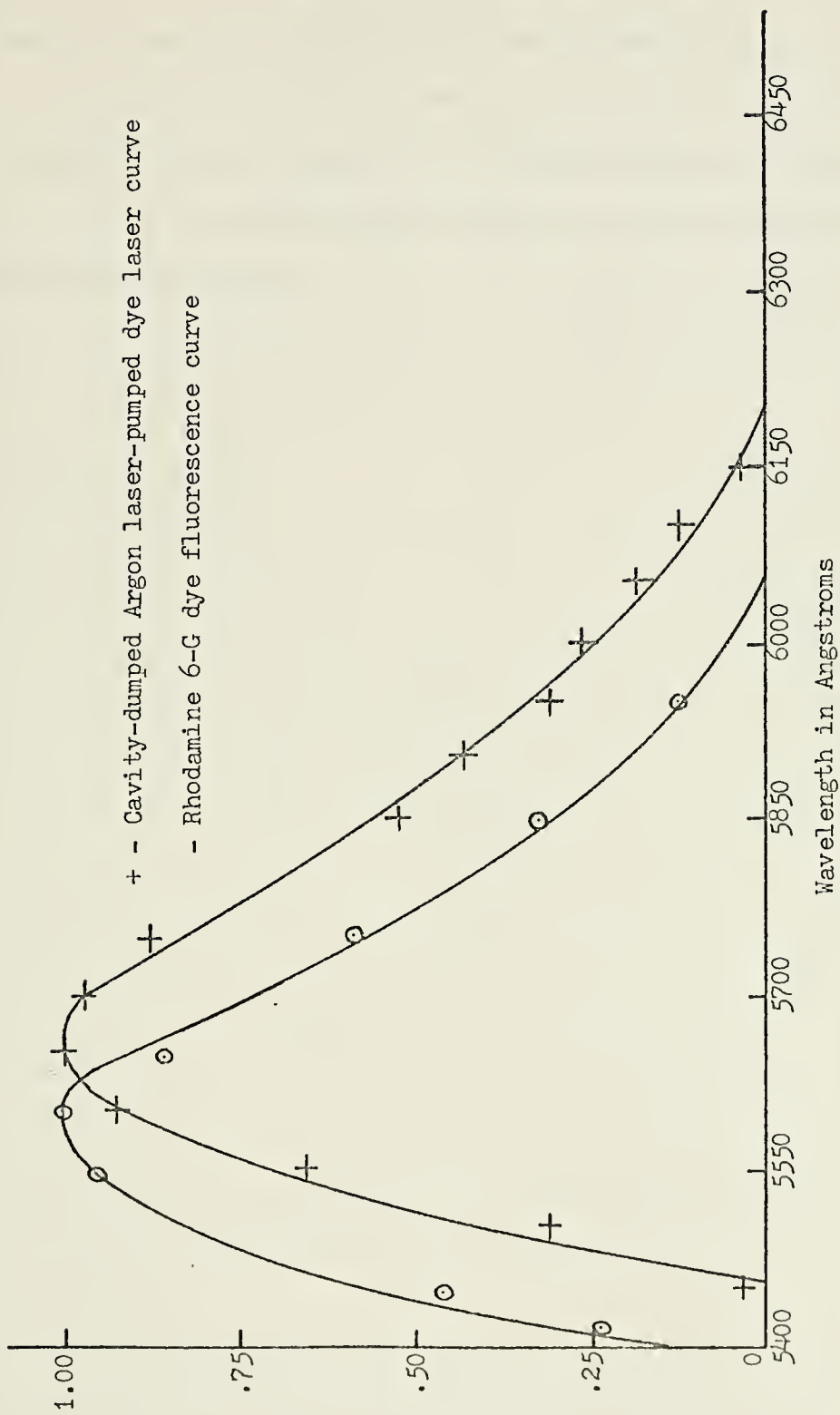


Figure 6: Wavelength-power curve of cavity-dumped Argon laser-pumped dye laser and Rhodamine 6-G fluorescence curve (Normalized).



A summary of the significant experimental observations is as follows:

(1) The pulsed dye laser wavelength spectrum exhibits a substantial green shift from the spectrum of the CW pumped dye laser.

(2) The shift increases and approaches the dye fluorescence curve as the pulse length from the cavity-dumper gets shorter.

(3) The efficiency is greater and the tuning range shorter for the pulsed dye laser.



#### IV. DYE LASER MODEL EVALUATION

Experimental results showed that the output wavelength spectrum of the cavity-dumped laser-pumped dye laser approached the dye fluorescence curve. The dye laser model of Figure 2 is analyzed in this section using equations (9), (10A), (10B) and (10D), which are the model gain and rate equations, and the pumping waveform of Figure 5 to determine if the spectrum shift can be predicted.

##### A. LITERATURE SURVEY

The first application of applying rate and gain equations to the dye laser was made by Sorokin, et al. [36, 22]. Their application assumed a Gaussian waveform pumping pulse and they attempted to predict the transient pulse shape of a pulsed Chloro-aluminum phthalocyanine (CAP) dye laser. They concluded from their results that a rapid pulse risetime was necessary or triplet state accumulation could prevent dye laser action.

The second analysis of Sorokin and his group again assumed a Gaussian pumping pulse and attempted to predict the efficiency of various dye lasers as functions of quantum efficiency and the excited singlet to triplet intersystem crossing rate constant ( $k_{st}$ ). The results from this work yielded the prediction that dye laser efficiency should increase as the intersystem crossing rate time constant decreased or as the quantum efficiency increased.

A year later, in 1968, Bass, Deutsch and Weber [37] used the rate and gain equations to predict lasing frequencies and time to lasing for Gaussian shaped laser and flashlamp pumped dye lasers. They had



experimentally observed that a laser pumped dye laser emitted a shorter wavelength than when flashlamp-pumped, however, different laser cavities were used for the two pumping conditions and this led to different parameters entered in the gain equation and to the conclusion that differing cavity Q's accounted for the higher lasing wavelength of the flashlamp pumped case.

The following year Bass and Weber [38] again used dye laser rate and gain equations. This time they assumed a rather long Gaussian shaped pumping pulse and included triplet state effects in their discussions. They still concluded that cavity Q was the main determinant of lasing wavelength, however, they conceded that determination of lasing wavelength could be more complex if triplet state effects were significant.

In mid 1970 Keller [39] solved for pseudo steady state solutions to the rate equations (all time derivatives were set equal to zero) to investigate the effects of quenching agents on singlet and triplet state populations. He arrived at the conclusion that a specific quencher for the triplet state would markedly improve laser efficiency and that a substance that quenched both singlet and triplet states would usually improve the efficiency.

The next group to use rate equations was McColgin, et al. [33]. They considered triplet state effects but only under steady state conditions as is found in some continuous wave dye lasers, and they concluded that the emission wavelength adjusts itself so that the ratio of mirror losses, scattering losses, etc. to singlet absorption losses remains approximately constant.

Pappalardo, Samuelson and Lempicki [40] in 1972 used the rate equations to calculate the efficiency of dye lasers as a function of pump





parameters and triplet state lifetime for pump pulses in excess of one microsecond in duration. They assumed a Gaussian shaped pumping pulse and predicted that long pulse operation of up to 30 microseconds was feasible for short triplet state lifetimes but that the efficiency would go down as the pumping pulse was lengthened.

The most recent use of rate equations has been by Strome and Tuccio [17] who used results to improve their original dye laser [16], by Atkinson and Pace [34] who used them to calculate the lineshape of a Fabry-Perot etalon tuned Rhodamine 6-G dye laser, and by Streifer and Saltz [41] who used the equations to analyze an acoustooptically tuned dyelaser.

