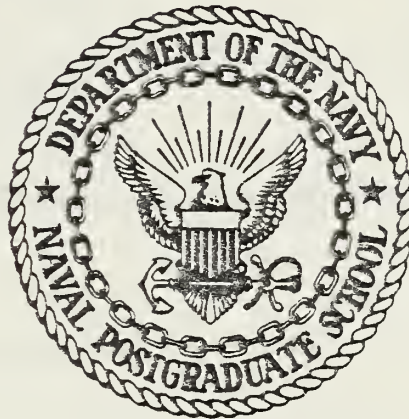


DUDLEY KNOX LIBRARY
NAVAL POSTGRADUATE SCHOOL
MONTEREY, CA 93940

NAVAL POSTGRADUATE SCHOOL

Monterey, California



THESIS

INVESTIGATION OF PIPE FLOW INSTABILITY
AND RESULTS FOR WAVE NUMBER ZERO

by

Michael James Arnold

December 1978

Thesis Advisor:

T. H. Gawain

Approved for public release; distribution unlimited.

T187434

REPORT DOCUMENTATION PAGE		READ INSTRUCTIONS BEFORE COMPLETING FORM
1. REPORT NUMBER	2. GOVT ACCESSION NO.	3. RECIPIENT'S CATALOG NUMBER
4. TITLE (and Subtitle) Investigation of Pipe Flow Instability and Results for Wave Number Zero		5. TYPE OF REPORT & PERIOD COVERED Master's Thesis; December 1978
7. AUTHOR(s) Michael James Arnold		6. PERFORMING ORG. REPORT NUMBER
9. PERFORMING ORGANIZATION NAME AND ADDRESS Naval Postgraduate School Monterey, California 93940		8. CONTRACT OR GRANT NUMBER(s)
11. CONTROLLING OFFICE NAME AND ADDRESS Naval Postgraduate School Monterey, California 93940		10. PROGRAM ELEMENT, PROJECT, TASK AREA & WORK UNIT NUMBERS
14. MONITORING AGENCY NAME & ADDRESS (if different from Controlling Office)		12. REPORT DATE December 1978
		13. NUMBER OF PAGES 122
		15. SECURITY CLASS. (of this report) Unclassified
		15a. DECLASSIFICATION/DOWNGRADING SCHEDULE
16. DISTRIBUTION STATEMENT (of this Report) Approved for public release; distribution unlimited.		
17. DISTRIBUTION STATEMENT (of the abstract entered in Block 20, if different from Report)		
18. SUPPLEMENTARY NOTES		
19. KEY WORDS (Continue on reverse side if necessary and identify by block number) Pipe Flow Instability		
20. ABSTRACT (Continue on reverse side if necessary and identify by block number) Past research by Harrison and Johnston on the stability of pipe flow yielded only tenuous results owing to errors in setup of the problem and in formulation of the complex axis boundary conditions. Recent advances in the formulation of these boundary conditions and application of generalized stability criteria allowed an accurate numerical solution to be made for angular wave number zero. The results show that flow for this case is characterized by certain		

ABSTRACT (Cont'd)

instabilities that have not been previously identified in linearized studies of this type.

A nonuniform computational mesh was developed which provided dramatic reductions in computational time on a limited basis.

Two data reduction programs were also developed to process and display data generated by the main program.

Investigation of Pipe Flow Instability
and Results for Wave Number Zero

by

Michael James Arnold
Lieutenant, United States Navy
B.S., University of Idaho, 1969

Submitted in partial fulfillment of the
requirements for the degree of

MASTER OF SCIENCE IN AERONAUTICAL ENGINEERING

from the

NAVAL POSTGRADUATE SCHOOL
December 1978

ABSTRACT

Past research by Harrison and Johnston on the stability of pipe flow yielded only tenuous results owing to errors in setup of the problem and in formulation of the complex axis boundary conditions.

Recent advances in the formulation of these boundary conditions and application of generalized stability criteria allowed an accurate numerical solution to be made for angular wave number zero. The results show that flow for this case is characterized by certain instabilities that have not been previously identified in linearized studies of this type.

A nonuniform computational mesh was developed which provided dramatic reductions in computational time on a limited basis.

Two data reduction programs were also developed to process and display data generated by the main program.

TABLE OF CONTENTS

I.	INTRODUCTION -----	9
II.	THE VORTICITY TRANSPORT EQUATION -----	12
III.	NUMERICAL METHODS -----	17
IV.	RESULTS -----	25
	A. STABILITY -----	25
	B. PERTURBATION VELOCITY PLOTS -----	27
	C. STABILITY CONTOUR PLOTS -----	28
	D. NONUNIFORM MESH EFFECTS -----	29
	E. NUMERICAL ACCURACY -----	31
V.	CONCLUSIONS AND RECOMMENDATIONS -----	46
APPENDIX A:	DERIVATION OF VORTICITY TRANSPORT EQUATION COEFFICIENTS -----	48
APPENDIX B:	FINITE DIFFERENCE EQUATIONS -----	51
APPENDIX C:	NON-UNIFORM MESH -----	57
APPENDIX D:	DERIVATION OF PERTURBATION VELOCITIES -	65
COMPUTER PROGRAMS	-----	69
LIST OF REFERENCES	-----	121
INITIAL DISTRIBUTION LIST	-----	122

LIST OF FIGURES

3-1	Finite Difference Mesh -----	18
3-2	Basic Composition of Coefficient Arrays and Vector of Unknowns -----	20
4-1	Normalized Perturbation Velocity -----	34
4-2	Normalized Perturbation Velocity -----	35
4-3	Normalized Perturbation Velocity -----	36
4-4	Normalized Perturbation Velocity -----	37
4-5	Normalized Perturbation Velocity -----	38
4-6	Stability Contour Plot -----	39
4-7	Stability Contour Plot -----	40
4-8	γ^* Versus Number of Mesh Points, N -----	41
4-9	γ^* Versus Number of Mesh Points, N -----	42
4-10	γ^* Versus Mesh Parameter, Lambda -----	43
4-11	Normalized Perturbation Velocity -----	44
4-12	Normalized Perturbation Velocity -----	45
C-1	R versus η for Four Selected Values of Lambda-Axis Offset -----	63
C-2	R versus η for Four Selected Values of Lambda-Wall Offset -----	64

TABLE OF SYMBOLS

C	Constant in non-uniform mesh functions given by equations (C-32) and (C-40)
D, D^2, \dots	Partial derivatives with respect to r .
D^*, D^{*2}, \dots	Partial derivatives with respect to η .
e	Base of natural logarithms.
$\bar{e}_x, \bar{e}_r, \bar{e}_\theta$	Unit vectors along the x , r and θ axes in cylindrical coordinates.
F, G, H	Components of the velocity vector potential defined in equation (2-6).
f_{11}, f_{22}, \dots	Coefficients of $D^*Q, D^{*2}Q, \dots$ in equations (C-9) through (C-12) as defined in equations (C-13) through (C-22).
i	$+\sqrt{-1}$, the imaginary unit. Also used as an index in Section III and Appendix D.
N	The number of interior points in the finite difference mesh of Section III.
O	Symbol denoting the phrase "of order".
Q	The component of the velocity vector potential derived from the component H by the change of variable, $H = rQ$.
R_e	Reynolds number based on mean velocity and pipe radius.
t	Time.
U	The streamwise velocity in Pipe Poiseuille Flow as defined by equation (2-11).
u, v, w	Components of the complex perturbation velocity defined in equation (D-1).
\bar{W}	Complex vector potential of perturbation velocity defined in equation (D-2).
x, r, θ	Cylindrical coordinates.
α	$\alpha_R + i\alpha_I$. Complex wave number of the perturbation in the x -direction.

β	in. Complex wave number of the perturbation in the θ direction, where $n = 0, 1, 2, 3, \dots$
δ	$1/(N+1)$. The r or η increment in the finite difference approximations of the derivatives of Q .
η	The independent variable replacing r in the nonuniform mesh of Appendix C.
γ	$\gamma_R + i\gamma_I$. Complex frequency of the perturbation.
$\bar{\Gamma}$	The vorticity transport equation expressed in abbreviated notation as defined in equation (2-7).
$\Gamma_x, \Gamma_r, \Gamma_\theta$	The components of $\bar{\Gamma}$ in cylindrical coordinates as defined in equation (2-7).
λ	Mesh offset parameter as defined in equations (C-32) and (C-40).
∇	Linear vector operator (nabla)
\times	Vector cross-product operator.
[]	Brackets enclosing a matrix.
{ }	Brackets enclosing a column vector.

I. INTRODUCTION

The problem of finding an analytical solution to the pipe flow stability problem has been pursued actively ever since the classical experiments of Osborne Reynolds [10] about 100 years ago. Up to now, however, no investigation has been able to satisfactorily predict flow instabilities, although many approaches have been taken.

Salwen and Grosch [11] studied pipe flow with various angular wave numbers and sinusoidal streamwise perturbations and concluded that it was stable for all axial and angular wave numbers. Perturbations with exponential growth in space but a purely sinusoidal time variation were researched by Garg and Rouleau [2] and those with both exponential growth in space and in time by Gill [3]. Both concluded that the flows were stable.

Because of this inability of linear theory to account for experimental fact, explanations by Davey and Drazin [1] involving finite disturbances and by Huang and Chen [5] and Leite [7] involving conditions at the pipe entrance have been offered. While these investigations have indeed shown instabilities to exist, a completely general solution to the linear problem has never been achieved.

Recently a more general theory was presented by Harrison [4] and further investigated by Johnston [6]. These two studies, however, failed to produce conclusive results due

to mathematical errors in the problem setup and inadequate formulation of the boundary conditions at the axis. Gawain [9] has subsequently formulated the axis boundary conditions in a new way which corrects the previous discrepancies and promises further advances.

For angular wave number, n , equal to zero, radical simplifications result in the governing equations (Section II), indicating that this case should be approached first. This investigation centers on that case.

Preliminary checks using the computer program of Ref. 6 revealed that, of the two eigenfunctions, G and H , which occur in this problem and which are uncoupled for $n = 0$, the latter appeared to be the more critical. Hence the present research was arbitrarily restricted to investigation of the stability of eigenfunction H . A similar study of the other eigenfunction, G , for $n = 0$ remains to be completed at some future time. Comparable calculations for other wave numbers ($n = 1, 2, 3, \dots$) also remain to be accomplished in the future. Extensive and systematic calculations of this type will be essential to provide the factual basis for a comprehensive theory of pipe flow stability.

Reverting to the case at hand, eigenfunction H for wave number $n = 0$, we note that the program of Ref. 6 was rewritten for this case, incorporating the newly formulated boundary conditions of Ref. 9. In addition, a new, generalized stability criteria was adopted. Moreover, a new technique was introduced which allows the use of nonuniform meshes to reduce computational time.

Lastly, two data reduction programs were written to process data produced by the main investigative program.

II. THE VORTICITY TRANSPORT EQUATION

Although a complete treatment of this subject is contained in Appendix A of Ref. 4 and further addressed in Ref. 6 and Ref. 9, it is felt that a brief overview is still required here to maintain continuity with previously referenced works. This discussion is an abbreviated version of Section II of Ref. 6.

Laminar flow of an incompressible fluid of constant viscosity is governed by the Navier-Stokes equation and the continuity equation. Taking the curl ($\nabla \times$) of the Navier-Stokes equation and introducing a perturbation velocity (\bar{v}) and vorticity ($\bar{\omega}$) gives the vorticity transport equation which is equation (A-10) of Appendix A, Ref. 4.

Expressing this equation in terms of the complex velocity vector potential, \bar{W} , gives

$$W(x, r, \theta, t) = (\bar{e}_x F(r) + \bar{e}_r G(r) + \bar{e}_\theta H(r)) e^X \quad (2-1)$$

where

$$X = \alpha x + \beta \theta + \gamma t \quad (2-2)$$

and

$$\bar{v} = \nabla \times \bar{W} \quad (2-3)$$

$$\bar{\omega} = \nabla \times \bar{v} . \quad (2-4)$$

It should also be noted that, as shown in part one of Appendix G in Ref. 4, α and γ are complex while β is a purely imaginary quantity defined by

$$\beta = i n \quad n = 0, 1, 2, \dots \quad (2-5)$$

When expressed in the form of equation (2-1), the vorticity transport equation becomes three simultaneous fourth-order differential equations of the form

$$\begin{aligned}
 & [M_4] \begin{Bmatrix} D^4 F \\ D^4 G \\ D^4 H \end{Bmatrix} + [M_3] \begin{Bmatrix} D^3 F \\ D^3 G \\ D^3 H \end{Bmatrix} + [M_2] \begin{Bmatrix} D^2 F \\ D^2 G \\ D^2 H \end{Bmatrix} \\
 & + [M_1] \begin{Bmatrix} DF \\ DG \\ DH \end{Bmatrix} + [M_0] \begin{Bmatrix} F \\ G \\ H \end{Bmatrix} - \gamma ([N_2] \begin{Bmatrix} D^2 F \\ D^2 G \\ D^2 H \end{Bmatrix} \\
 & + [N_1] \begin{Bmatrix} DF \\ DG \\ DH \end{Bmatrix} + [N_0] \begin{Bmatrix} F \\ G \\ H \end{Bmatrix}) = \begin{Bmatrix} 0 \\ 0 \\ 0 \end{Bmatrix} \quad (2-6)
 \end{aligned}$$

Equations (2-5) may be further expressed in the abbreviated form

$$\bar{\Gamma} = \begin{Bmatrix} \bar{\Gamma}_x \\ \bar{\Gamma}_r \\ \bar{\Gamma}_\theta \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \\ 0 \end{Bmatrix} \quad (2-7)$$

where $\bar{\Gamma}$ appears to be a set of three coupled equations in the components of \bar{W} . As given in Appendix B of Ref. 4, equations (2-7) actually represent only two independent conditions and by an appropriate linear combination of Γ_x and Γ_θ , equations (2-6) can be expressed as a set of two equations in three unknowns. The appropriate linear combination is given in Appendix B of Ref. 4 and yields the set of equations

$$\begin{aligned} \Gamma_r &= 0 \\ -\frac{in}{r} \Gamma_x + \alpha \Gamma_\theta &= 0 . \end{aligned} \quad (2-8)$$

Except for the case where n is equal to zero, equations (2-8) do not uncouple. The linear combination given by the second of equations (2-8) does, however, reduce the highest order derivative of $G(r)$ in equations (2-6) to second order. Appendix C of Ref. 4 illustrates the redundancy of the three components of \bar{W} , allowing one of these components to be arbitrarily set to zero for all r . The maximum benefits of equations (2-8) are obtained if

$$F(r) = 0 \quad (2-9)$$

Incorporating equations (2-8) and (2-9) into equations (2-6) results in the form

$$\begin{aligned}
 & [M'_4] \begin{Bmatrix} D^4 G \\ D^4 H \end{Bmatrix} + [M'_3] \begin{Bmatrix} D^3 G \\ D^3 H \end{Bmatrix} + [M'_2] \begin{Bmatrix} D^2 G \\ D^2 H \end{Bmatrix} \\
 + & [M'_1] \begin{Bmatrix} DG \\ DH \end{Bmatrix} + [M'_0] \begin{Bmatrix} G \\ H \end{Bmatrix} - \gamma ([N'_2] \begin{Bmatrix} D^2 G \\ D^2 H \end{Bmatrix} \\
 + & [N'_1] \begin{Bmatrix} DG \\ DH \end{Bmatrix} + [N'_0] \begin{Bmatrix} G \\ H \end{Bmatrix}) = \begin{Bmatrix} 0 \\ 0 \end{Bmatrix} \quad (2-10)
 \end{aligned}$$

where the coefficient matrices are given by equations (2-10) through (2-17) of Ref. 6. It is appropriate to note that these same coefficient matrices appear in Ref. 9, equations (A1) through (A9), in a slightly different form resulting from the substitutions

$$U = 2(1 - r^2) \quad (2-11)$$

$$t = \alpha^2 + \frac{\beta^2}{r^2} \quad \text{and} \quad (2-12)$$

$$T = \alpha U - \frac{1}{R_e} \left(\alpha^2 + \frac{\beta^2}{r^2} \right) . \quad (2-13)$$

As discussed in the previous section, the case where

$$\beta = \text{in} , \quad n = 0 \quad (2-14)$$

leads to great simplifications in equations (2-10), (2-12) and (2-13). In particular, equations (2-10) uncouple and allow an independent investigation of either H or G. As a result of the findings discussed in Section I, it was decided to explore the function H only. This reduced equation (2-10) to that of equation (A-6) of Appendix A, which is a linear, homogeneous fourth order differential equation in $H(r)$.