## B. GAIN AND RATE EQUATION PARAMETERS

The evaluation of the dye laser system rate and gain equations for a cavity such as the one in Figure 2 requires a lot of physical data. The actual dye used to obtain the experimental data had a mixture of methanol and water as the dye solvent; however, there is no information available on the dye in this particular solvent hence physical data available for ethanol solutions was used. It was assumed that any predictions or conclusions based on the results would apply to Rhodamine 6-G in the actual solvent. (The main difference in the solvents is that both the singlet absorption and fluorescence spectrum peaks are further in the green region of the visible spectrum when the solvent is 100 percent ethanol.)

The fluorescence lineshape and singlet absorption cross section of a  $10^{-4}$  Molar solution of Rhodamine 6-G in ethanol are shown in Figure 7. These curves were obtained by F. Grum of Kodak Research Laboratories and were reproduced in Reference 8. Snavely [8] measured the fluorescence yield of the dye and obtained a value of  $\phi = 0.83$  and the fluorescence spectrum of Figure 7 is normalized to this value; i.e.



$$\int_{-\infty}^{\infty} E(\lambda) d\lambda = \phi = 0.83.$$

There is no data available on the triplet absorption cross section for an alcohol solution of Rhodamine 6-G. The only available curve is shown in Figure 8 which was determined by Buettner [42] using a flash photolysis technique for Rhodamine 6-G in polymethyl methacrylate.

Several values of singlet and triplet state lifetimes are given in the literature. The various values given for  $T_s$  are 7.4 nsec in Weber and Bass [38], 4.8 nsec in McColgin, et al. [33], 7.3 nsec in Snavely and Peterson [31], and 5.5 nsec in Mack [43]. The reason for the use of different values by the groups mentioned above is due to different measured values. An average value of 6.3 nsec was used in this analysis because there were no significant reasons why one value should be preferred.

There are also many values of  $T_t$  given in the literature, but the best evidence available suggests that the most reasonable value is 100 nsec. Keller [39] arrived at this value in his analysis of Oxygen quenching in an alcohol solution and Snavely and Schafer [9] obtained 100 nsec as a measured value.

The value of the intersystem crossing rate constant,  $k_{st}$ , can be determined from the formula

$$k_{st}^{-1} = T_s \phi / (1 - \phi) \quad (11)$$

which is found in Snavely [8] and others. The value of 0.84 from Bass and Steinfeld [44] is the most common value found for  $\phi$  and since the value of 6.3 nsec is assumed for  $T_s$ , (11) yields the value of 29.4 nsec for  $k_{st}^{-1}$ . The fluorescence spectrum of Figure 7 is normalized using  $\phi = 0.83$  so that spectrum and equation (11) are compatible.



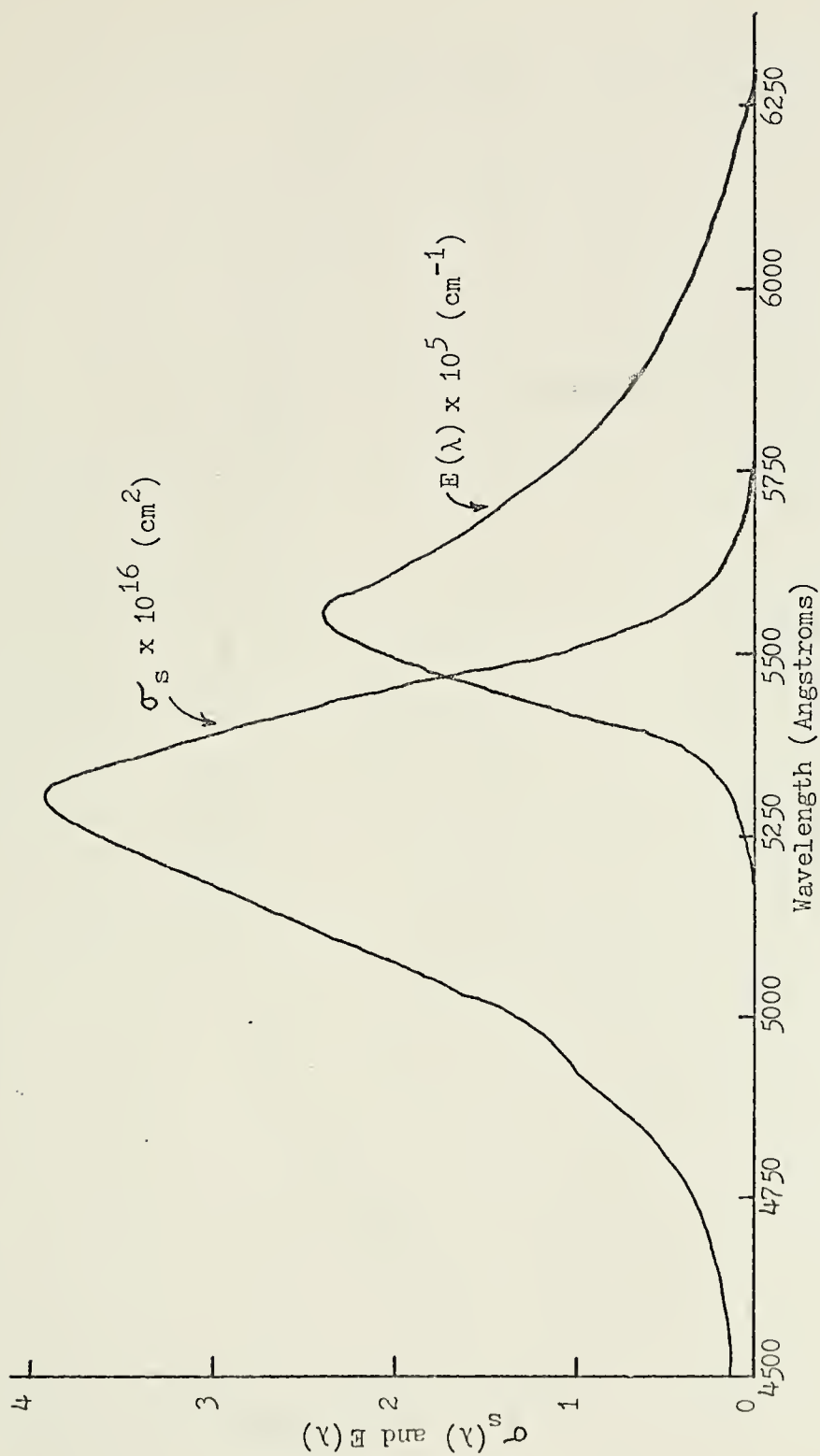


Figure 7: Singlet state absorption cross section and fluorescence spectrum of  $10^{-4}$  Molar Rhodamine 6-G in ethanol [8].



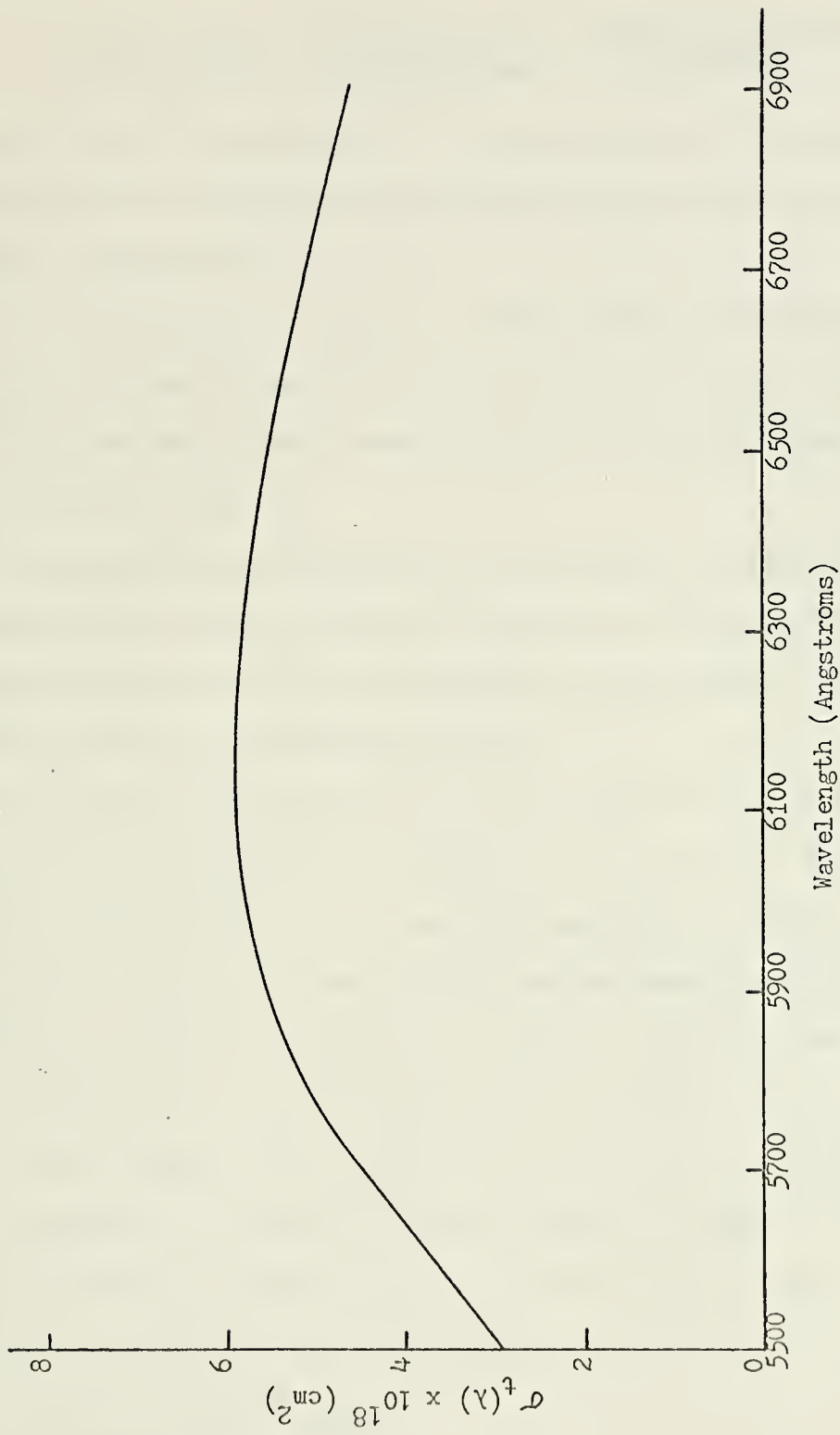


Figure 8: Triplet state absorption coefficient of Rhodamine 6-G in polymethyl methacrylate [42].





The easiest constant to evaluate is  $N$ , the total molecular density. The actual dye solution was  $3 \times 10^{-4}$  Molar so that

$$N = 3 \times 10^{-4} \text{ Molar} \times \frac{1 \text{ gram-mole}}{1 \text{ Molar-liter}} \times \frac{6.023 \times 10^{23} \text{ molecules}}{\text{gram-mole}} \times \frac{1 \text{ liter}}{10^3 \text{ cm}^3}$$

which is  $1.8 \times 10^{17}$  molecules/cm<sup>3</sup>. This assumes that the density of the solvent at room temperature is approximately the same as at the standard reference temperature.

The values of  $L_1 = .1$  cm,  $R_1 = .995$  and the curve of  $R_2$  versus wavelength were obtained from the manufacturer [45]. The dye cell windows were anti-reflection (AR) coated so typical values for AR surfaces of .98 were used for  $T_1$  and  $T_2$ .

The remaining constants needed for evaluation of  $\sigma_e(\lambda)$  are  $c$ , the speed of light, and  $\bar{n}$ , the index of refraction. The value used for the speed of light was  $3 \times 10^{10}$  centimeters/sec and the index of refraction of the dye solution was approximately 1.33.

The optical pumping rate,  $P(t)$ , may be written as  $\sigma_p I_p / hf_p$  where  $\sigma_p$  is the dye absorption cross section for  $f_p$ , the pump frequency,  $h$  is Plancks constant and  $I_p$  is the pump irradiance in watts/cm<sup>2</sup>. The pump wavelength was 5145 Angstroms and the corresponding frequency was  $5.82 \times 10^{14}$  sec<sup>-1</sup>. Plancks constant has a value of  $6.625 \times 10^{-34}$  sec<sup>-1</sup>. The value of  $2.37 \times 10^{-16}$  cm<sup>2</sup> obtained from Figure 7 was used for  $\sigma_p$ .

Pump irradiance can be written as power per unit area. The model of the pump power as a function of time is shown in Figure 5. The average power output of the Argon laser was 0.5 watts and the value of 42.7 watts for  $P_0$  was obtained from the formula

$$P_{\text{average}} = \frac{1}{T} \int_0^T P(t) dt.$$



There is a problem in defining the effective area and volume in which the pump beam is absorbed. Jacobs, Samuelson, and Lempicki [46] assumed uniform pumping in their analysis of losses in a continuous wave dye laser and used the beam waist area as the effective pumped area. At least some of the Argon laser beam must exit the dye volume before uniform pumping can be assumed [40], otherwise the amount of beam penetration into the dye and the size of the pumped volume cannot be determined. The dye laser used in the experiment allowed approximately five percent of pump power to exit the dye cell [45], therefore, uniform pumping was assumed. The beam waist radius was six microns [45] so the assumed effective area was  $1.13 \times 10^{-6} \text{ cm}^2$ .

The assumption of uniform pumping is one of the most marginal in dye laser analysis. The complexities introduced by trying to evaluate mode-matching of the pump and dye laser beams in the dye cell and nonuniform excitation of dye molecules, however, makes the assumption practical.

## C. GAIN AND RATE EQUATION SOLUTIONS

### 1. Solution Method

The first step in the theoretical analysis was the solution of the rate equations for the state population densities as a function of time. The solution was carried out using a program entitled INTEGL of the U.S. Naval Postgraduate School Computer Library. (See Appendix A) INTEGL uses a fourth order Runge-Kutta method with programmable step size changes to achieve the solution to simultaneous ordinary differential equations. The only modifications made to the INTEGL program were those necessary to enable it to handle the wide range of parameter values used in the rate equations.



The values of the state population densities were then used to evaluate the gain equation. This was accomplished by first selecting values of state densities at a particular time and then evaluating the gain using the wavelength dependent parameters over a large range of wavelengths. The succeeding steps followed the same pattern and this was continued until the point of maximum gain versus time and wavelength reached a constant relationship; i.e., the wavelength of maximum gain remained unchanged with time.

The purpose of evaluating the gain in this manner was that the lasing wavelength of the laser model should be where the gain is maximum. This was then compared to the dye fluorescence spectrum to see where the lasing wavelength was in relation to the peak wavelength of the spectrum.

The last step was to predict where the experimental laser should lase based on the theoretical comparison and the fluorescence spectrum of the experimental dye. Theoretical and experimental results should agree if theory was sound.

## 2. Theoretical Results

A graph of the peak gain wavelength versus time is shown in Figure 9. It can be seen from the figure that lasing starts at the longer wavelengths and sweeps toward the shorter wavelengths. Figure 9 also contains a graph of the ratio of triplet state population to ground singlet state population ( $N_t/N_o$ ) versus time.

The results of evaluation of the gain versus wavelength at steady state is shown in Figure 10. It can be seen that the peak wavelength (5570 Angstroms) corresponds to the peak wavelength of the Rhodamine 6-G fluorescence curve of Figure 7.