III. NUMERICAL METHODS

Substituting the change of variable $H = rQ$ as given in equation (A-1) and the coefficients defined in equations (A-11) through (A-18) into the vorticity transport relation, equation (A-6), gives the expression

$$\begin{aligned} M_4 D^4 Q + M_3 D^3 Q + M_2 D^2 Q + M_1 DQ + M_0 Q \\ - \gamma [N_2 D^2 Q + N_1 DQ + N_0 Q] = 0 , \end{aligned} \quad (3-1)$$

which is a homogeneous fourth order differential equation in $Q(r)$. The boundary conditions for this case are derived in detail in Ref. 9 as

$$\begin{aligned} Q(1) &= 0 \\ DQ(1) &= 0 \\ DQ(0) &= 0 \\ D^3 Q(0) &= 0 . \end{aligned} \quad (3-2)$$

The boundary finite difference equations derived in Appendix B from equations (3-2), along with the standard central difference equations given in Ref. 6, allow the function $Q(r)$ to be approximated by a finite number of discrete unknowns. As shown by Figure 3-1 below, the non-dimensionalized radius of the pipe is divided into a one-dimensional

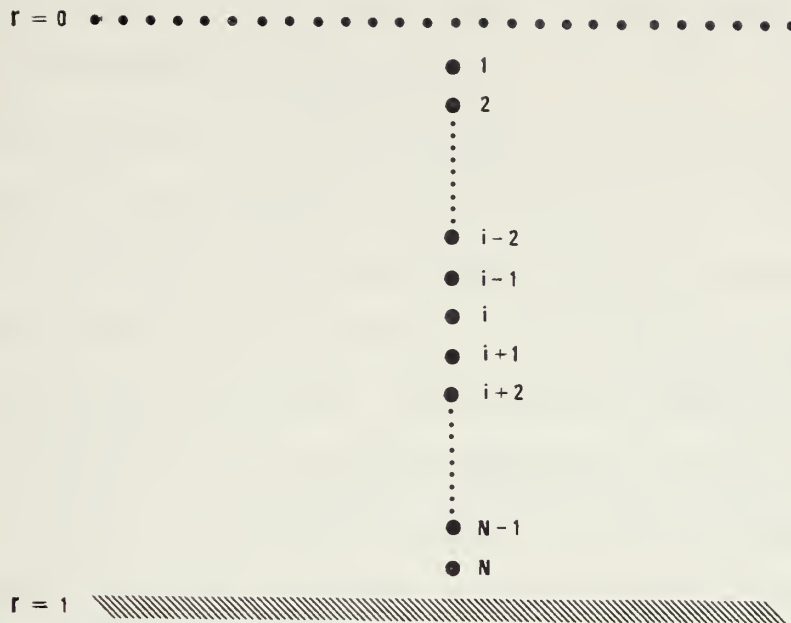


Figure 3-1 Finite Difference Mesh

computational mesh consisting of N interior points, $N+1$ intervals, and $N+2$ total points, including the boundary points at $r = 1$ and $r = 0$. As will be discussed later, the spacing between these points may or may not be uniform. For the uniform case, the spacing is defined by

$$\delta = 1/(N+1) . \quad (3-3)$$

For the nonuniform case, a change of independent variable is performed. The spacing of the new independent variable, η , is still given by equation (3-3).

With a nonuniform mesh, the points shown in Figure 3-1 will be concentrated near the axis or near the wall according

to the type of offset specified. These effects are discussed in detail in Section IV.

Substitution of the finite difference equations of Appendix B into equation (3-1) results in a set of N , linear, algebraic difference equations in terms of the unknown value of Q at each of the N interior points of the computational mesh. Since each of these equations is of the form of a linear combination of the i th, central, point and the two, three or four adjacent points (depending on the order of the derivative being approximated), this system of equations consists of a coefficient array multiplying a vector containing the unknown value of the function Q at each of the N interior points. This technique allows the problem to be converted into an eigenvalue problem of the form

$$[X] \{Q\} - \gamma [Y] \{Q\} = 0 \quad (3-4)$$

with the basic composition of the arrays $[X]$ and $[Y]$ and the vector $\{Q\}$ as illustrated in Figure 3-2 below.

It should be noted at this point that Figure 3-2 differs somewhat from the normal finite difference banded matrix in the first two rows and last row because of the method of deriving the finite difference approximations at the boundaries. Additionally, the order of the N unknowns has been reversed from that of Ref. 6. This was done to conform to standard matrix notation.

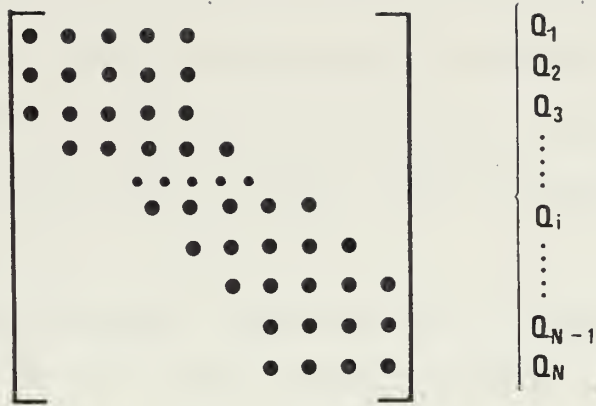


Figure 3-2 Basic Composition of Coefficient Arrays and Vector of Unknowns

This array is established by the subroutine MSET2 in conjunction with the subroutine MSET1 and function subprograms CQM1E1 and CQM2E1, which compute the numerical value for each element in the array. Subroutine MSET1 provides the coefficients given by equations (A-11) through (A-18) of Appendix A or by equations (C-24) through (C-31) if a nonuniform mesh is specified. Function CQM1E1 then computes the values for each of the elements of array [X] in equation (3-4) using the coefficients passed from subroutine MSET1 in vector CQM1. Function subprogram CQM2E1 performs the same function for matrix [Y] in equation (3-4) using the coefficients passed in vector CQM2.

The solution of the eigenvalue problem as formulated to this point is carried out by the controlling subroutine of program PIPE0, subroutine STAB, by the following steps:

- 1) Subroutine MSET2 is called twice to set up the coefficient matrices [X] and [Y] of equation (3-4).

- 2) Subroutine CDMTIN is then called to invert matrix $[Y]$, the second coefficient array in equation (3-4). CDMTIN was obtained from the IBM Library routine CMTRIN by modifying it to accept double precision arrays.
- 3) Both coefficient arrays, $[X]$ and $[Y]$, are then pre-multiplied by $[Y]^{-1}$. Since multiplication of an array by its inverse invariably results in the identity matrix, $[I]$, only the product $[Y]^{-1}[X]$ is computed using subroutine MULM. This converts the eigenvalue problem of equation (3-4) to the more conventional form

$$([Z] - \gamma[I])\{Q\} = 0 \quad (3-5)$$

where

$$[Z] = [Y]^{-1}[X] \quad (3-6)$$

- 4) Since all programs currently available for solving equations (3-5) require that the real and imaginary parts of the elements of $[Z]$ be presented in separate arrays, subroutine DSPLIT is called to accomplish this.
- 5) The eigenvalues and eigenvectors of equations (3-5) are computed using subroutines EBALAC, EHESSC, ELRH2C and EBBCKC which are available through the International

Math and Statistics Library. Subroutine EBALAC balances matrix [Z] by equalizing the exponents of all terms. The details of this transformation are retained for later use. The balanced matrix is then passed to subroutine EHSSC where it is reduced into the complex upper Hessenberg form. Subroutine ELRH2C then solves for the eigenvalues and eigenvectors. To transform the eigenvectors back into the original unbalanced form, EBBCKC is finally called using information passed from subroutine EBALAC.

For each solution, subroutine STAB determines the least stable eigenvalue (largest algebraic value) and then writes the values of N , R_e , α_R , α_I , λ , γ_{RL} , γ_{IL} and KSET to file FT02F001. The eigenvector corresponding to the least stable eigenvalue is also written to FILE FT02F001 when MODENO is set equal to one.

Control of subroutine STAB is accomplished by the main program, PIPE0. This program is a time-sharing (CP/CMS) program. Modes one and three compute the stability of the flow for a given set of input conditions. Mode one writes the least stable eigenvector to FILE FT02F001 while this output is inhibited when MODENO is set equal to three. To generate data for program EIGFCN, program PIPE0 must be run with MODENO equal to one.

Mode two operation generates a grid of stability values (stability map) based on parameters read in from FILE FT01F001. Due to the long run time in this mode, only small meshes can be generated under CP/CMS. Longer runs must be accomplished under batch, with changes to the program as specified in the comments section. Data is output to file FT03F001 when MODENO is equal to two and is compatible with program STBCONT.

The plotting programs EIGFCN and STBCONT were used to process the data generated by program PIPE0 in modes one and three, respectively. Program EIGFCN generates normalized plots of the perturbation velocity, u , as a function of radius, r . The perturbation velocities generated in accordance with Appendix D were normalized in two steps. First the perturbation velocity of largest magnitude was determined. Letting this velocity be termed u_C , a normalizing constant producing unit magnitude and zero phase angle in u_C was found in the following manner:

If

$$u_C = u_{RC} + iu_{iC} , \quad (3-7)$$

then

$$Cu_C = 1 + i(0) \quad (3-8)$$

where C is the normalizing constant. Thus,

$$C = \frac{1}{u_{RC} + iu_{iC}} = \frac{u_{RC} - u_{iC}}{(u_{RC}^2 - u_{iC}^2)} \quad (3-9)$$

$$= \frac{\bar{u}_C}{|u_C|^2} \quad (3-10)$$

where \bar{u}_C is the complex conjugate of u_C .

The nondimensionalized radius values were taken directly from the data cards for uniform meshes or computed from equations (C-32) or (C-40) in the case of a nonuniform mesh.

Program STBCONT plots the stability contours against α_R and α_I . The stability map generated by program PIPE0 is searched columnwise and rowwise for sign changes for each of the three stability criteria discussed in Section V and Ref. 9. The points are then plotted, producing contours of incipient, critical and fully developed instability and areas that denote stable flow and subcritical, supercritical and hypercritical instability.

Both programs, EIGFCN and STBCONT, utilize the NPS VERSATEC plotter, certain built-in VERSATEC subroutines, and subroutine PLOTG. These routines are only accessible when running under FORTCLGW.

IV. RESULTS

A. STABILITY

Since an understanding of the term stability is necessary to interpret the results of this investigation, a brief discussion is presented here. A complete discussion of the generalized criteria of stability is given by Gawain [9].

The characteristics of the flow for the case $n = 0$ are set by the parameters R_e and α . For fixed values of these parameters, the solution of equations (3-5) is a set of N eigenvalues, γ , and their corresponding eigenvectors, Q . As can readily be seen from equation (2-1), the value of the real part of the complex eigenvalue γ will determine the growth or decay rate in time of the perturbation. Since positive values of the real part of γ represent an exponential growth rate in time, the most important γ is the one having the largest algebraic value for its real part. This root is termed the least stable root and will be represented by the symbol γ_{RL} . As the stability represented by γ_{RL} is that seen by a fixed observer, it is not the most general criterion. As derived in Ref. 9, a more appropriate stability criterion is that based on an axis system moving at the average volumetric velocity of the flow. This criteria is termed γ_{RL}^* and is defined by Ref. 9 as

$$\gamma_{RL}^* = \gamma_{RL} + \alpha_R . \quad (4-1)$$

For this and subsequent discussions, the subscript will be dropped and γ^* will refer to the quantity defined by equation (4-1). Three stability cases arise from this equation. The first is termed incipient instability and is defined by

$$\gamma^* = -|\alpha_R| . \quad (4-2)$$

The second case, termed critical instability, is given by

$$\gamma^* = 0 \quad (4-3)$$

and, lastly, the case termed fully developed instability is said to exist when

$$\gamma^* = +|\alpha_R| . \quad (4-4)$$

The transition from stable flow to fully developed instability is progressive and several distinct stages are given in Ref. 9 to describe this transition. The region from incipient to critical instability is termed subcritical instability, that from critical instability to fully developed instability is called supercritical instability while that beyond fully developed instability is termed hypercritical instability.

B. PERTURBATION VELOCITY PLOTS

Initial investigation of the function Q was centered around plotting its appearance in the region of interest. A Reynolds number of 1150 (2300 based on diameter) was chosen as this value is generally accepted as the nominal value for transition to turbulent flow. The value of α was set at $-0.5 + i 10.0$ for the major part of the investigation as preliminary checks revealed that supercritical instabilities were present for this value. A secondary Reynolds number of 4000 was chosen to show trends.

The quantity chosen as the most realistic and representative of the eigenfunction Q is the axial perturbation velocity, u . This quantity was derived from the elements of the least stable eigenvector as outlined in Appendix D. Initially, R_e and α_I were held fixed and α_R was varied over a range of positive and negative values. For values of α_R below about two, the normalized perturbation velocity was found to have all activity near the axis with a decay essentially to zero by $r = 0.3$. A typical plot of u versus r for an α_R in this range is shown in Figure 4-1. When α_R was made sufficiently positive, the plot changed significantly in both appearance and region of activity. Figure 4-2 shows a plot of u for $\alpha_R = 2.5$. The activity can now be seen to be concentrated near the wall, with most of the activity occurring at r values greater than 0.7.

Although no particular relationship between the nature of u and the stability of the flow was evident or expected,

the plots were nevertheless valuable as indicators for various parameters involved in the investigation.

First, as can be seen by the differences in Figures 4-1 and 4-2, the plots were ideal indicators of changes in the nature of the function Q . Secondly, the adequacy of the mesh could be directly observed by noting the number of points defining the curves in regions of high activity. Figures 4-3, 4-4 and 4-5 show the same conditions as Figure 4-1 but with decreasing number of mesh points, N . Lastly, the effects of nonuniform meshes could be observed as will be discussed later in this section.

C. STABILITY CONTOUR PLOTS

The principal results of this investigation are shown in Figures 4-6 and 4-7. Although these two figures pertain to only a limited portion of the complex α plane, they do represent a significant advance in the investigation of pipe flow stability. As can be seen in these figures, the flow is characterized by regions of differing stability, ranging from stable through supercritical instability. Note that these two figures correspond to Reynolds numbers of 1150 and 4000, respectively. This is a result that has not, to this writer's knowledge, been heretofore achieved by a linearized analysis of fully developed pipe flow. The figures also show that, as has been born out by previous investigations, flow for purely sinusoidal oscillations ($\alpha_R = 0$) is stable. Additionally, a comparison of Figures 4-6

and 4-7 shows the effect of Reynolds number on the flow stability. It is clear from this comparison that an increase in Reynolds number reduces the size of the stable regions in the complex α plane; in other words, stability decreases with increasing Reynolds number. This trend agrees with our general experience pertaining to fluid flow. Lastly, the effect of the real and imaginary parts of the wave number α can readily be seen. For α_R , increasingly negative values produce successively greater levels of instability. While a contour plot was not produced for positive values of α_R , point checks of stability in this region suggest that somewhat similar contours exist in the right half-plane also. For α_I , increasing values produce increasing stability. This effect is also more pronounced at the lower Reynolds number.

D. NONUNIFORM MESH EFFECTS

One of the difficulties in this investigation was the relatively long computing time required to obtain an accurate solution, especially when operating under CP/CMS (time-sharing). The major factor controlling computing time was the number of interior mesh points, N . As an example, an increase in N of 50 percent resulted in a fourfold increase in computing time. Therefore, the desired objectives of rapidity and accuracy were in direct conflict. Additionally, follow-on investigations for values of angular wave number n other than zero involve matrices twice the order required for this case because of the coupling of equations (2-8).

For these reasons, a nonuniform mesh was developed to obtain increased accuracy at lower values of N . The nature of the velocities as seen in Figures 4-1 and 4-2 shows that a high degree of resolution in the computational mesh is only required in the vicinity of the axis (α_R less than about 2) or the wall (α_R greater than about 2). It was therefore theoretically possible to redistribute the points at moderate values of N to attain resolutions equivalent to much finer (and more time-consuming) uniform meshes.

As can be seen from Figures 4-8 and 4-9, the value of γ^* varies with the number of mesh points, N . Theoretically, each of these curves would approach some limiting value if N were increased without bound, and it is this theoretical limit that represents the required solution. In practice, it is adequate to approximate the unknown limit by a point that lies on the relatively flat portion of the curve at a value of N which is practically attainable and which does not involve a prohibitively long computing time. It has been found in this investigation that $N = 79$ fulfills these conditions.

The conversion to a nonuniform mesh involved a change of independent variable and the introduction of an analytical function to control the distribution of the mesh points. The details of these steps are given in Appendix C. By varying the mesh offset parameter, λ , it was possible to vary γ^* over a wide range. To determine when the high

accuracy solution ($N = 79$) and the nonuniform solutions were approximately equal, γ^* was plotted versus λ for fixed values of R_e , α and N with the value of γ^* for $N = 79$ as a reference. Figure 4-10 shows a plot of this type for $N = 31$. The appropriate value of λ can be seen to be approximately 1.1. Figure 4-11 is the perturbation velocity plot of the solution for $N = 31$ and $\lambda = 1.1$ for the same R_e and α as Figure 4-1. Note that the γ^* values are equal for these two figures. While the resolution of Figure 4-11 is not quite as fine as that of Figure 4-1, a comparison of Figure 4-11 with Figure 4-5 makes the improved resolution obvious. Figures 4-2 and 4-12 are similar to Figures 4-1 and 4-11 except that a wall offset was used. Note that for this case $\lambda = 1.2$, which points to a drawback of the nonuniform mesh, that of dependence on input conditions. While a check of λ dependence on α was not made, it most probably exists. There is also, however, the possibility that for small regions of the complex α plane, the variations in λ are small enough to allow an average value of λ to be nearly optimum for the entire region. While not used for the main results of this study, the method as developed here may well prove to be of maximum utility in follow-on investigations of higher angular wave numbers.

E. NUMERICAL ACCURACY

To ensure that the solutions presented here were of sufficient accuracy, two separate checks were made. The

first, γ^* dependence on N , is the most commonly used criterion.

For a solution to be accurate, it should be virtually independent of mesh fineness, that is, of N . The required magnitude of N for an accurate solution was found by plotting γ^* against N . Figures 4-8 and 4-9 both show that the solution is well converged for $N = 79$ at Reynolds numbers of 1150 and 4000, as γ^* changes by only .001 to .003 from $N = 31$ to $N = 79$ for both values of Reynolds number.

The second verification of the solution, so obvious that it is sometimes overlooked, involves simply substituting the numerical solution (least stable eigenvector) into the governing equation to ensure that it is indeed being satisfied. A short program was independently written to check the finite difference representations of equation (3-1) at the first and last interior stations and at a mid-radius station. Initial checks of numerical solutions yielded unsatisfactory results and led to the discovery of various programming errors. In particular, it was discovered that four double precision constants in the finite difference approximations were lacking the required "D0" exponent. Elimination of these seemingly trivial errors resulted in a surprising four order-of-magnitude improvement in the accuracy of the solution, with the left side of equation (3-1) improving from order 10^{-4} to order 10^{-8} .

It is instructive to note at this point that the order of magnitude of the left side of equation (3-1) is not the

true measure of its satisfaction. A more correct procedure is to compare this value with the largest term in the equation. When examined from this viewpoint, the relative error for solutions at $R_e = 1150$ and $R_e = 4000$ are found to be of order 10^{-11} to 10^{-12} , a very satisfactory result.

Therefore, by these results, the solutions presented here are both virtually independent of N and satisfy the governing differential equation to a high degree. The efforts expended to reach these conclusions were well worth the result and also point out that attention to detail is fundamental to accurate numerical results.