The results displayed in Figure 9 and the gain equation, equation (9) of Section II, can be used to explain the theoretically obtained dye



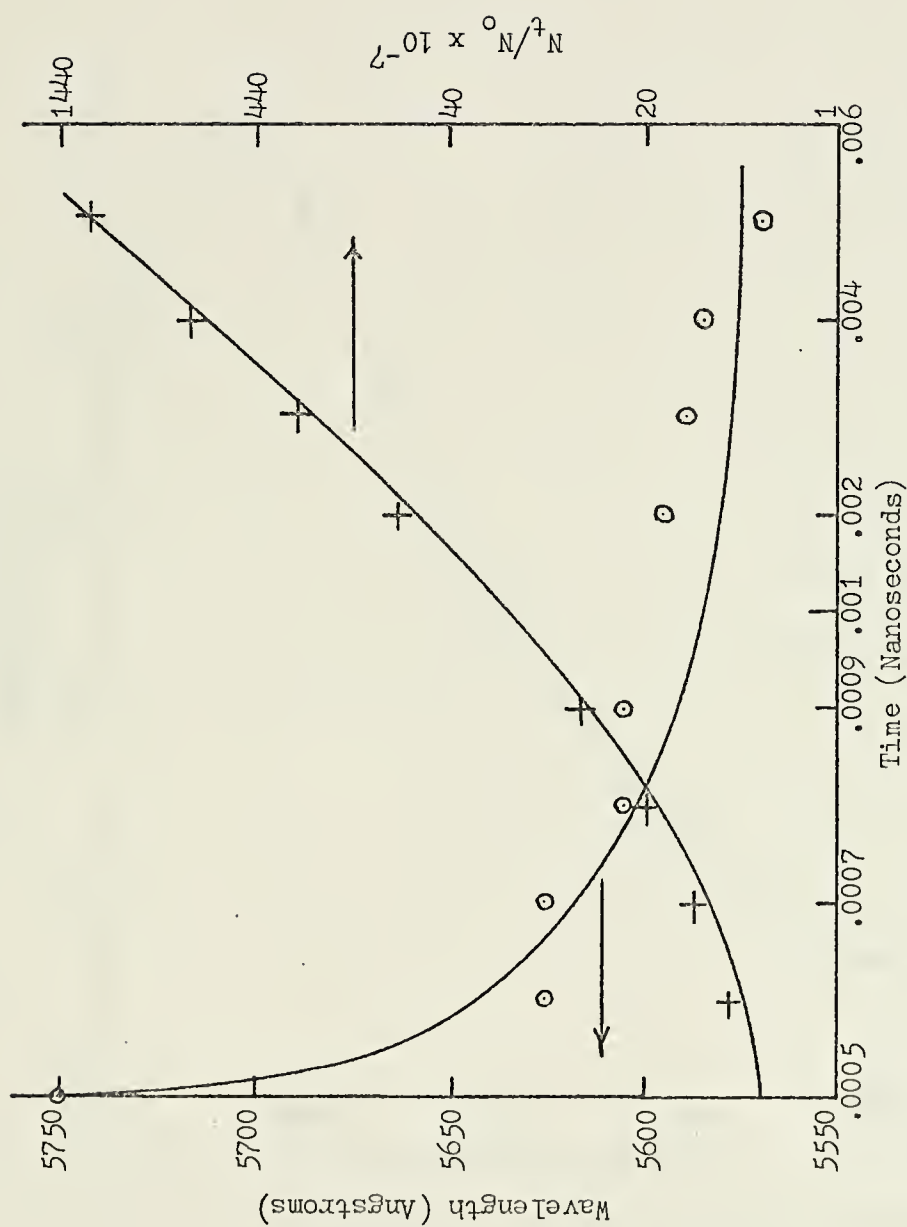


Figure 9: Peak gain wavelength and ratio of triplet to ground singlet state population for  $3 \times 10^{-4}$  Molar Rhodamine 6-G laser model.





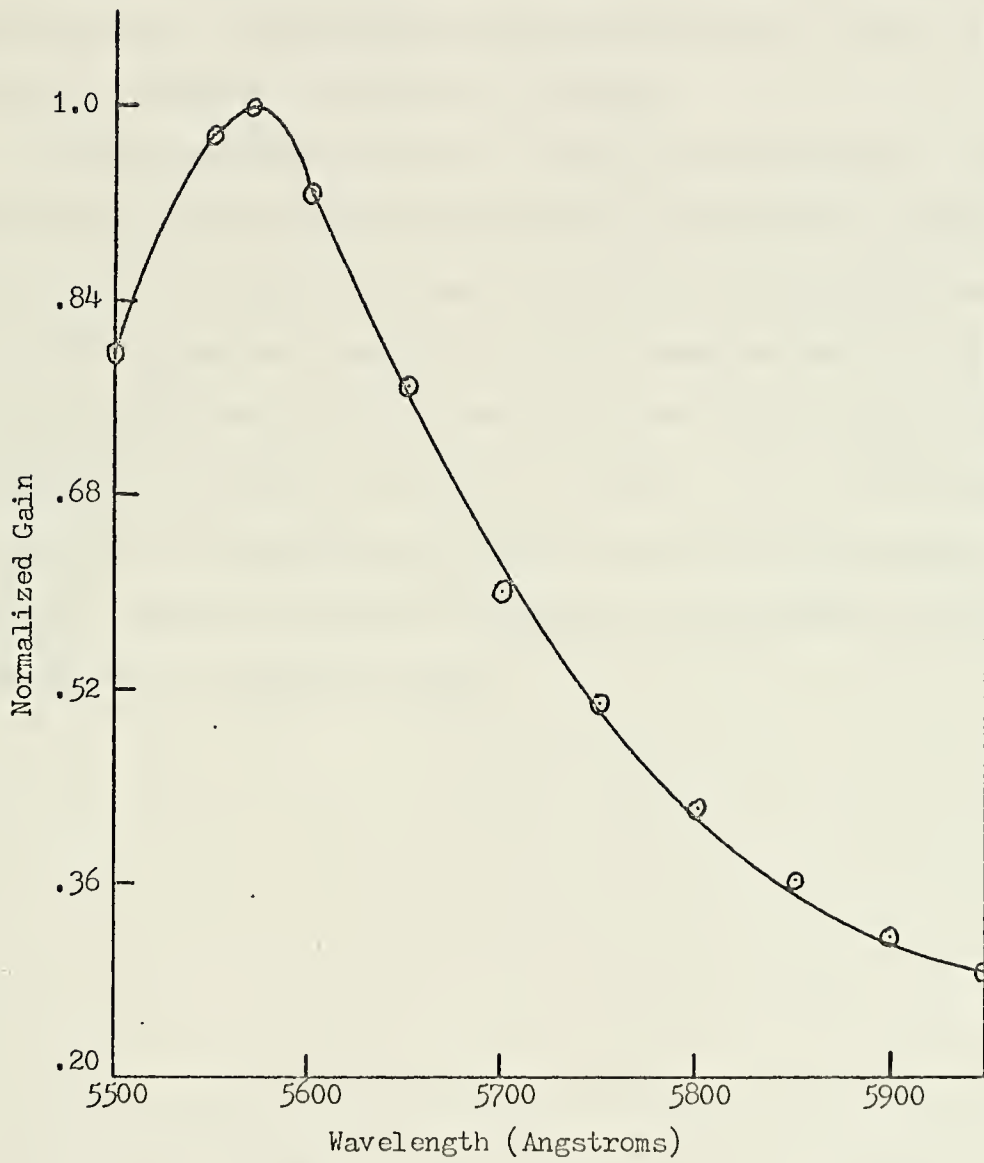


Figure 10: Theoretical normalized gain versus wavelength for Rhodamine 6-G dye laser model.



laser wavelength sweep. The two wavelength dependent losses in the gain equation are caused by singlet and triplet absorption. The ratio of  $N_t/N_o$  at the onset of lasing is very small indicating domination of the singlet state absorption loss. (Singlet absorption is associated with  $N_o$  and triplet absorption with  $N_t$ .) Figure 7 shows that singlet absorption is greatest for shorter wavelengths, hence there should be a tendency for long lasing wavelengths at the onset of lasing.