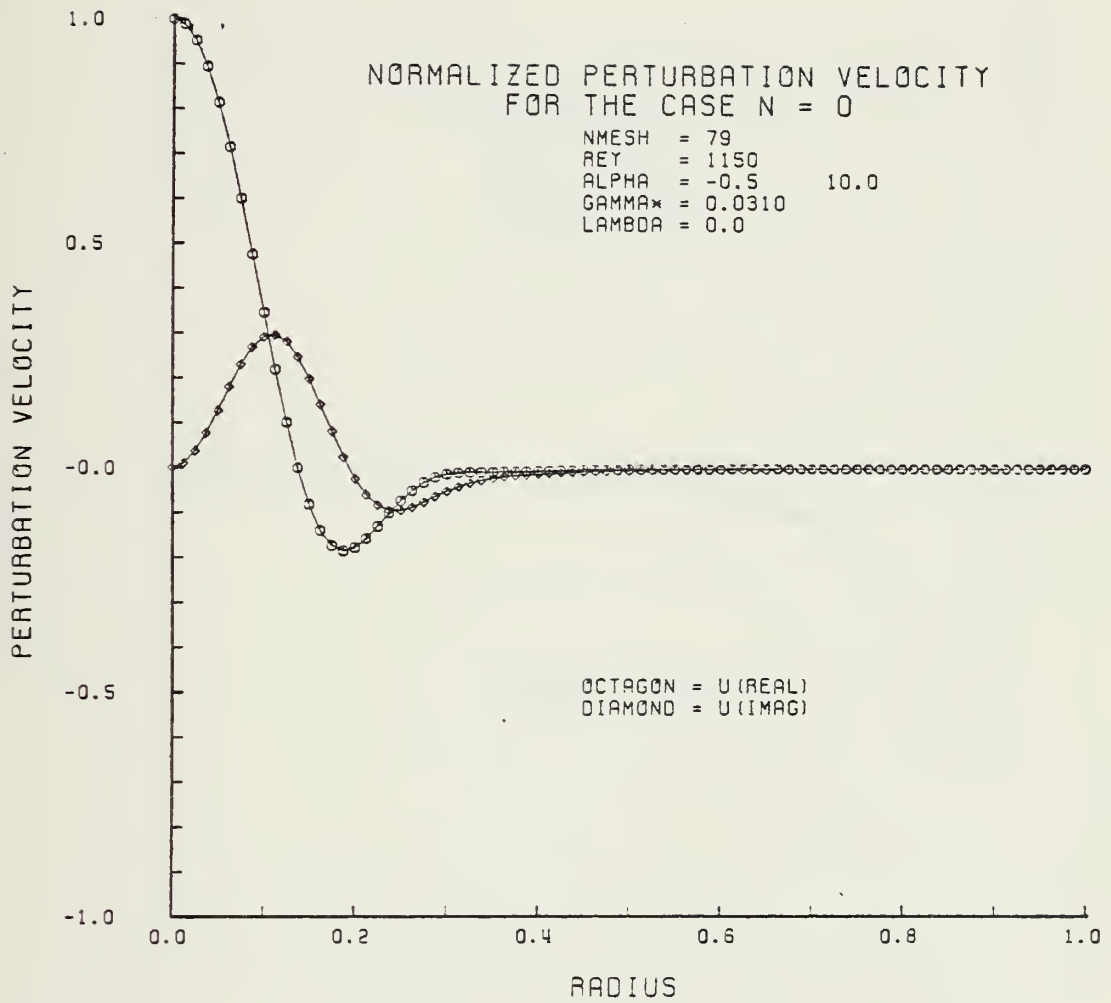


FIGURE 4-1

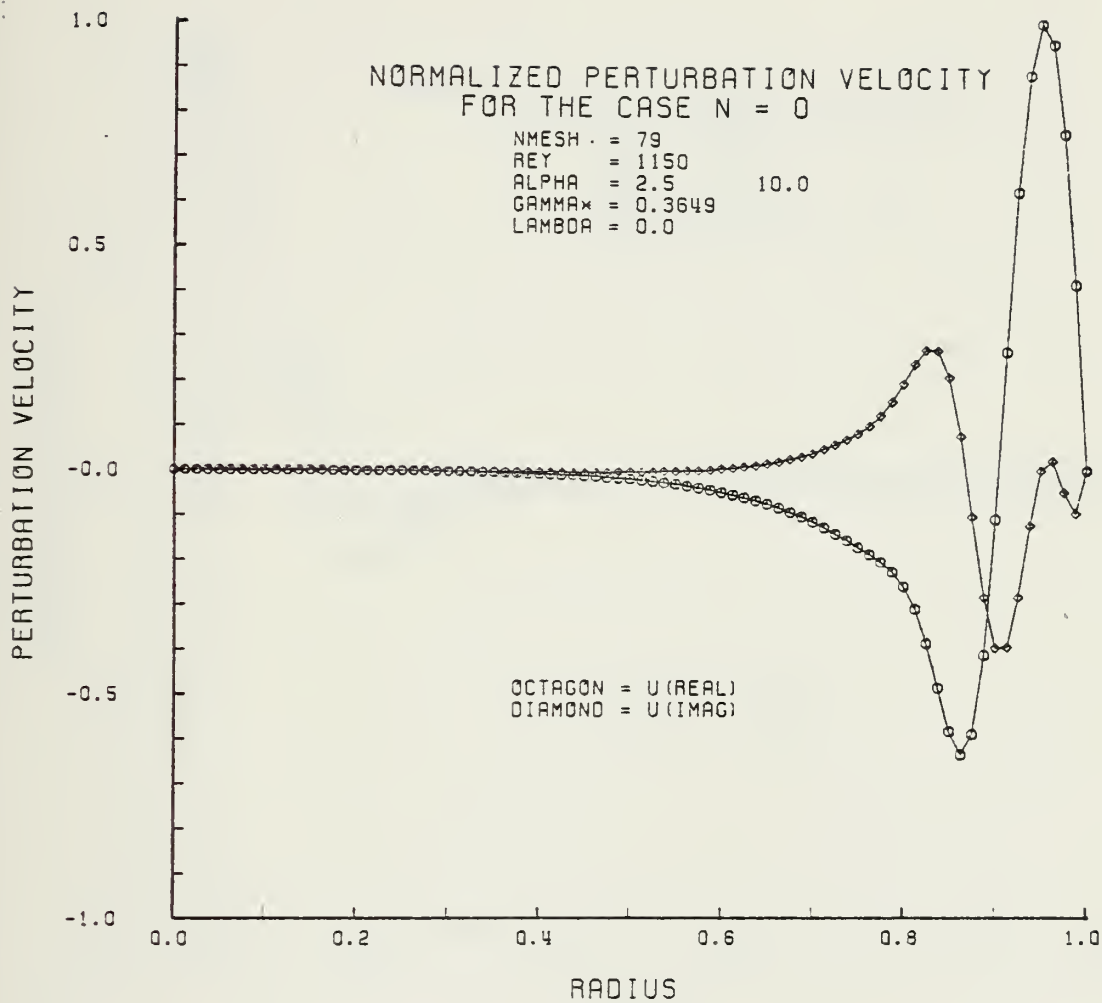


FIGURE 4-2

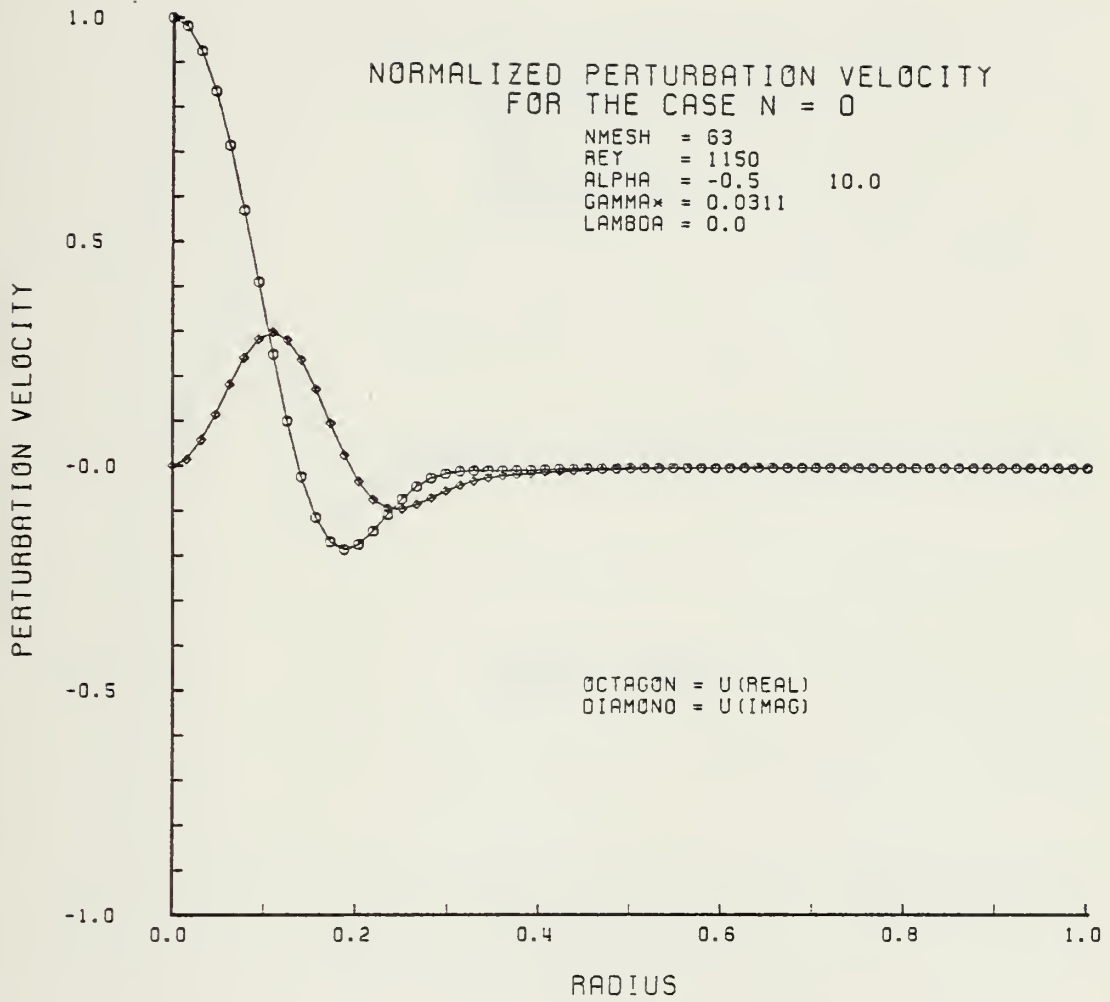


FIGURE 4-3

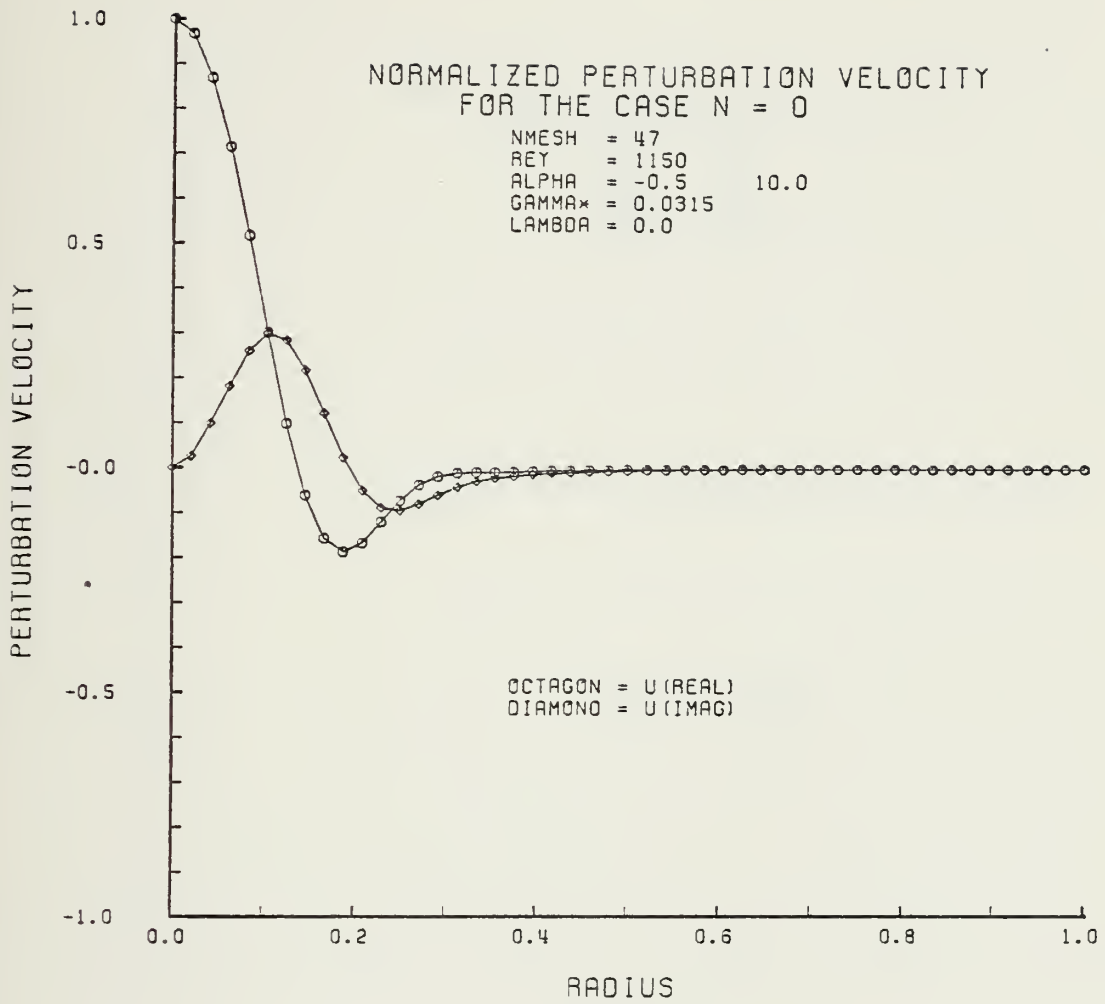


FIGURE 4-4

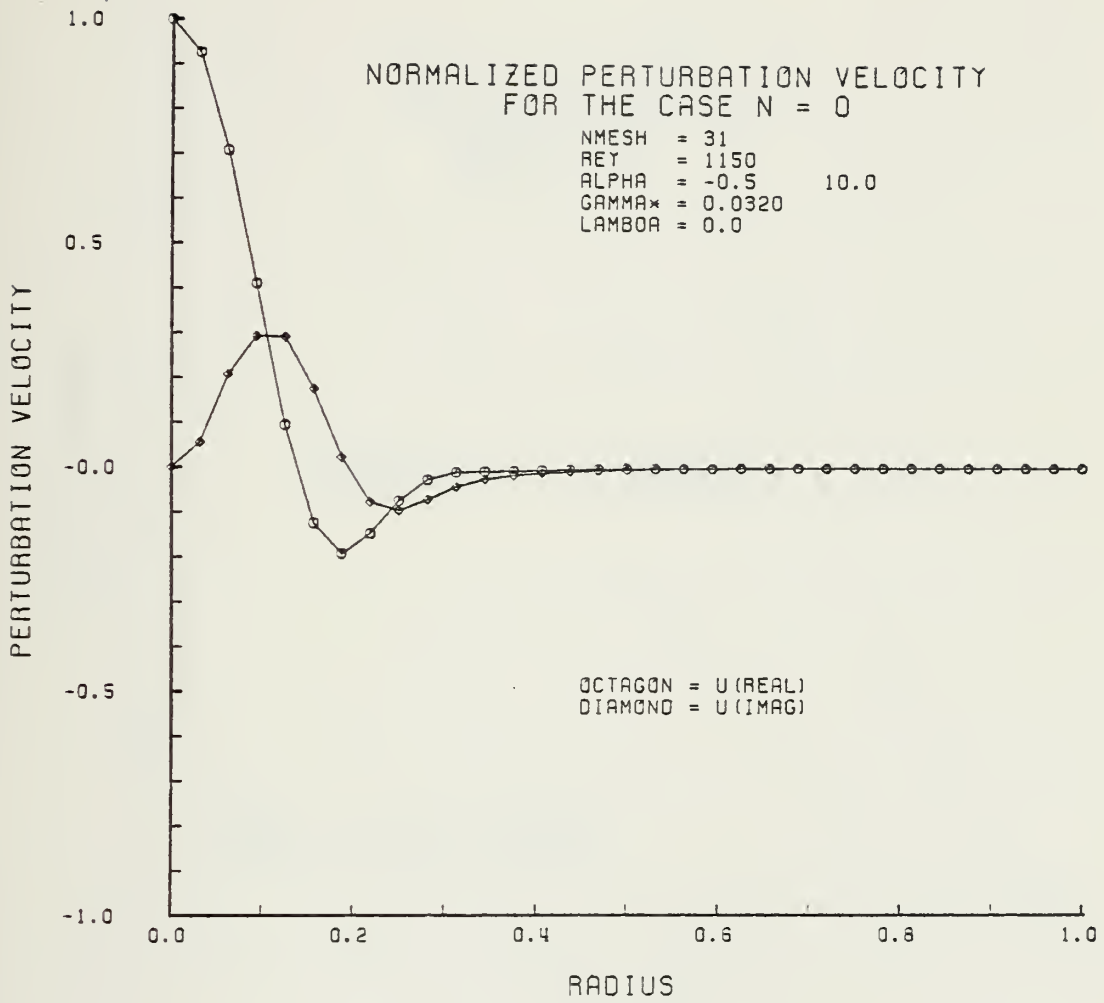


FIGURE 4-5

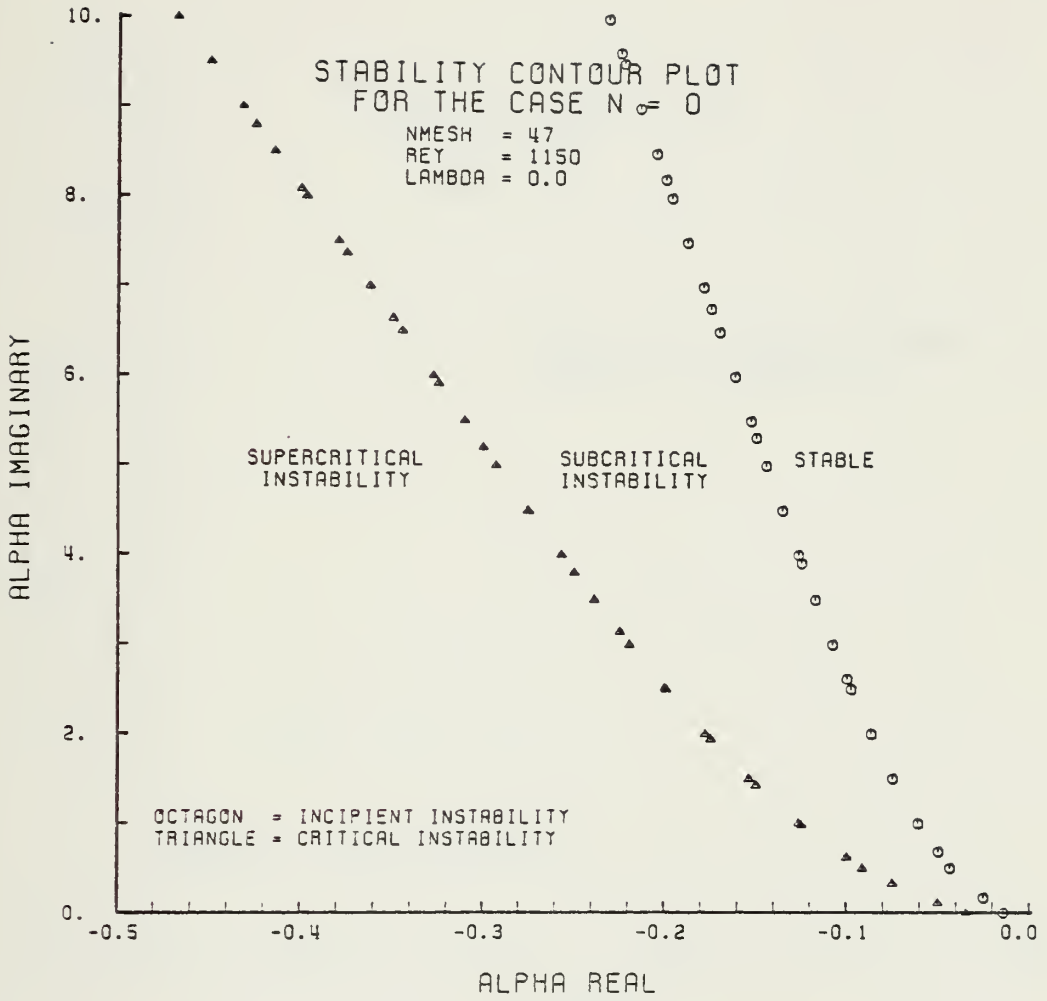


FIGURE 4-6

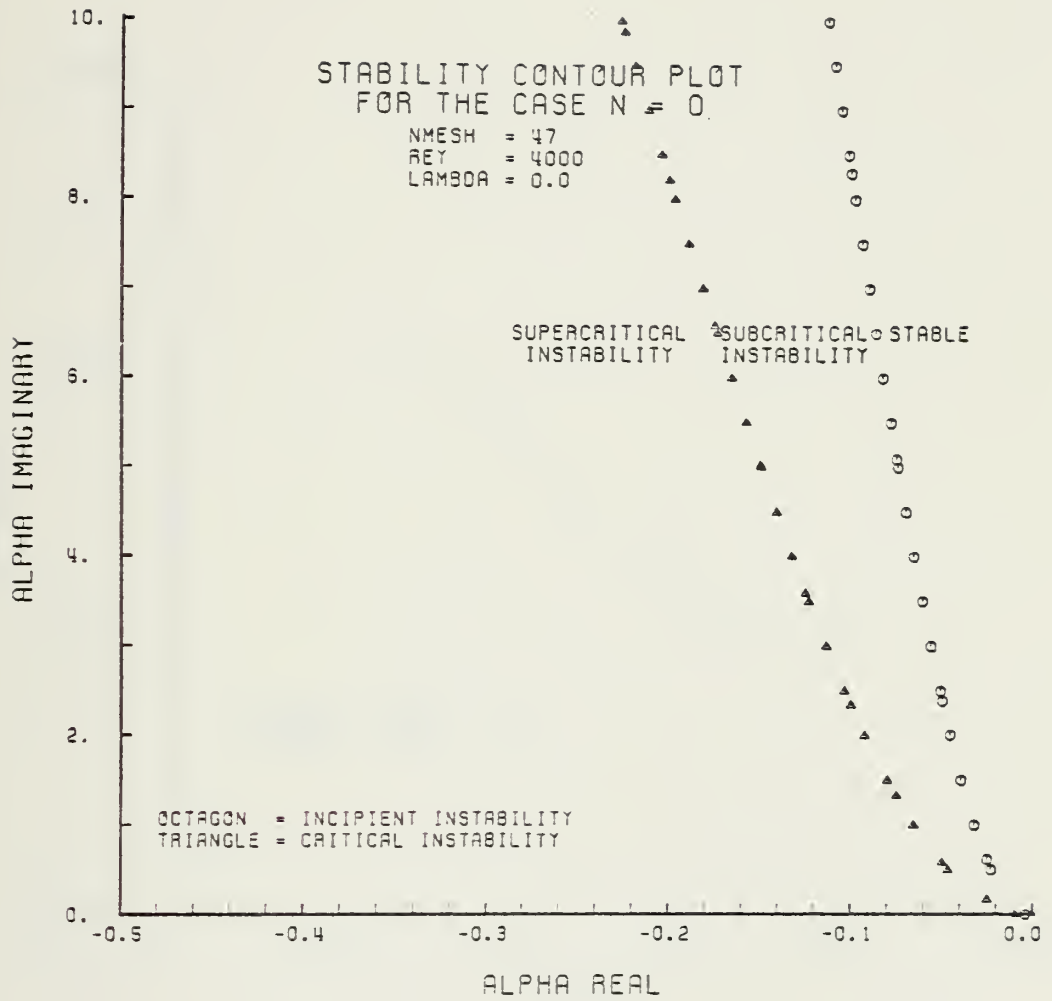


FIGURE 4-7

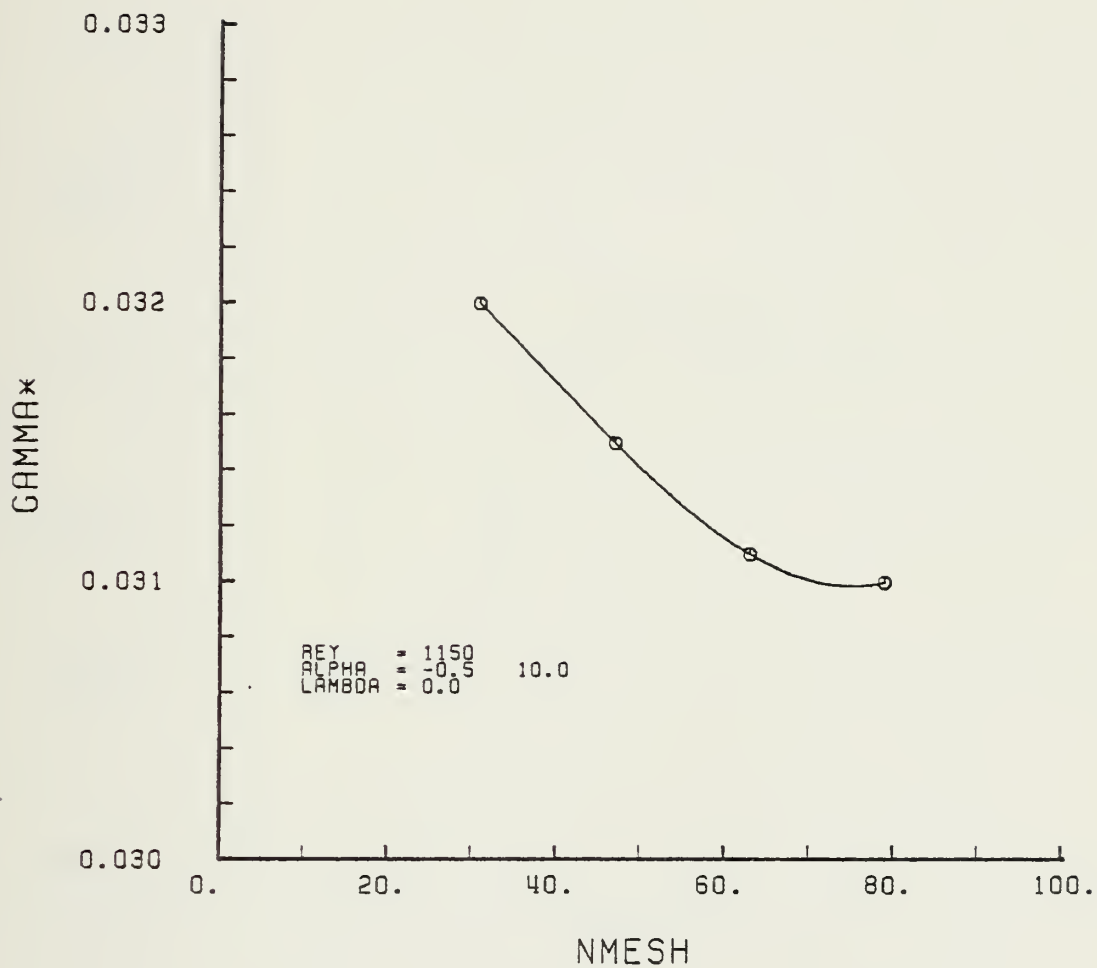


FIGURE 4-8. γ^* Versus Number of Mesh Points, N.

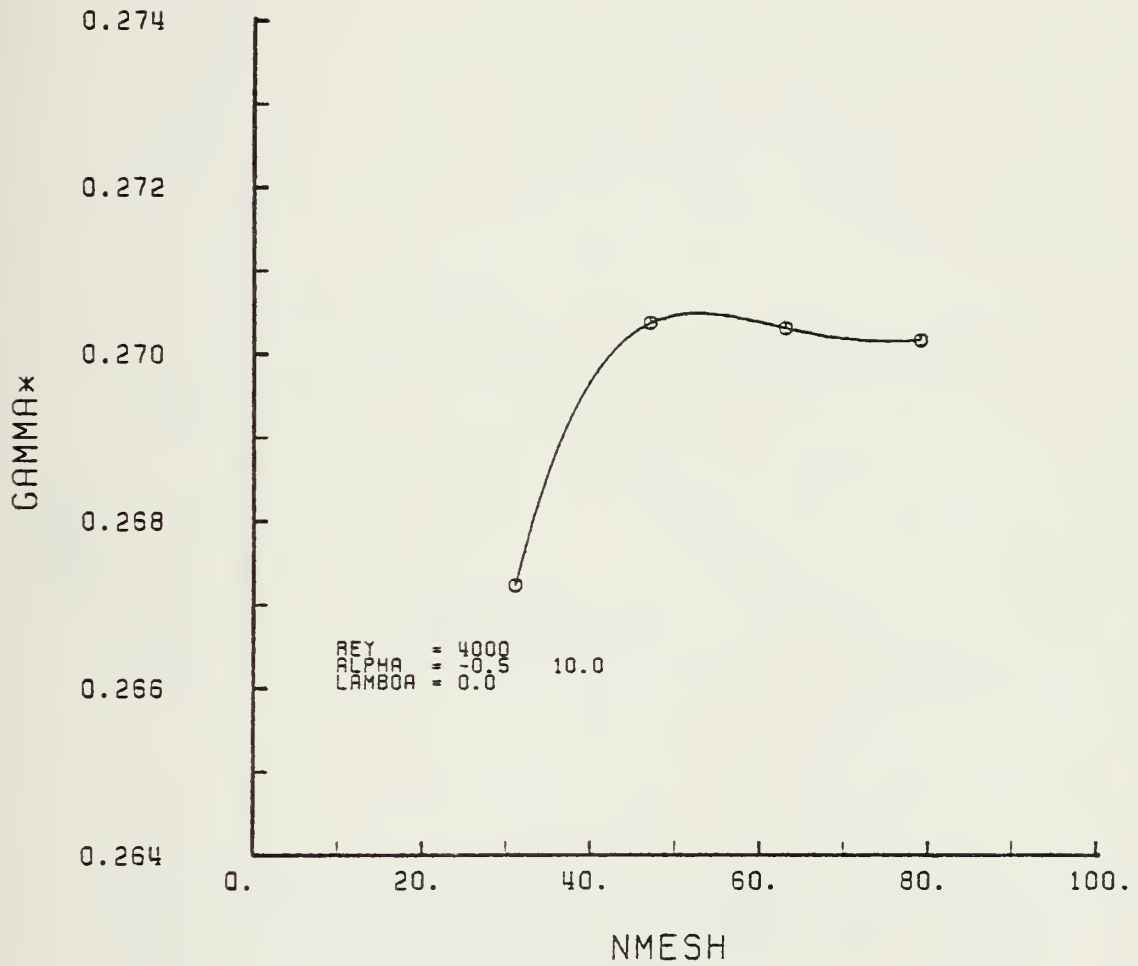


FIGURE 4-9. γ^* Versus Number of Mesh Points, N.

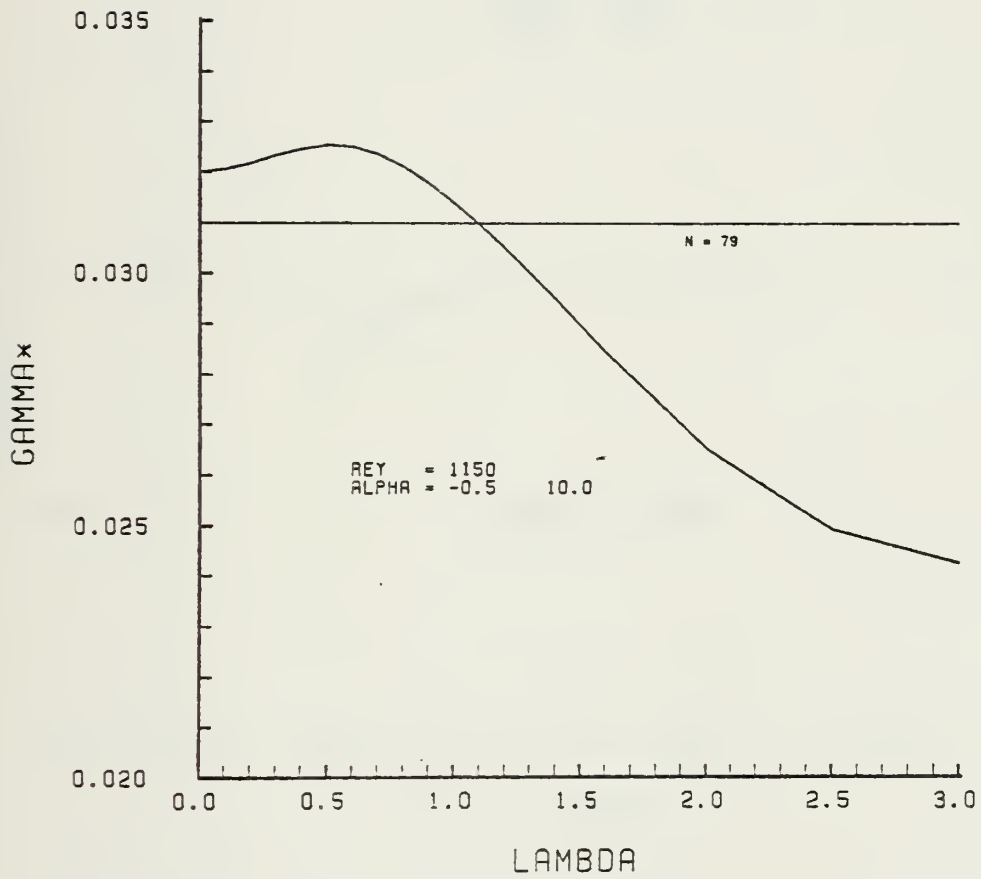


FIGURE 4-10. γ^* Versus Mesh Parameter, Lambda

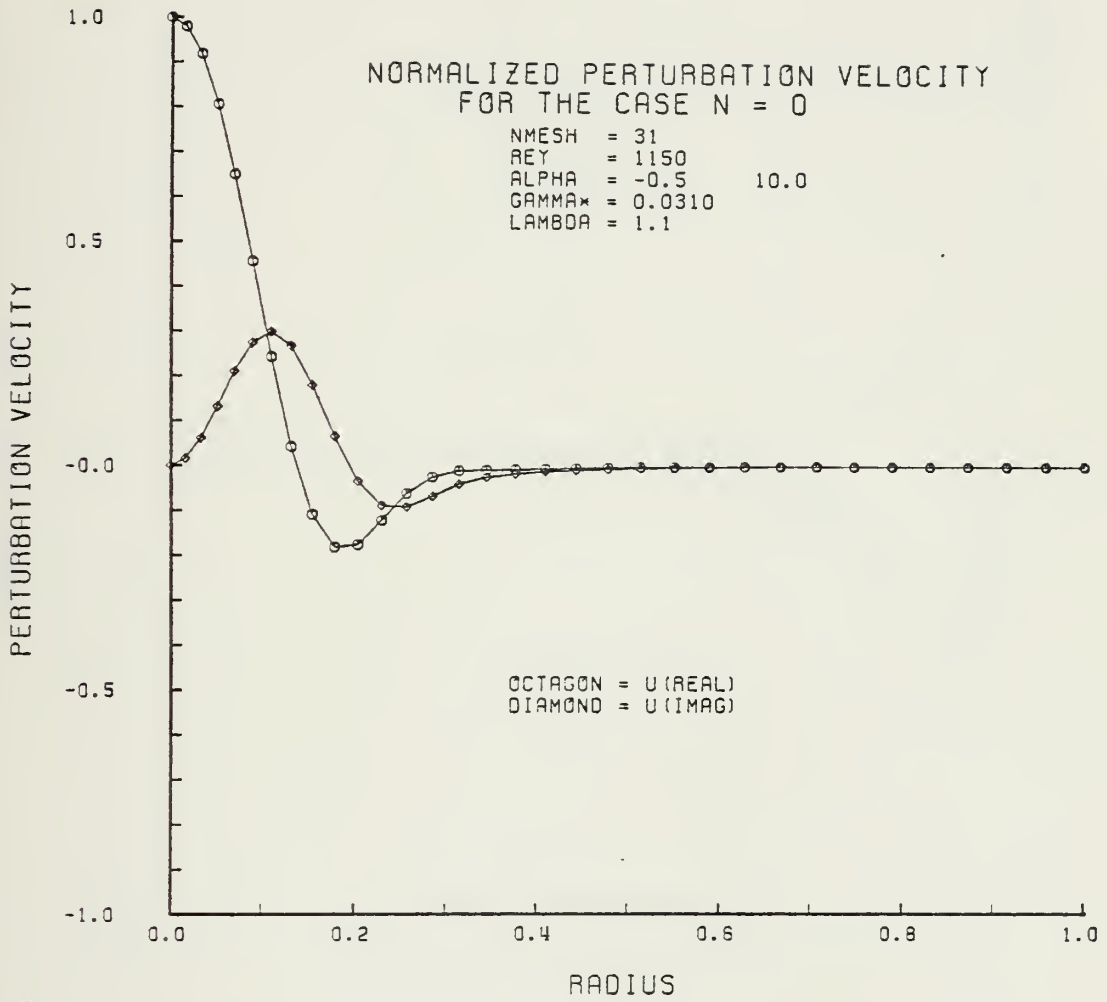


FIGURE 4-11

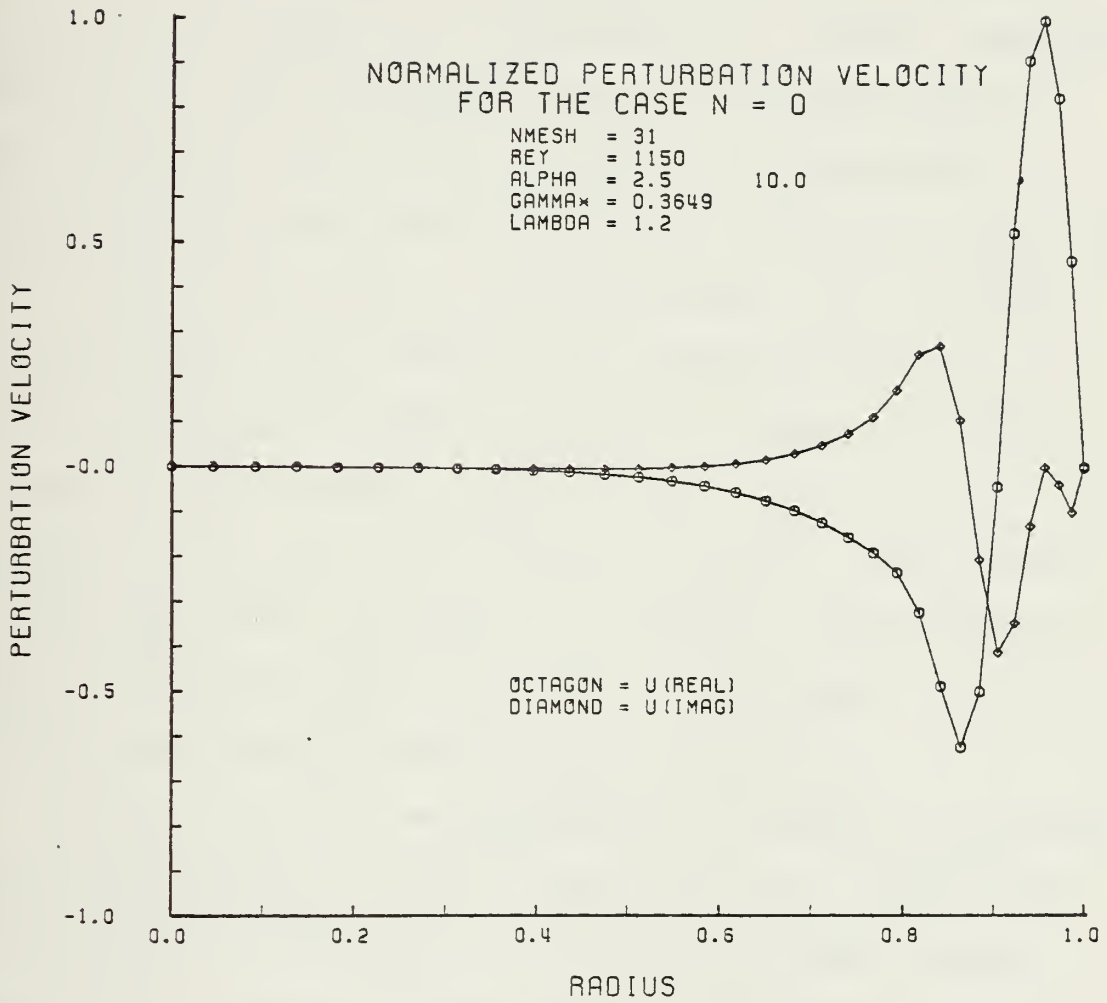


FIGURE 4-12

V. CONCLUSIONS AND RECOMMENDATIONS

The implementation of the newly developed boundary conditions of Gawain [9] has permitted a stable, numerical solution to the linearized vorticity transport equation. The results of the numerical solution are presented in Section IV and show that the stability of pipe Poiseuille flow is governed by the three parameters, α_R , α_I and R_e . In particular, both positive and negative values of α_R , that is, streamwise growth and decay in space, if sufficiently large, produce unstable growth rates in time. This result is new and it is consistent with the known experimental fact that transition to turbulent flow depends not only on Reynolds number but also on the general character of the perturbations which exist in the flow.

The perturbation velocity plots of Section IV represent the first practical look at the function Q . These plots were valuable indicators for adequacy of mesh fineness, that is, N , changes in the nature of the function Q and effects of a nonuniform mesh.

No instabilities were discovered for purely sinusoidal perturbations ($\alpha_R = 0$). This is consistent with the previous investigation of Ref. 11, but should not be assumed for investigations of other angular wave numbers, ($n = 1, 2, 3, \dots$).