As pumping continues the ratio of  $N_t/N_o$  becomes greater. This indicates that as time passes there should be a tendency for triplet absorption losses to dominate. Figure 8 shows that triplet absorption is greater at the longer wavelengths (up to 6000 Angstroms which is within the lasing range of Rhodamine 6-G) and, therefore, the dye laser should end up lasing at wavelengths closer to the peak of the dye fluorescence curve. (The above statements are true in this case for the specific pulse model of Figure 5. Pulses of different shape, duration, and intensity may or may not cause the sweep.)



## V. CONCLUSIONS

### A. EXPERIMENTAL AND THEORETICAL COMPARISON

The theoretical analysis predicted that for the pumping pulse model of Figure 5 and after a wavelength sweep of very short duration the dye laser would lase at or near the peak of the dye fluorescence spectrum. It was observed experimentally that the dye laser lased approximately 50 Angstroms from the peak of the dye fluorescence curve in the pulsed case (Figure 6) rather than 200 Angstroms as in the continuous wave pumped case (Figures 4 and 6). It may be concluded, therefore, that theory and experimentation are in agreement within the limits of the assumptions made in the study.

The variable of least confidence in the entire analysis was the experimental dye fluorescence curve. The peak of the fluorescence curve for Rhodamine 6-G in an ethanol solution is typically about 5570 Angstroms [8] and in a water solution is about 5750 Angstroms [46]. The actual solvent was determined by nuclear magnetic resonance measurements to be approximately two-thirds water and one-third methanol. This should have caused the fluorescence peak of the dye to be somewhere around 5660 to 5700 Angstroms which is much closer to the observed peak wavelength than the 5600 Angstrom peak of the fluorescence curve of Figure 6. One explanation for the difference between the measured and expected fluorescence peaks is the fact that the spectrometer calibration curve used to measure the experimental dye fluorescence spectrum had a steep slope and was very sensitive to spectrometer dial readings.



## B. REMARKS

The decrease in efficiency of the dye laser as the pulse length was widened can be explained very simply in terms of the rate equations. The lengthening of the pulse decreased the peak power of the pumping pulse which in turn caused a longer time to reach the onset of lasing. As a result, the triplet state had more time to build up and extract molecules from the singlet lasing system - efficiency had to decrease. There are several things which hinder the analysis of the dye laser. There is not much data available on dye properties in various solvents and the properties of the triplet state are still pretty much unknown. There is also a lot to be learned about the effects of thermal heating and nonuniform pumping of the dye and no one as yet has begun to investigate the effect of chemical reaction and relaxation rates on the dye laser. Stability and noise considerations also require investigation.





## APPENDIX A

### DYE LASER RATE EQUATION COMPUTER PROGRAM LISTING

The computer listing on the following pages contains the program for the solution of the dye laser rate equations. The main part of the program is SUBROUTINE INTEGL and the comments preceding its operational statements completely describe its use. The rate equations in the program correspond to equations (10A), (10B) and (10D) of Section II. The data cards used in the actual solution and described in the usage comments of INTEGL are located at the end of the listing.

INTEGL had to be modified to some extent so it could handle a wider range of parameter values than was originally allowed. These modifications included input and output format and an increase in the number of allowed integration steps and pages of output, but no changes were made to the basic fourth order Runge-Kutta integration and other computational algorithms.



DYE LASER RATE EQUATION SOLUTION

```

//SNYS2C83 JOB (2083,029C,DC12), 'SNYDER, G. W.', TIME=2
// EXEC FOR TCG, REGION.GC=100K
//FCRT DIMENSION X(30),XDCT(30),C(15)
C(10)=1.0
1 CALL INTEG4(T,X,XDCT,C)
IF(T.GE.0.001) GO TO 2
C(11)=1.0C
C(12)=1.0C
GO TO 3
2 C(11)=1
C(12)=1
IF(T.LE.0.1) GO TO 3
C(11)=1.0C
C(12)=1.0C
3 CONTINUE
Y = POWER(T)

```

CCCCCCCC

THE THREE EQUATIONS BELOW ARE THE DYE LASER RATE EQUATIONS AND CORRESPOND TO EQUATIONS 10A, 10B AND 10D OF SECTION III. IN THE EQUATIONS X(1) IS THE EXCITED SINGLET STATE POPULATION, X(2) IS THE TRIPLET STATE POPULATION AND X(3) IS THE GROUND SINGLET STATE POPULATION.

$$\begin{aligned}
 XDCT(1) &= -C(2)*X(1) + C(6)*C(3)*C(7)*Y - .542*Y*X(1) - \\
 .542*Y*X(2) \\
 XDCT(2) &= -C(4)*X(2) + C(5)*X(1) \\
 X(3) &= 18.0*C(7)*C(7)*C(7) - X(1) - X(2) \\
 GO TO 1 \\
 END
 \end{aligned}$$

CCC

FUNCTION POWER(Z) BELOW IS THE ARGON LASER POWER AS A FUNCTION OF TIME IN ACCORDANCE WITH THE TWO SEGMENT CURVE OF FIGURE 5.

```

FUNCTION POWER(Z)
IF(Z.GE.12.0) GO TO 10
POWER = (42.7*Z)/12.0
10 POWER = 42.7*EXP((12.0-Z)/6.0)
RETURN
END

```

1C







EQUATION STATEMENTS

THE NUMBER OF EQUATIONS TO BE SOLVED, "N", MUST BE LESS THAN OR EQUAL 30. THEY MUST BE REDUCED TO STANDARD FORM AND THEN EXPRESSED AS WELL AS ADDITIONAL NON-SUBSCRIPTED VARIABLES MAY BE USED AS STANDARD LIBRARY FUNCTIONS AND/OR SUBROUTINES. THE ORDER OF THE EQUATIONS IS USUALLY UNIMPORTANT. A TYPICAL SET MIGHT READ:

```

ERROR=SIN(C(1)*T)-X(1)
XDOT(1)=X(2)-.1*X(1)
XDOT(2)=ERROR+C(2)*XDOT(1)

```

THE X ARRAY CONTAINS THE DEPENDENT VARIABLES. THE XDOT ARRAY CONTAINS THE FIRST DERIVATIVE WITH RESPECT TO THE INDEPENDENT VARIABLE ("T" IN THIS EXAMPLE). THE XDOT ARRAY BY THE USER DEFINE "XDOTI" IN TERMS OF THE INDEPENDENT AND DEPENDENT VARIABLES, CONSTANTS AND OTHER VALUES IN THE "XDOT" ARRAY.

CONSTANTS (C(1), I, LE, 8) ARE READ FROM A DATA CARD. CCN (C(10)) MUST NEVER BE USED EXCEPT AS IN THE SAMPLE DECK SETUP ABOVE. CCN(11), C(12), AND C(13) CONTROL FREQUENCY OF PRINT-OUT DURING INTEGRATION, THE FREQUENCY OF THAT POINTS GENERATED, AND THE PRINT-OUT EVERY STEP AND STEP-SIZED IS GENERATED AT EACH STEP. THE INSTANTANEOUS STATE DEFINED ON A DATA CARD. THE CCN(11) AND "GO TO 1" STATEMENTS REDEFINING THEM BETWEEN THE "CALL INTEGRATE" AND "GO TO 1" STATEMENTS ABOVE (AT ANY POINT DURING THE SOLUTION) SUCH AS:

```

C(11)=10.
C(12)=1.