Adequate numerical accuracy was proven by demonstrating that the solution was virtually independent of the number of mesh points, N , and that it satisfied to a high degree an independent check of the governing differential equation. This procedure should also be carried out in future investigations prior to conducting full scale data runs.

This study suggests that similar, and perhaps even more rewarding results will be obtained for the higher angular wave numbers. Although lengthy, programming is straightforward if approached systematically. The general organization of the programs of Ref. 4 or Ref. 6 should be helpful in this task. It is recommended that the case for $n = 1$ be undertaken as a follow-on to this study.

The nonuniform computational mesh was shown to be a powerful tool in the reduction of computational time. At the same time, however, the dependence of the mesh offset parameter, λ , on input conditions needs to be investigated further to realize the full potential of this technique.

APPENDIX A

DERIVATION OF VORTICITY TRANSPORT EQUATION COEFFICIENTS

From the change of variable introduced in Ref. 9, the function H for the case $n = 0$ is expressed by

$$H = rQ \quad (A-1)$$

Taking derivatives

$$DH = rDQ + Q \quad (A-2)$$

$$D^2H = rD^2Q + 2DQ \quad (A-3)$$

$$D^3H = rD^3Q + 3D^2Q \quad (A-4)$$

$$D^4H = rD^4Q + 4D^3Q \quad (A-5)$$

Let the '*' superscript denote element (2,2) of matrices (A1) through (A9) of Ref. 9. Since for $n = 0$, only the function H was investigated, equations (2-10) become

$$\begin{aligned} M_4^* D^4H + M_3^* D^3H + M_2^* D^2H + M_1^* DH + M_0^* H \\ - \gamma [N_2^* D^2H + N_1^* DH + N_0^* H] = 0 \end{aligned} \quad (A-6)$$

Substituting for H, equation (A-6) becomes

$$\begin{aligned}
& M_4^* \{rD^4Q + 4D^3Q\} + M_3^* \{rD^3Q + 3D^2Q\} + M_2^* \{rD^2Q + 2DQ\} \\
& + M_1^* \{rDQ + Q\} + M_0^* \{rQ\} - \gamma [N_2^* \{rD^2Q + 2DQ\} \\
& + N_1^* \{rDQ + Q\} + N_0^* \{rQ\}] = 0 \tag{A-7}
\end{aligned}$$

Before proceeding further, it should be noted that the Ref. 9 matrices from which the coefficients for equation (A-7) were taken were obtained from matrices (2-10) through (2-17) of Ref. 6 by means of the following substitutions:

$$U = 2(1 - r^2) \tag{A-8}$$

$$t = \alpha^2 \frac{n_2}{r^2} \tag{A-9}$$

$$T = \alpha U - \frac{1}{R_e} \left(\alpha^2 - \frac{n_2}{r} \right) \tag{A-10}$$

Defining the new coefficients for equation (A-7) as M_0 through M_4 and N_0 through N_2

$$M_4 = rM_4^* = -\frac{r}{R_e} \tag{A-11}$$

$$M_3 = 4M_4^* + rM_3^* = -\frac{6}{R_e} \tag{A-12}$$

$$M_2 = 3M_3^* + rM_2^* = r\alpha U - \frac{1}{R_e} \left\{ \frac{3}{r} + 2\alpha^2 r \right\} \tag{A-13}$$

$$M_1 = 2M_2^* + rM_1^* = 3\alpha U + \frac{3}{R_e} \left\{ \frac{1}{r^2} - 2\alpha^2 \right\} \tag{A-14}$$

$$M_0 = M_1^* + rM_0^* = r\alpha^3 U - \frac{\alpha^4 r}{R_e} \tag{A-15}$$

$$N_2 = rN_2^* = -r \quad (\text{A-16})$$

$$N_1 = 2N_2^* + rN_1^* = -3 \quad (\text{A-17})$$

$$N_0 = N_1^* + rN_0^* = -\alpha^2 r \quad (\text{A-18})$$

Upon making use of the foregoing substitutions, the governing relation can finally be reduced to the form previously shown in equation (3-1).

APPENDIX B

FINITE DIFFERENCE EQUATIONS

Improved finite difference equations for the boundaries were obtained by not using the virtual point method of Ref. 4 and Ref. 6 and deriving the forms directly from the boundary conditions of Appendix A. The equations thus formed are also of consistent order truncation error, significantly improving the accuracy of the solution [Ref. 8].

Because of a peculiarity in the form of the consistent second order truncation error equations at the axis, a singularity resulted for α equal to zero. Consistent third order truncation error equations eliminated this problem.

From Appendix A, the axis boundary conditions are

$$DQ(0) = 0 \quad \text{and} \quad D^3Q(0) = 0 \quad (\text{B-1})$$

Representing Q by a power series and applying equations (B-1) yields

$$\begin{aligned} Q(r) = & Q(0) + D^2Q(0)\frac{r^2}{2!} + D^4Q(0)\frac{r^4}{4!} + D^5Q(0)\frac{r^5}{5!} \\ & + D^6Q(0)\frac{r^6}{6!} + \dots \end{aligned} \quad (\text{B-2})$$

Using five mesh points at $r = \delta, 2\delta, 3\delta, 4\delta$ and 5δ results in the matrix

$$\begin{pmatrix} Q_1 \\ Q_2 \\ Q_3 \\ Q_4 \\ Q_5 \end{pmatrix} = \begin{bmatrix} 1 & \frac{1}{2} & \frac{1}{24} & \frac{1}{120} & \frac{1}{720} \\ 1 & 2 & \frac{16}{24} & \frac{32}{120} & \frac{64}{720} \\ 1 & \frac{9}{2} & \frac{81}{24} & \frac{243}{120} & \frac{729}{720} \\ 1 & 8 & \frac{256}{24} & \frac{1024}{120} & \frac{4096}{720} \\ 1 & \frac{25}{2} & \frac{625}{24} & \frac{3125}{120} & \frac{15625}{720} \end{bmatrix} \begin{pmatrix} Q(0) \\ \delta^2 D^2 Q(0) \\ \delta^4 D^4 Q(0) \\ \delta^5 D^5 Q(0) \\ \delta^6 D^6 Q(0) \end{pmatrix} + O\delta^7$$

(B-3)

Differentiating equation (B-2) and substituting $r = \delta$ gives
(in matrix form)

$$\begin{pmatrix} Q(\delta) \\ \delta D Q(\delta) \\ \delta^2 D^2 Q(\delta) \\ \delta^3 D^3 Q(\delta) \\ \delta^4 D^4 Q(\delta) \end{pmatrix} = \begin{bmatrix} 1 & \frac{1}{2!} & \frac{1}{4!} & \frac{1}{5!} & \frac{1}{6!} \\ 0 & 1 & \frac{1}{3!} & \frac{1}{4!} & \frac{1}{5!} \\ 0 & 1 & \frac{1}{2!} & \frac{1}{3!} & \frac{1}{4!} \\ 0 & 0 & 1 & \frac{1}{2!} & \frac{1}{3!} \\ 0 & 0 & 1 & 1 & \frac{1}{2!} \end{bmatrix} \begin{pmatrix} Q(0) \\ \delta^2 D^2 Q(0) \\ \delta^4 D^4 Q(0) \\ \delta^5 D^5 Q(0) \\ \delta^6 D^6 Q(0) \end{pmatrix} + O\delta^7$$

(B-4)

Let [A] and [B] denote the coefficient matrices of equations (B-3) and (B-4) respectively. The values of $Q(0)$, $\delta^2 D^2 Q(0)$, $\delta^4 D^4 Q(0)$, $\delta^5 D^5 Q(0)$ and $\delta^6 D^6 Q(0)$ may be solved for by

$$\begin{pmatrix} Q(0) \\ \delta^2 D^2 Q(0) \\ \delta^4 D^4 Q(0) \\ \delta^5 D^5 Q(0) \\ \delta^6 D^6 Q(0) \end{pmatrix} = [A]^{-1} \begin{pmatrix} Q_1 \\ Q_2 \\ Q_3 \\ Q_4 \\ Q_5 \end{pmatrix} + O\delta^7 \quad (\text{B-5})$$

Putting equation (B-5) into equation (B-4),

$$\begin{pmatrix} Q(\delta) \\ \delta D Q(\delta) \\ \delta^2 D^2 Q(\delta) \\ \delta^3 D^3 Q(\delta) \\ \delta^4 D^4 Q(\delta) \end{pmatrix} = [B][A]^{-1} \begin{pmatrix} Q_1 \\ Q_2 \\ Q_3 \\ Q_4 \\ Q_5 \end{pmatrix} + O\delta^7 \quad (\text{B-6})$$

The last line of this set of equations gives

$$\begin{aligned} D^4 Q(\delta) = \frac{1}{\delta^4} & (-.911564626Q_1 + 2.750242955Q_2 - 3.043731779Q_3 \\ & + 1.42468416Q_4 - .219630709Q_5) + O\delta^3 \end{aligned} \quad (\text{B-7})$$

To solve for $D^3 Q(\delta)$, the rightmost column and bottom row are eliminated from matrices [A] and [B] then these new matrices are inserted into equations (B-5) and (B-6).

The bottom line of equation (B-6) will now give the expression for $D^3Q(\delta)$ with a consistent third order truncation error. $D^2Q(\delta)$ and $DQ(\delta)$ were solved for in a similar manner.

$$D^3Q(\delta) = \frac{1}{\delta^3}(1.825165563Q_1 - 3.250331126Q_2 + 1.660927152Q_3 - .235761589Q_4) + O\delta^3 \quad (B-8)$$

$$D^2Q = \frac{1}{\delta^2}(-\frac{35}{60}Q_1 + \frac{8}{15}Q_2 + \frac{1}{20}Q_3) + O\delta^3 \quad (B-9)$$

$$DQ = \frac{1}{\delta}(-\frac{2}{3}Q_1 + \frac{2}{3}Q_2) + O\delta^3 \quad (B-10)$$

Due to the complexity of the boundary conditions, it was decided that consistent third order truncation error equations should also be used at $r = 2\delta$. For this the [B] matrix only need be changed as equation (B-2) is unchanged at this station. The new matrix [B] is formed by differentiating equation (B-2) and making the substitution $r = 2\delta$. Proceeding as for $r = \delta$ gives the following finite difference approximations

$$D^4Q(2\delta) = \frac{1}{\delta^4}(-3.10340136Q_1 + 6.903012634Q_2 - 5.342274053Q_3 + 1.66083577Q_4 - 0.123420797Q_5) + O\delta^3 \quad (B-11)$$

$$D^3Q(2\delta) = \frac{1}{\delta^3}(.868874172Q_1 - .937748345Q_2 - .254304636Q_3 + .323178808Q_4) + O\delta^3 \quad (B-12)$$

$$D^2Q(2\delta) = \frac{1}{\delta^2} \left(\frac{11}{12}Q_1 - \frac{28}{15}Q_2 + \frac{19}{20}Q_3 \right) + O\delta^3 \quad (\text{B-13})$$

$$DQ(2\delta) = \frac{1}{\delta} \left(-\frac{4}{3}Q_1 + \frac{4}{3}Q_2 \right) + O\delta^3 \quad (\text{B-14})$$

It should also be noted that the value of Q at $r = 0$ may be solved for from the top line of equations (B-5)

$$Q(0) = (1.795918367Q_1 - 1.24781341Q_2 + .606413994Q_3 - .177842566Q_4 + .023323615Q_5) + O\delta^3 \quad (\text{B-15})$$

The central difference equations given by Ref. 6 were already consistent second order truncation error equations as confirmed by Ref. 8 and were retained.

For the wall, the clamped end, consistent second order equations (5) through (8) of Table II, Ref. 8 were modified for the "right boundary" using the procedure given in Section 5 of that reference.

$$D^4Q(1-\delta) = \frac{1}{\delta^4} \left(-\frac{1}{4}Q_{N-3} + \frac{8}{3}Q_{N-2} - 9Q_{N-1} + 16Q_N \right) + O\delta^2 \quad (\text{B-16})$$

$$D^3Q(1-\delta) = \frac{1}{\delta^3} \left(-\frac{1}{3}Q_{N-2} + 3Q_N \right) + O\delta^2 \quad (\text{B-17})$$

$$D^2Q(1-\delta) = \frac{1}{\delta^2} (Q_{N-1} - 2Q_N) + O\delta^2 \quad (\text{B-18})$$

$$DQ(1-\delta) = \frac{1}{\delta} \left(-\frac{1}{2}Q_{N-1} \right) + O\delta^2 \quad (\text{B-19})$$

Since the wall finite difference approximations were of only second order truncation error, the approximations for DQ through D^4Q at $r = 1-2\delta$ were obtained directly from the central difference equations with $Q(1) = 0$.

$$D^4Q(1-2\delta) = \frac{1}{\delta^4}(Q_{N-3} - 4Q_{N-2} + 6Q_{N-1} - 4Q_N) + O\delta^2 \quad (B-20)$$

$$D^3Q(1-2\delta) = \frac{1}{\delta^3}\left(-\frac{1}{2}Q_{N-3} + Q_{N-2} - Q_N\right) + O\delta^2 \quad (B-21)$$

$$D^2Q(1-2\delta) = \frac{1}{\delta^2}(Q_{N-2} - 2Q_{N-1} + Q_N) + O\delta^2 \quad (B-22)$$

$$DQ(1-2\delta) = \frac{1}{\delta}\left(-\frac{1}{2}Q_{N-2} + \frac{1}{2}Q_N\right) + O\delta^2 \quad (B-23)$$

APPENDIX C

NONUNIFORM MESH

To control the distribution of a fixed number of mesh points, a change of the independent variable from r to η was performed.

$$Q = Q(\eta) \quad (C-1)$$

$$r = r(\eta) \quad (C-2)$$

The derivative with respect to r becomes

$$D = (D^* r)^{-1} D^* \quad (C-3)$$

where

$$D^* = \frac{d}{d\eta} \quad \text{and} \quad D = \frac{d}{dr} \quad (C-4)$$

$DQ, D^2Q \dots$ can now be expressed in terms of the new independent variable, η .

$$DQ = (D^* r)^{-1} D^* Q \quad (C-5)$$

$$\begin{aligned} D^2Q &= D(DQ) = (D^* r)^{-1} D^* (DQ) \\ &= (D^* r)^{-2} D^{*2} Q - (D^* r)^{-3} (D^{*2} r) D^* Q \end{aligned} \quad (C-6)$$

$$\begin{aligned}
D^3 Q &= D(D^2 Q) = (D^* R)^{-1} D^* (D^2 Q) \\
&= (D^* r)^{-3} D^{*3} Q - 3(D^* r)^{-4} (D^{*2} r) D^{*2} Q \\
&\quad - [(D^* r)^{-4} (D^{*3} r) - 3(D^* r)^{-5} (D^{*2} r)^2] D Q \quad (C-7)
\end{aligned}$$

$$\begin{aligned}
D^4 Q &= D(D^3 Q) = (D^* r)^{-1} D^* (D^3 Q) \\
&= (D^* r)^{-4} D^{*4} Q - 6(D^* r)^{-5} (D^{*2} r) D^{*3} Q \\
&\quad + [15(D^* r)^{-6} (D^{*2} r) - 4(D^* r)^{-5} (D^{*3} r)] D^{*2} Q \\
&\quad - [15(D^* r)^{-7} (D^{*2} r)^3 - 10(D^* r)^{-6} (D^{*2} r) (D^{*3} r) \\
&\quad\quad + (D^* r)^{-5} (D^{*4} r)] D Q \quad (C-8)
\end{aligned}$$