```

CAUSES PRINT-OUT TO TAKE PLACE AT EVERY 10 STEPS AND PLOT POINTS TO BE GENERATED AT EACH STEP. NOTE THAT MORE THAN 4500 INTEGRATION STEPS, 450 LINES OF PRINT-OUT, OR 900 PLOT POINTS CAN BE GENERATED IN ANY ONE RUN.

ANY LEGAL FORTRAN STATEMENT OR TECHNIQUE CAN BE USED IN THE CALLING PROGRAM, PROVIDED THAT A TRANSFER OF CONTROL DOES NOT PREVENT THE PROCESSING OF ALL RELEVANT STATEMENTS DURING EACH STEP OF THE INTEGRATION. FOR EXAMPLE THE FOLLOWING WOULD BE ALLOWED:

```

IF(T-10.0) 3,2,2
3 C(1)=0.
2 CONTINUE

```

CC









TIONAL CARDS ARE REQUIRED IF N.GT.8.

NEXT TO  
LAST:

CONTROLS THE PRINT-OUT. A BLANK CARD WILL SUPPRESS THE  
CUT-OUT UP TO 8 QUANTITIES. EACH QUANTITY IS A 2-CHARACTER  
INTEGRATION COLUMN BEING PRINTED UNDER THE VARI-  
THE VARIABLE IS PUNCHED IN COLUMNS 1-8 AND COLUMNS  
COLUMN HEADING IS PUNCHED IN COLUMNS 9-10. THE  
PCNDING SUBSCRIPT IN COLUMNS 11-18 AND COLUMNS  
ING IS PUNCHED IN COLUMNS 19-20. THE SUBSCRIPT  
CONTINUE IN THIS MANNER UP TO A MAXIMUM OF 8 QUANTITIES.  
THE SUBSCRIPT VARIABLE (OR BLANK) IS USED TO SPECIFY THE  
INDEPENDENT VARIABLE, I.

LAST:

CONTROLS THE PLOT OUTPUT. A BLANK CARD WILL SUPPRESS  
THE OUTPUT. UP TO 4 CURVES CAN BE PLOTTED SEPARATELY.  
BE PLOTTED SPECIFIED BY A PUNCH IN COLUMNS 1-16. CHAR-  
CURVE IS IDENTIFIED BY THE SUBSCRIPT (2 DIGITS), AND  
ACTERS, SUBSCRIPT OF THE X-COORDINATE (2 DIGITS), AND  
TITLE IS PUNCHED IN COLUMNS 17-18, AND THE ORIGINATE  
COLUMNS 19-20. THE FIRST Y-COORDINATE IS IN  
SECOND COLUMN AND THE SECOND Y-COORDINATE IS IN  
COLUMNS 37-38, AND THROUGH 4 GRAPHS. WHEN ALL  
AND SO ON GRAPH, SPECIFY ONLY THE FIRST TITLE AND LEAVE  
ON ONE GRAPH, SPECIFY ONLY THE FIRST TITLE AND LEAVE  
THE OTHER TITLE FIELDS BLANK.

INTEGRAL USES SUBROUTINES PLOTP AND UTPLCT TO GENERATE  
THE POINT PLOTTING.

MULTIPLE RUNS

IF SEVERAL SOLUTIONS OF THE SAME EQUATIONS ARE DESIRED, THE  
NUMBER OF RUNS IS SPECIFIED IN THE SECOND DATA CARD. FOR EACH  
RUN AFTER THE FIRST, ONLY THE FOURTH THROUGH THE LAST CARDS ARE  
REQUIRED. NO INFORMATION, BESIDES THE FIRST 3 CARDS, IS  
RETAINED BETWEEN RUNS.

CAUTIONS TO USER

TOO LARGE A STEP-SIZE CAN RESULT IN UNACCEPTABLE INTEGRATION  
ERRORS AND EVEN INSTABILITY OF THE INTEGRATION PROCESS. THE  
ONLY (MODERATELY) SAFE TEST IS TO RUN A PROBLEM WITH AT LEAST  
TWO DIFFERENT STEP-SIZES AND THEN COMPARE THE RESULTS.  
VALUES OF STEP-SIZE LARGER THAN UNITY ARE USUALLY NOT  
RECOMMENDED.

114450  
114460  
114470  
114480  
114490  
114500  
114510  
114520  
114530  
114540  
114550  
114560  
114570  
114580  
114590  
114600  
114610  
114620  
114630  
114640  
114650  
114660  
114670  
114680  
114690  
114700  
114710  
114720  
114730  
114740  
114750  
114760  
114770  
114780  
114790  
114800  
114810  
114820  
114830  
114840  
114850  
114860  
114870  
114880  
114890  
114900









```

5000 FCRMAT(2F15.6)
      TF = TF1
      IF(DT2.NE.0.) GO TO 9
      WRITE(6,206) I, TF
      FCRMAT(22H INITIAL TIME
206 1 FCRMAT(22H FINAL TIME
      WRITE(6,207) DT
      FCRMAT(22H STEP SIZE
      GC TO 12
      IF(DT3.NE.0.) GO TO 11
      TF = TF2
      WRITE(6,206) II, TF
      WRITE(6,208) DT, II, TF1, DT2, TF1, TF2, TF
208 1 FCRMAT(22H STEP SIZE
      9H AND T = ,E19.11, 13H BETWEEN T = ,E19.11,
      GC TO 12
      TF=TF3
      11 WRITE(6,208) DT, II, TF1, DT2, TF1, TF2, DT3, TF2, TF
      12 WREAD(5,50C1) (C(I), I=1,8)
5001 FCRMAT(8F10.4)
      106 READ(5,10C6) (X(I), I=1,NN)
      J = 0
      DC 14 I=1,8
      IF(C(I).NE.0.) J=J+1
      CCNTINUE
      14 K = 0
      DC 16 I=1,NN
      IF(X(I).NE.0.) K=K+1
      CCNTINUE
      16 IF(J = 1) 17,18,19
      17 WRITE(6,209)
      209 FCRMAT(17,34H ALL THE CCNSTANTS, C(I), ARE ZER0 )
      GC TO 423
      18 WRITE(6,210)
      210 FCRMAT(17,30H THE ONLY NON-ZER0 CCNSTANT IS )
      GC TO 420
      19 WRITE(6,211)
      211 FCRMAT(17,35H THE NON-ZER0 CCNSTANTS, C(I), ARE )
      420 DC 422 I=1,8
      IF(C(I).NE.0.) WRITE(6,212) I,C(I)
      212 FCRMAT(14X,2HC(,12,4H) = ,E12.4)
      CCNTINUE
      423 IF(K = 1) 424,425,426
      424 WRITE(6,1209)
      1209 FCRMAT(17,36H ALL THE INITIAL CCNCITIONS ARE ZER0 )
      GC TO 20
      425 WRITE(6,1210)

```

```

INT12390
INT12400
INT1241C

```

```

INT12440
INT12460
INT12470
INT12480
INT12490
INT12500

```

```

INT12530
INT12540
INT12550

```

```

INT12580
INT12590
INT12600
INT12610
INT12620
INT12630
INT12640
INT12650
INT12660
INT12670
INT12680
INT12690
INT12700
INT12710
INT12720
INT12730
INT12740
INT12750