The derivatives of Q with respect to r can now be written

$$D Q = f_{11} D^* Q \quad (C-9)$$

$$D^2 Q = f_{22} D^{*2} Q + f_{21} D^* Q \quad (C-10)$$

$$D^3 Q = f_{33} D^{*3} Q + f_{32} D^{*2} Q + f_{31} D^* Q \quad (C-11)$$

$$D^4 Q = f_{44} D^{*4} Q + f_{43} D^{*3} Q + f_{42} D^{*2} Q + f_{41} D^* Q \quad (C-12)$$

where

$$f_{11} = (D^* r)^{-1} \quad (C-13)$$

$$f_{22} = (D^* r)^{-2} \quad (C-14)$$

$$f_{21} = -(D^* r)^{-3} (D^{*2} r) \quad (C-15)$$

$$f_{33} = (D^* r)^{-3} \quad (C-16)$$

$$f_{32} = -3(D^* r)^{-4} (D^{*2} r) \quad (C-17)$$

$$f_{31} = 3(D^* r)^{-5} (D^{*2} r)^2 - (D^* r)^{-4} (D^{*3} r) \quad (C-18)$$

$$f_{44} = (D^* r)^{-4} \quad (C-19)$$

$$f_{43} = -6(D^* r)^{-5} (D^{*2} r) \quad (C-20)$$

$$f_{42} = 15(D^* r)^{-6} (D^{*2} r)^2 - 4(D^* r)^{-5} (D^{*3} r) \quad (C-21)$$

$$f_{41} = -15(D^* r)^{-7} (D^{*2} r)^3 + 10(D^* r)^{-6} (D^{*2} r) (D^{*3} r) \\ - (D^* r)^{-5} (D^{*4} r) \quad (C-22)$$

Substituting equations (C-9) through (C-12) into the vorticity transport equation (A-6) yields

$$M_4^* D^{*4} Q + M_3^* D^{*3} Q + M_2^* D^{*2} Q + M_1^* D^* Q + M_0^* Q \\ - \gamma [N_2^* D^{*2} Q + N_1^* D^* Q + N_0^* Q] = 0 \quad (C-23)$$

where

$$M_4^* = M_4 f_{44} \quad (C-24)$$

$$M_3^* = M_4 f_{43} + M_3 f_{33} \quad (C-25)$$

$$M_2^* = M_4 f_{42} + M_3 f_{32} + M_2 f_{22} \quad (C-26)$$

$$M_1^* = M_4 f_{41} + M_3 f_{31} + M_2 f_{21} \quad (C-27)$$

$$M_0^* = M_0 \quad (C-28)$$

$$N_2^* = N_2 f_{22} \quad (C-29)$$

$$N_1^* = N_2 f_{21} + N_1 f_{11} \quad (C-30)$$

$$N_0^* = N \quad (C-31)$$

In order to concentrate the mesh points at the axis, the function

$$r = 1 - C \tanh \lambda(1-\eta) \quad (C-32)$$

was chosen where λ is a parameter controlling the degree of concentration of mesh points near the axis. Equation (C-32) must satisfy the two conditions

$$r = 0 \quad \text{at} \quad \eta = 0 \quad (C-33)$$

and

$$r = 1 \quad \text{at} \quad \eta = 1 .$$

Substituting equation (C-33) into (C-32) gives

$$C = 1/\tanh \lambda . \quad (C-35)$$

Computing derivatives

$$D^* r = C\lambda/\cosh^2 \lambda(1-\eta) \quad (C-36)$$

$$D^{*2} r = 2C\lambda^2 [\tanh \lambda(1-\eta)/\cosh^2 \lambda(1-\eta)] \quad (C-37)$$

$$D^{*3} r = -2C\lambda^3 \{ [1-2\sinh^2 \lambda(1-\eta)]/\cosh^4 \lambda(1-\eta) \} \quad (C-38)$$

$$D^{*4} r = 8C\lambda^4 [\tanh^3 \lambda(1-\eta)/\cosh^2 \lambda(1-\eta)] \quad (C-39)$$

To shift the mesh point concentration to the wall, the function

$$r = C \tanh \lambda \eta \quad (C-40)$$

was selected. Satisfying equations (C-33) and (C-34) for this equation also gives equation (C-35). The derivatives

of (C-40) are given by equations (C-36) through (C-39) if η is substituted for all occurrences of $(1-\eta)$ and the signs of equations (C-37) and (C-39) are reversed. Figures C-1 and C-2 show equations (C-32) and (C-40) for four selected values of the parameter λ .

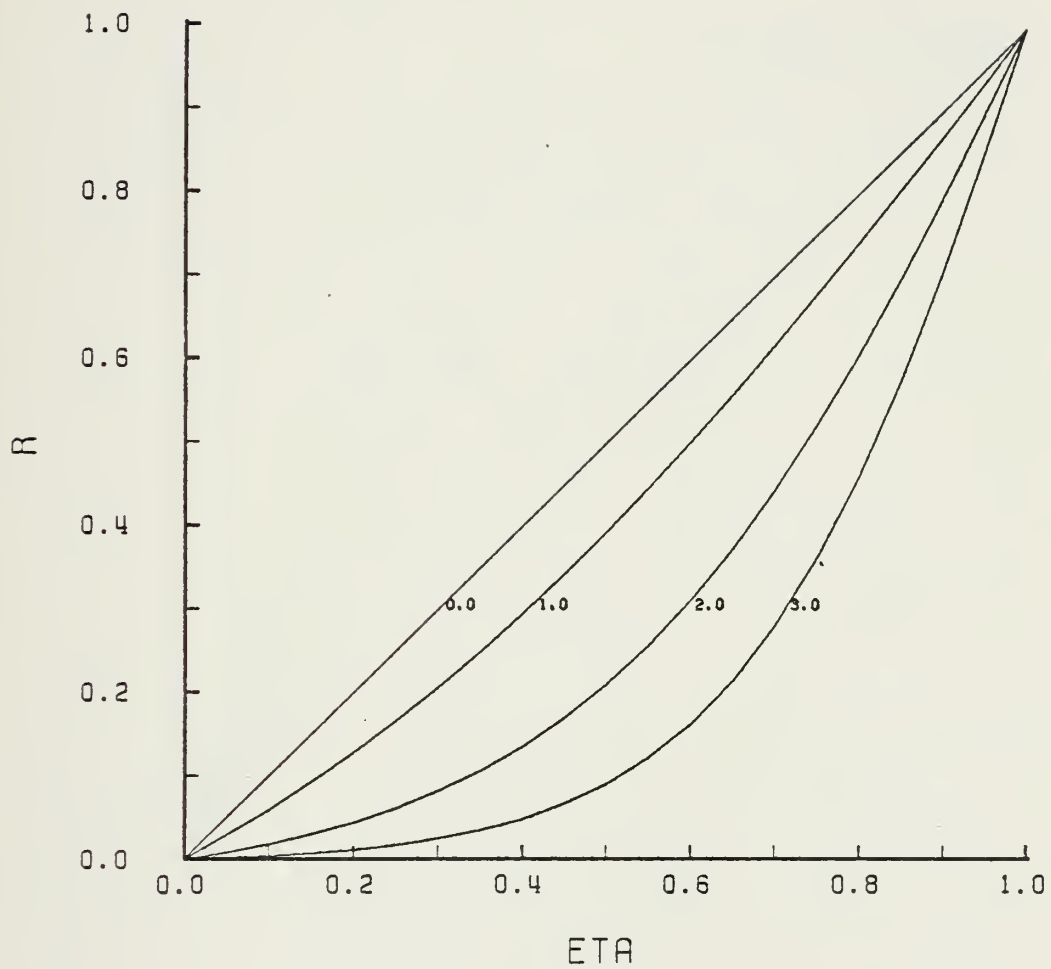


FIGURE C-1. R Versus η for Four Selected Values of Lambda - Axis Offset

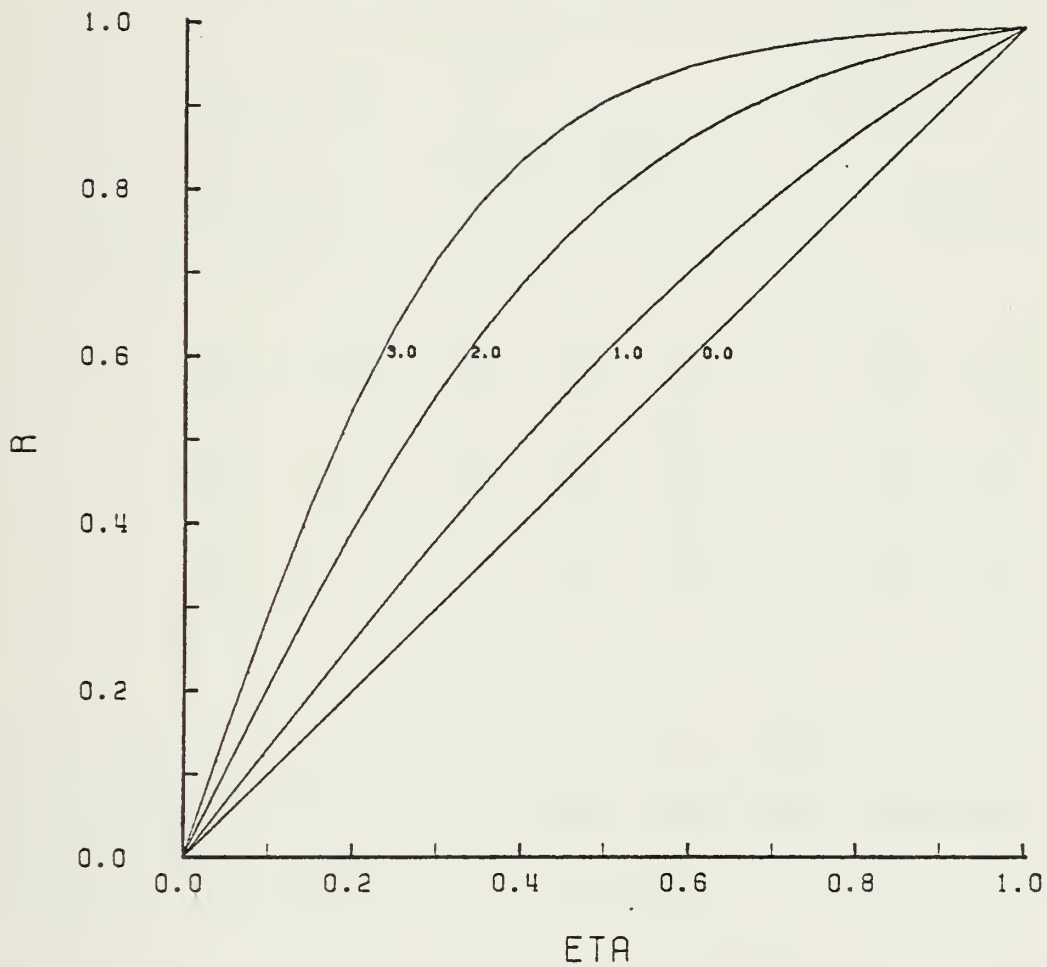


FIGURE C-2. R Versus η for Four Selected Values of Lambda - Wall Offset

APPENDIX D

DERIVATION OF PERTURBATION VELOCITIES

From Ref. 4, Appendix E, equations E-6 through E-8:

$$\begin{Bmatrix} u(r) \\ v(r) \\ w(r) \end{Bmatrix} = [A]\bar{W} + [B]D\bar{W} \quad (D-1)$$

$$= \begin{bmatrix} 0 & -\frac{\beta}{r} & \frac{1}{r} \\ \frac{\beta}{r} & 0 & -\alpha \\ 0 & \alpha & 0 \end{bmatrix} \begin{Bmatrix} F \\ G \\ H \end{Bmatrix} + \begin{bmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ -1 & 0 & 0 \end{bmatrix} \begin{Bmatrix} DF \\ DG \\ DG \end{Bmatrix} \quad (D-2)$$

For this case $\beta = \eta_i = 0$ and $F = DF = 0$. Restricting the investigation to the function H for the reason expressed in Section I and solving for $u(r)$ gives

$$u(r) = \frac{H}{r} + DH \quad (D-3)$$

Performing the change of variable

$$H = rQ \quad (D-4)$$

$$DH = Q + rDQ \quad (D-5)$$

$$u(r) = \frac{rQ}{r} + (Q + rDQ) = 2Q + rDQ \quad (D-6)$$

In order to implement this derivation in a numerical analysis, equation (D-6) was rewritten as

$$u_i = 2Q_i + r_i DQ_i \quad (D-7)$$

Performing the change of independent variable (Appendix C) to accommodate a nonuniform mesh

$$Q_i = Q(\eta_i) \quad (D-8)$$

$$r_i = r(\eta_i) \quad (D-9)$$

$$DQ_i = (D^* r_i)^{-1} D^* Q(\eta_i) \quad (D-10)$$

Substituting equations (D-8), (D-9) and (D-10) into equation (D-7) gives

$$u_i = 2Q(\eta_i) + r(\eta_i) (D^* r_i)^{-1} D^* Q(\eta_i) \quad (D-11)$$

For the axis offset nonuniform mesh, $r(\eta)$ is given by equation (C-32) and $(D^* r)$ by equation (C-36). Substituting into equation (D-11) using equation (C-35) results in

$$\begin{aligned}
u_i &= 2Q(\eta_i) + \left\{ 1 - \frac{\tanh[\lambda(1-\eta_i)]}{\tanh \lambda} \right\} \left\{ \frac{\cosh^2[\lambda(1-\eta_i)]}{C\lambda} \right\} D^* Q(\eta_i) \\
&= 2Q(\eta_i) + \left\{ 1 - \frac{\tanh[\lambda(1-\eta_i)]}{\tanh \lambda} \right\} \frac{\tanh \lambda \cosh^2[\lambda(1-\eta_i)]}{\lambda} \left\{ D^* Q(\eta_i) \right\}
\end{aligned}
\tag{D-12}$$

For the wall offset mesh, equation (C-40) is substituted for equation (C-32) and all occurrences of the term $1-\eta_i$ are replaced by the term η_i .

The value of u at the axis (u_0) and at the wall (u_{N+1}) were solved for by using the boundary conditions specified in Ref. 9, namely

$$Q(1) = 0 \tag{D-13}$$

$$DQ(1) = 0 \tag{D-14}$$

$$DQ(0) = 0 \tag{D-15}$$

$$D^3Q(0) = 0 \tag{D-16}$$

From equations (D-13) and (D-14), using equation (D-7) it is obvious that

$$u_{N+1} = 0 \tag{D-17}$$

and from equations (D-15) and (D-7), it is similarly found that

$$u_0 = 2Q(0) ,$$

(D-18)

where the finite difference approximation for $Q(0)$ is given by equation (B-15).


```

10 PIP0
20 PIP0
30 PIP0
40 PIP0
50 PIP0
60 PIP0
70 PIP0
80 PIP0
90 PIP0
100 PIP0
110 PIP0
120 PIP0
130 PIP0
140 PIP0
150 PIP0
160 PIPC
170 PIP0
180 PIP0
190 PIP0
200 PIP0
210 PIP0
220 PIP0
230 PIP0
240 PIP0
250 PIP0
260 PIP0
270 PIP0
280 PIP0
290 PIP0
300 PIP0
310 PIP0
320 PIP0
330 PIP0
340 PIP0
350 PIP0
360 PIP0
370 PIP0
380 PIP0
390 PIP0
400 PIP0
410 PIP0
420 PIP0
430 PIP0
440 PIP0
450 PIP0
460 PIP0
470 PIP0
480 PIP0

.....
PROGRAM PIPEO (CP/CMS VERSION)
PROGRAM TO INVESTIGATE FLOW STABILITY AND CHARACTERISTICS
FOR THE 3-D CYLINDRICAL FLOW PROBLEM
NI=0; FOR THE FUNCTION Q

TO OBTAIN A FLOW CHART OF THIS PROGRAM, CONSULT NAVAL
POSTGRADUATE SCHOOL TECHNICAL NOTE IN 0141-25, "USER'S
GUIDE TO THE PROGRAMMING AIDS LIBRARY", UNDER PROGRAM
FLOWCH.

-----
IMPLICIT REAL*(A-H,O-Z)
DIMENSION AIPLT(20), REYPLT(20)
COMPLEX *16A,G
INTEGER *4CLOCK(6)
COMMON /COEFNT/ A,G,REY,DEL,R,AMDA

INPUT DESIRED MODE NUMBER

WRITE (6,6) MODENO
READ (5,7) MODENO
CALL IXCLOCK (CLOCK)
WRITE (6,8) CLOCK(3),CLOCK(4)
IF (MODENO.EQ.2) GO TO 2

*****
CALCULATE STABILITY AT A POINT (MODENO = 1 & 3)

OUTPUTS DATA TO FILE FT02F001 COMPATIBLE WITH PROGRAM EIGFCN &
PRINTS THE VALUE OF THE LEAST STABLE EIGENVALUE(GAMMA*) AT THE
CONSOLE FOR EACH SET OF INPUT CONDITIONS. DATA IS ENTERED IN
'D' FORMAT WITH A ZERO EXPONENT (I.E.,1.D0,1150.D0, .5D0,ETC.).
RUN IS TERMINATED WHEN A NEGATIVE VALUE DETERMINED BY THE VALUE
ENTERED. THE TYPE OF MESH OFFSET IS SET EQUAL TO -1; IF
OF LAMBDA. IF LAMBDA < -1.D-10, KSET IS SET EQUAL TO 0.
LAMBDA > 1.D-10, KSET IS SET EQUAL TO 1; OTHERWISE KSET = 0.
IF MODENO = 1, THE EIGENVECTOR CORRESPONDING TO THE LEAST
STABLE EIGENVALUE (GAMMA*) WILL BE WRITTEN TO FILE FT02F001.