```

```

INT12780
INT12790
INT12800
INT12810
INT12820
INT12830

```





INT112840  
 INT112850  
 INT112860  
 INT112870  
 INT112880  
 INT112890  
 INT112900  
 INT112910

INT112940  
 INT112950  
 INT112960  
 INT112970  
 INT112980  
 INT112990  
 INT113000  
 INT113010  
 INT113020  
 INT113030  
 INT113040  
 INT113050  
 INT113060  
 INT113070  
 INT113080  
 INT113090  
 INT113100  
 INT113110  
 INT113120  
 INT113130  
 INT113140  
 INT113150  
 INT113160  
 INT113170  
 INT113180  
 INT113190  
 INT113200  
 INT113210  
 INT113220  
 INT113230  
 INT113240  
 INT113250  
 INT113260  
 INT113270  
 INT113280  
 INT113290  
 INT113300  
 INT113310

```

1210 FCRMAT (/ , 39H THE ONLY NON-ZERO INITIAL CCNCDITION IS )
      GC TO 427
426  WRITE (6, 1211)
1211 FCRMAT (/ , 36H THE NON-ZERO INITIAL CONDITICNS ARE )
427  CC 429 I=1, NN
      IF (X(I) .NE. 0) WRITE(6, 1212) I, X(I)
1212 FCRMAT (14X, 2HX( , 12, 4H) = , E25.5)
429  CCNTINUE
      READ (5, 104) (JTITL(I), IP(I), I=1, 8)
104  FCRMAT(8(A8, I2))

C     CHECK FOR THE NUMBER OF COLUMNS CALLED FOR BY LOCATING FIRST
C     BLANK CCLUMN HEADING
C
      CC 21 J=1, 8
      IF (JTITL(J) .EQ. IBLANK) GO TO 22
      CCNTINUE
      J = 9
      JJ = J - 1

C     JJ IS NOW THE NUMBER OF COLUMNS. REPEAT WITH THE GRAPHS.
C
      READ (5, 105) (KTITL(I), KTITL(I+1), IG(I), IG(I+1), I=1, 7, 2)
105  FCRMAT (4(2A8, 2I2))
      CC 24 K=1, 7, 2
      IF (KTITL(K) .EQ. IBLANK .AND. KTITL(K+1) .EQ. IBLANK) GC TO 25
      CCNTINUE
      K = 8
      KK = K/2
      KKK = KK*2
      MULTIP = 0
      IF (KK .NE. 1) GO TO 306
      IF (IG(3) .NE. 0) GO TO 306
      IF (IG(5) .NE. 0) GO TO 306
      MULTIP = 2
      KKK = 4
      GC TO 306
      IF (IG(7) .NE. 0) GO TO 305
      MULTIP = 3
      KKK = 6
      GC TO 306
      MULTIP = 4
      KKK = 8

C     IF MULTIP = 0, KK IS THE NUMBER OF SINGLE CLURVE GRAPHS. OTHERWISE
C     MULTIP IS THE NUMBER OF CURVES ON A SINGLE GRAPH.
C
      CC 306 IF (JJ .EQ. 0) GO TO 27

```



```

WRITE(6,214) (JTITLE(I),IP(I),I=1,JJ)
FCRMAT(///,56H THE COLUMN HEADINGS AND THE CRRRESPONDING VARIABLES AND THE CRRRESPONDING VARIABLES AR
IS ARE ;//,(15X,A8,9X,2HX(,I2,1H))
GC TO 28
WRITE(6,215)
FCRMAT(///,25H NC PRINTOUT IS REQUIRED )
IF(KK.EG.0) GC TO 308
IF(MULTIP.NE.0) GC TO 309
IF(KK.NE.1) GO TO 307
WRITE(6,216) KTITLE(1),KTITLE(1),KTITLE(2),IG(1),IG(2)
FCRMAT(///,52H THE GRAPH TITLE AND THE CRRRESPONDING VARIABLES AR
LE ;//,10X,2A8,4X,2FX(,I2,8H) VS. X(,I2,1H))
GC TO 31
WRITE(6,217) (KTITLE(I),KTITLE(I+1),IG(I),IG(I+1),I=1,KKK,2)
FCRMAT(///,64H THE INDIVIDUAL GRAPH TITLES AND THE CRRRESPONDING
VARIABLES ARE ;//,(10X,2A8,4X,2HX(,I2,8H) VS. X(,I2,1H))
GC TO 31
WRITE(6,1217)
FCRMAT(///,24H NO GRAPHS ARE REQUIRED )
GC TO 31
WRITE(6,1220)
FCRMAT(///,52H THE GRAPH TITLE AND THE CRRRESPONDING VARIABLES AR
LE ;//)
WRITE(6,1221) KTITLE(1),KTITLE(2),(IG(I),IG(I+1),I=1,KKK,2)
FCRMAT(10X,2A8,4X,2HX(,I2,8H) VS. X(,I2,1H),/, (30X,2FX(,I2,
8H) VS. X(,I2,1H))
C THIS ENDS THE BOOK-KEEPING. INITIALIZE BEFCRE ENTERING MAIN LCOP.
C
C 21 IPAGE = 0
C T = TI = 0
C NCPTS = 0
C NUMPTS = 0
C ITITLE(8) = IBLANK
C ITITLE(11) = IBLANK
C ITITLE(12) = IBLANK
C RUN(2) = BT(ARC)
C (11) = 20.
C (12) = 5.
C (13) = DT
C 42 I = 1, NN
C XC(1) = X(I)
C TC = T
C (10) = 2.
C RETURN
C 2000 IF(JJ.EG.0) GO TO 54
C

```



INT113800  
 INT113810  
 INT113820  
 INT113830  
 INT113840  
 INT113850  
 INT113860  
 INT113870  
 INT113880  
 INT113890  
 INT113900  
 INT113910  
 INT113920

INT113960  
 INT113970  
 INT113980  
 INT113990  
 INT114000  
 INT114010  
 INT114020  
 INT114030  
 INT114040  
 INT114050  
 INT114060  
 INT114070  
 INT114080  
 INT114090  
 INT114100  
 INT114110  
 INT114120  
 INT114130  
 INT114140  
 INT114150  
 INT114160  
 INT114170  
 INT114180  
 INT114190  
 INT114200  
 INT114210  
 INT114220  
 INT114230  
 INT114240  
 INT114250  
 INT114260  
 INT114270

```

INCPR = C(11)+0.0000001
C(11) = 20.
IF( MOD (NCPTS, 50*INCPR).EQ.0) GC TO 46
IF( MOD (NCPTS, 10*INCPR).EQ.0) GC TC 47
IF( MOD (NCPTS, 1) IPAGE + 1
IPAGE = IPAGE + 1
IF(NR.EG.1) GO TO 1047
WRITE(6,218) (JTITLE(I), I=1,6), IPAGE, JTITLE(7), (JTITLE(I), I=1,8)
WRITE(9,219)
GC TO 47
1047 WRITE(6,1218) (ITITLE(I), I=1,6), IPAGE, (JTITLE(I), I=1,8)
WRITE(6,219)
47 WRITE(6,219)
218 WRITE(6,219) (1H1,///,23X,6A8,7X,5HPAGE, I2,14H CF OUTPUT FOR,A8,////,
1218 FCRMAT(1H1,///,23X,6A8,27X,5HPAGE, I2,////,21X,4(A8,17X))
219 FCRMAT(1H )
48 DC 49 I=1,NN
49 TC(I) = X(I)
TC = T
C(10) = 3.
RETURN
C 50 CC 53 I=1, JJ
C
53 PR(I) = T
IF(IP(I).NE.0) PR(I)=XC(IP(I))
CCONTINUE
WRITE(6,220) (PR(I), I=1, JJ)
220 FCRMAT(7X, 8E25.5)
54 IF(KK.EG.0) GO TO 62
INCPR = C(12)+0.0000001
C(12) = 5
IF( MOD (NCPTS, INCPR).NE.0) GO TC 62
62 XC(I) = X(I)
TC = T
C(10) = 4.
RETURN
C 58 CC 61 I=1, KKK
C
61 GR(I) = T
IF(IG(I).NE.0) GR(I)=XC(IG(I))
CCONTINUE
IF(KKK.GE.8) GO TO 1610
KPI = KKK + 1
CC 1612 I=KPI, 8

```



```

1612 GR(I) = C./NUMPTS + 1
161C NUMPTS = NUMPTS + 1
      Y1(NUMPTS) = GR(1)
      Y2(NUMPTS) = GR(2)
      Y3(NUMPTS) = GR(3)
      Y4(NUMPTS) = GR(4)
      Y5(NUMPTS) = GR(5)
      Y6(NUMPTS) = GR(6)
      Y7(NUMPTS) = GR(7)
      Y8(NUMPTS) = GR(8)
      NUMPTS = NUMPTS + 1
62 IF(NUMPTS.LT.1000) GO TC 64
      WRITE(6,221)
221 FORMAT (//////,26H STOP AT 1000 GRAPH POINTS )
      GC TO 91
64 IF(NUMPTS.LT.10000) GO TC 66
      WRITE(6,222)
222 FORMAT (//////,32H STOP AT 10000 INTEGRATION STEPS )
      GC TO 91
66 IF(IPAGE - 15) 69,67,68
67 IF(MOD(NUMPTS,50)NE.0) GC TO 69
68 WRITE(6,223)
223 FORMAT(//////,27H STOP AT 10 PAGES OF OUTPUT )
      GC TO 91
69 DO I=1,NN
7C IF(ABS(X(I)).GT.1.0E+18) GO TO 71
      ICNTINUE
72 GC TO 72
71 WRITE(6,224)
224 FORMAT (//////,76H STOP WITH THE ABSOLUTE VALUE OF A DEPENDENT VARIABLE
      1 GR(1) GREATER THAN 1.0E+18. ,//,57H INTEGRATION PROBABLY UNSTABLE.
      2 TRY A SMALLER STEP SIZE. ,26HNO GRAPHS WILL BE PLOTTED.
      GC TO 330
72 DT = C(13)
73 IF(TI.GT.TF) GO TO 80
74 IF(T.LT.TF) GO TO 75
225 WRITE(6,225)
      FORMAT (//////,26H NORMAL STOP AT FINAL TIME )
75 GC TO 91
76 IF(T.GE.TF1) GO TO 77
      C(13) = DT
77 GC TO 87
78 IF(T.GE.TF2) GO TO 79
      C(13) = DT2
79 GC TO 87
      C(13) = DT3
8C IF(TF.GE.T) GO TC 74

```

```

INT14280
INT14290
INT14300
INT14310
INT14320
INT14330
INT14340
INT14350
INT14360
INT14370
INT14380
INT14400
INT14420
INT14440
INT14460
INT14480
INT14490
INT14510
INT14520
INT14540
INT14550
INT14560
INT14570
INT14580
INT14590
INT14600
INT14610
INT14620
INT14630
INT14640
INT14650
INT14660
INT14670
INT14680
INT14690
INT14700
INT14710
INT14720
INT14730
INT14740
INT14750

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INT114760
INT114770
INT114780
INT114790
INT114800
INT114810
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INT115130
INT115140
INT115150
INT115160
INT115170
INT115180
INT115190
INT115200
INT115210
INT115220
INT115230

```

```

87 C(10) = 5.
88 CALL RKUTTA (NN,T,X,DT,C,TC,XC,CX)
90 IF(C(10).EG.6.) RETURN
91 T = T + DT
92 GC TO 2000
93 IF(KK.EG.0) GO TO 330
94 IF(MULTIP.NE.0) GC TO 97
95 PRINT PLCT UP TO 4 INDIVIDUAL CURVES
96
97 NUMPTS=-NUMPTS
98 DC 310 II=1, KK
99 WRITE(6,99998)
100 FCORMAT(IH1)
101 ITITLE(9)=KTITLE(2*II-1)
102 ITITLE(10)=KTITLE(2*II)
103 GC TO (311,312,313,314), II
104 CALL PLCTP(X1,Y1,NUMPTS,0)
105 CALL TC 310
106 CALL PLCTP(X2,Y2,NUMPTS,0)
107 GC TO 310
108 CALL PLCTP(X3,Y3,NUMPTS,0)
109 GC TO 310
110 CALL PLCTP(X4,Y4,NUMPTS,0)
111 WRITE(6,99999) ITITLE
112 FCORMAT(IH0,8X,12A8)
113 GC TO 330
114
115 PLCT DUMMY CURVE ALCNG AXES TO SET SCALES FOR MULTIPLE PLOT
116
117 BIGX = 0.
118 BIGY = 0.
119 SMLX = 0.
120 SMLY = 0.
121 DC 1970 I=1, NUMPTS
122 XMAX=AMAX1 ( X1(I), X2(I), X3(I), X4(I) )
123 YMAX=AMAX1 ( Y1(I), Y2(I), Y3(I), Y4(I) )
124 XMIN=AMIN1 ( X1(I), X2(I), X3(I), X4(I) )
125 YMIN=AMIN1 ( Y1(I), Y2(I), Y3(I), Y4(I) )
126 IF(BIGX.LT.XMAX) BIGX=YMAX
127 IF(BIGY.LT.YMAX) BIGY=YMIN
128 IF(SMLX.GT.XMIN) SMLX=YMIN
129 IF(SMLY.GT.YMIN) SMLY=YMIN
130 CCNTINUE
131
132 197C

```







INT1157C0  
 INT115710  
 INT115720  
 INT115730  
 INT115740  
 INT115750  
 INT115760  
 INT115770  
 INT115780  
 INT115790  
 INT115800  
 INT115810  
 INT115820  
 INT115830  
 INT115840  
 INT115850

```

CT(4) = 1.0D0
II=0
II=II+1
TC = T + CT(II)*DT
XC(J) = X(J) + CT(II)*AK(II-1, J)
C(LC) = 6.C
RETURN
RC 4 J=1, NN
AK(II, J) = DT*DX(J)
IF(II.LT.4) GO TO 7
C 5 J=1, NN
X(J)=X(J)/(AK(1, J)+2.0*(AK(2, J)+AK(3, J))+AK(4, J))/6.C
C(LC) = 7.C
RETURN
END
  
```

```

//GC.SYSIN DD *
C
C
C THE FOLLOWING LISTING CONTAINS THE DATA CARDS AS DESCRIBED UNDER SUBROUTINE
C INTEGR USAGE. CARDS FOUR AND SIX ARE MODIFIED TO ACCEPT LARGER PARAMETER
C VARIATIONS. SEE STATEMENTS FOLLOWING INT12360 AND INT12550 FOR THE NEW FORMAT.
C
C
C ORGANIC DYE LASER ANALYSIS
  
```

```

1 3 C00001 .00001 .001 .001 1.0
1 C0000 C00.154C000C018000.0000.010000C00.034000C0C054200.0CC01C000.0C00
0 C0000 C000.C000C000C0000C0000C0000C00000000000000000000000000000000
TIME CO SINGLET 1 TRIPLET 2 GRUNC 3 GRUNC STATE 2 GRUNC STATE 3
SINGLET STATE 1 TRIPLET STATE 2 GRUNC STATE 3
  
```



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<p>The wavelength spectrum of a CW pumped Rhodamine 6-G organic dye laser is observed to shift approximately 150 Angstroms toward shorter wavelengths when pumped by 30 nsec pulses at a one megahertz repetition rate from a cavity-dumped Argon laser. Experimental evidence of the shift is presented, gain and rate equations are developed for a simplified dye laser model, and theoretical results are obtained for a computer simulation of the experiment. A comparison is made of the theoretical and experimental results and satisfactory</p>		



20. (cont'd)

agreement is obtained within the limits of the values of the parameters used and the assumptions made in formulating the model.



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