```


IF MODENO = 3, THIS OUTPUT IS INHIBITED. MODENO MUST BE
SET EQUAL TO ONE TO GENERATE CORRECT DATA FOR PROGRAM EIGFCN.

```
1 WRITE (6,9) AMDA
READ (5,11) AMDA
KSET = 0
IF (AMDA.LT.-1.D-10) KSET=-1
IF (AMDA.GT.1.D-10) KSET=1
WRITE (6,10)
READ (5,11) AR
WRITE (6,12) AR
READ (5,11) AI
WRITE (6,13) REY
READ (5,11) REY
IF (REY.LE.0.D0) GO TO 5
CALL STAB (AR,AI,GRMAX,KSET,MODENO)
WRITE (6,14) GRMAX
GO TO 1
```

COMPUTE STABILITY MAP (MODENO = 2)

COMPUTES A STABILITY MAP AT MESH POINTS ESTABLISHED BY
THE FOLLOWING PARAMETERS READ FROM FILE FT01FO01:

NXSTP - NO OF MESH PTS IN X-DIRECTION
NYSTP - NO OF MESH PTS IN Y-DIRECTION
N - DIMENSION OF MATRICES X & Y IN SUBROUTINE STAB
DELAR - MAGNITUDE OF THE X-DIRECTION STEP
DELAJ - MAGNITUDE OF THE Y-DIRECTION STEP
REY - REYNOLDS NUMBER

**NOTE - RUN TIME IS LONG IN THIS MODE, SO CP/CMS
RUNS SHOULD BE LIMITED TO 10X10 MESHES OR LESS WITH
N = 31. LARGER RUNS SHOULD BE MADE UNDER BATCH.
LITTLE OS MAY ALSO BE USED AS LESS THAN 196K OF
CORE IS REQUIRED IN THIS MODE FOR N <= 47.

**NOTE - TO RUN THE MAPPING PORTION UNDER OS OR LITTLE OS,
PERFORM THE FOLLOWING:

- 1) RETAIN ALL MAIN PROGRAM SECTIONS BRACKETED BY '-----'.
- 2) CHANGE ALL READ DEVICES TO '5' VICE '1' IN MAIN PROGRAM.
- 3) CHANGE THE DEVICE CODE OF THE LAST WRITE STATEMENT

PIPO 490
PIPO 500
PIPO 510
PIPO 520
PIPO 530
PIPO 540
PIPO 550
PIPO 560
PIPO 570
PIPO 580
PIPO 590
PIPO 600
PIPO 610
PIPO 620
PIPO 630
PIPO 640
PIPO 650
PIPO 660
PIPO 670
PIPO 680
PIPO 690
PIPO 700
PIPO 710
PIPO 720
PIPO 730
PIPO 740
PIPO 750
PIPO 760
PIPO 770
PIPO 780
PIPO 790
PIPO 800
PIPO 810
PIPO 820
PIPO 830
PIPO 840
PIPO 850
PIPO 860
PIPO 870
PIPO 880
PIPO 890
PIPO 900
PIPO 910
PIPO 920
PIPO 930
PIPO 940
PIPO 950
PIPO 960


```

C PRIOR TO THE 'STOP' IN THE MAIN PROGRAM TO '7' VICE '3'. PIP0 970
C CHANGE ALL OTHER MAIN PROGRAM WRITE STATEMENTS TO DEVICE PIP0 980
C CODE '6'. PIP0 990
C IMMEDIATELY AFTER STATEMENT 116 IN MAIN PROGRAM, INSERT PIP01000
C THE FOLLOWING: 'MODENO = 2' PIP01010
C DELETE PORTION BETWEEN '---' MARKINGS IN PIP01020
C SUBROUTINE STAB. PIP01030
C PIP01040
C PIP01050
C PIP01060
C PIP01070
C PIP01080
C PIP01090
C ***** PIP01100 *****
C PIP01110 *****
C PIP01120 -----
C 2 READ (1,15) NXSTP,NYSTP,N,DELAR,DELA1,REY PIP01130
C READ (1,16) ARSTR1,AIRST1 PIP01140
C READ (1,16) AMDA PIP01150
C KSET = 0 PIP01160
C IF (AMDA.LT.-1.D-10) KSET=-1 PIP01170
C IF (AMDA.GT.1.D-10) KSET=1 PIP01180
C AR = ARSTR1 PIP01190
C WRITE (6,17) REY PIP01200
C WRITE (3,18) NXSTP PIP01210
C DO 4 I=1,NXSTP PIP01220
C AI = AIRST1 PIP01230
C DO 3 J=1,NYSTP PIP01240
C CALL STAB (AR,AI,GRMAX,KSET,MODENO) PIP01250
C WRITE (6,19) AR,AI,GRMAX PIP01260
C WRITE (3,20) AR,AI,GRMAX PIP01270
C 3 AI = AI+DELA1 PIP01280
C 4 AR = AR+DELAR PIP01290
C PIP01300
C PIP01310
C PIP01320
C PIP01330
C PIP01340
C PIP01350
C PIP01360
C PIP01370
C PIP01380
C PIP01390
C PIP01400
C PIP01410
C PIP01420
C PIP01430
C PIP01440

```



```

10 FORMAT (I5X, 'INPUT ALPHA REAL')
11 FORMAT (D20.10)
12 FORMAT (I5X, 'INPUT ALPHA IMAG')
13 FORMAT (I5X, 'INPUT REYNOLDS NO')
14 FCRMAT (I5X, 'STAB=', D20.10)
15 FORMAT (3I2, 3D20.10)
16 FCRMAT (2D20.10)
17 FORMAT (I5X, 'REYNOLDS NUMBER = ', F8.1, '/', I0X, 'AR', '18X,
18 'AI', '16X, 'STAB', '/')
18 FORMAT (I2)
19 FORMAT (I5X, 'STOP TIME=', D20.10)
20 FORMAT (3E20.10)
21 FORMAT ('0', 'STOP TIME=', 2A4)
END

```

```

.....SUBROUTINE STAB(AR,AI,GRMAX,KSET,MODENO).....
PURPOSE
  RETURNS THE REAL PORTION OF THE LEAST STABLE EIGENVALUE
  FOR THE GIVEN INPUT CONDITIONS. THIS VALUE DETERMINES THE
  STABILITY OF THE FLOW.
USAGE
  CALL STAB(AR,AI,GRMAX,KSET,MODENO)
DESCRIPTION OF PARAMETERS
  AR - THE REAL PART OF ALPHA
  AI - THE IMAGINARY PART OF ALPHA.
  GRMAX - THE REAL PCRTION OF THE LEAST STABLE EIGENVALUE.
  THIS VALUE IS RETURNED TO THE CALLING PROGRAM.
  N - THE NUMBER OF INTERIOR MESH POINTS
  KSET - AN INTEGER DENOTING THE TYPE OF MESH OFFSET
  USED.
  KSET = -1 FOR WALL OFFSET
  KSET = 0 FOR UNIFORM MESH
  KSET = 1 FOR AXIS OFFSET

```


STABI250
 STABI260
 STABI270
 STABI280
 STABI290
 STABI300
 STABI310
 STABI320
 STABI330
 STABI340
 STABI350
 STABI360
 STABI370
 STABI380
 STABI390

```

C      DO 2 I=1,MDIM
C      RADIUS(I) = DFLOAT(I)*DELL0
C      WRITE (2,7) RADIUS(I),ZR(I,NEIG),ZI(I,NEIG)
C-----
C      2 CONTINUE
C      3 RETURN
C      4 FCRMAT ('0* * * ERROR NUMBER',I7,' ON EIGENVALUE',
C      1 I7,' * * *',//)
C      5 FCRMAT ('12',3D20.10)
C      6 FCRMAT ('12')
C      7 FCRMAT ('F15.7,2(1PD20.10)')
C      END

```

MSTI 10
 MSTI 20
 MSTI 30
 MSTI 40
 MSTI 50
 MSTI 60
 MSTI 70
 MSTI 80
 MSTI 90
 MSTI 100
 MSTI 110
 MSTI 120
 MSTI 130
 MSTI 140
 MSTI 150
 MSTI 160
 MSTI 170
 MSTI 180
 MSTI 190
 MSTI 200
 MSTI 210
 MSTI 220
 MSTI 230
 MSTI 240
 MSTI 250
 MSTI 260
 MSTI 270
 MSTI 280
 MSTI 290
 MSTI 300
 MSTI 310

```

C.....SUBROUTINE MSET1(ETA,CQM1,CQM2,KSET).....
C
C      PURPOSE
C      MSET1 GENERATES THE COEFFICIENTS FOR THE FINITE DIFFERENCE
C      APPROXIMATION OF THE COMPONENT Q.
C      USAGE
C      CALL MSET1(ETA,CQM1,CQM2,KSET)
C
C      DESCRIPTION OF PARAMETERS
C      ETA - INDEPENDENT VARIABLE REPLACING R
C            IN NONUNIFORM MESH.
C      CQM1 - COEFFICIENTS OF Q AND ITS DERIVATIVES
C            IN THE FINITE DIFFERENCE APPROXIMATION OF
C            THE NON-GAMMA TERMS. CQM1(1) IS THE
C            COEFFICIENT FOR Q, CQM1(2) IS THE COEF-
C            FICIENT FOR DQ, AND SO ON TO CQM1(5).
C      CQM2 - SAME AS CQM1 EXCEPT GAMMA TERMS AND
C            DIMENSIONED 3 INSTEAD OF 5.
C      KSET - MESH OFFSET PARAMETER AS DESCRIBED FOR
C            SUBROUTINE STAB.
C
C      OTHER ROUTINES NEEDED
C      NONE

```



```

C.....
C.....
C.....
C.....
SUBROUTINE MSET1 (ETA,CQM1,CQM2,KSET)
IMPLICIT REAL*8(A-H,O-Z)
COMPLEX *16A,G,CQM1(5),CQM2(3),M4,M3,M2,M1,M0,N2,N1,N0
COMMON /COEFNT/ A,G,REY,DEL,AMDA
C.....
C.....
C.....
C.....
AMDA = DABS(AMDA)
IF (KSET.EQ.0) GO TO 1
TETA = ETA
IF (KSET.EQ.1) TETA=1D0-ETA
ETAP = AMDA*TETA
CNST = 1D0/DTANH(AMDA)
R = 1D0-CNST*DTANH(ETAP)
IF (KSET.EQ.-1) R = CNST*DTANH(ETAP)
C.....
C.....
C.....
C.....
S1 = DSINH(ETAP)
C1 = DCOSH(ETAP)
T1 = DTANH(ETAP)
C.....
C.....
C.....
C.....
D4N = (8D0*CNST*AMDA**4)*(T1**3/C1**2)
IF (KSET.EQ.-1) D4N=-D4N
D3N = (-2D0*CNST*AMDA**3)*(1D0-2DC*S1**2)/C1**4
D2N = (2D0*CNST*AMDA**2)*T1/C1**2
IF (KSET.EQ.-1) D2N=-D2N
D1N = (CNST*AMDA)*(1D0/C1**2)
C.....
C.....
C.....
C.....
F11 = 1D0/D1N
F21 = 1D0/D1N**2
F31 = (-1D0/D1N**3)*D2N
F41 = 1D0/D1N**4
F12 = 1D0/D1N**2
F22 = (-1D0/D1N**3)*D2N
F32 = 1D0/D1N**3
F42 = (-3D0/D1N**4)*D2N
F13 = (3D0/D1N**5)*D2N**2-((1D0/D1N**4)*D3N
F23 = (-6D0/D1N**6)*D2N**2-(4D0/D1N**5)*D3N
F33 = (15D0/D1N**7)*D2N**3+(1D0/D1N**6)*D2N**3N-((1D0/D1N**5)*D4N
F43 = (-15D0/D1N**4)*D2N
C.....
C.....
C.....
C.....
1 IF (KSET.EQ.0) R = ETA
      DEFINE THE RECURRING PARAMETER U(R).
C.....
C.....
C.....
C.....
      U = 2D0*(1D0-R**2)
C.....
C.....
C.....
C.....
      COEFFICIENTS OF THE COMPONENT C.
C.....
C.....
C.....
C.....

```


MST1 800
 MST1 810
 MST1 820
 MST1 830
 MST1 840
 MST1 850
 MST1 860
 MST1 870
 MST1 880
 MST1 890
 MST1 900
 MST1 910
 MST1 920
 MST1 930
 MST1 940
 MST1 950
 MST1 960
 MST1 970
 MST1 980
 MST1 990
 MST1 1000
 MST1 1010
 MST1 1020
 MST1 1030
 MST1 1040
 MST1 1050
 MST1 1060
 MST1 1070
 MST1 1080
 MST1 1090
 MST1 1100
 MST1 1110
 MST1 1120
 MST1 1130
 MST1 1140
 MST1 1150
 MST1 1160

```

C      = -R/REY
M3    = -6D0/REY
M2    = R*A*U-1D0/REY*(3D0/R+2D0*A**2*R)
M1    = 3D0*A*U+1D0/REY*(3D0/R**2-6D0*A**2)
M0    = R*A**3*U-1D0/REY*(A**4*R)

C      N2 = -R
N1    = -3D0
N0    = -A**2*R

C      IF (KSET.EQ.0) GO TO 2

C      CGM1(5) = M4*F44
CGM1(4) = M4*F43+M3*F33
CGM1(3) = M4*F42+M3*F32+M2*F22
CGM1(2) = M4*F41+M3*F31+M2*F21+M1*F11
CGM1(1) = M0

C      CQM2(3) = N2*F22
CQM2(2) = N2*F21+N1*F11
CQM2(1) = N0
RETURN

C      2  CGM1(5) = M4
CGM1(4) = M3
CGM1(3) = M2
CGM1(2) = M1
CGM1(1) = M0

C      CQM2(3) = N2
CQM2(2) = N1
CQM2(1) = N0

C      RETURN
END

C.....SUBROUTINE MSET2(X,N,MDIM,CFMAT,KSET).....
C
C      PURPOSE
C      MSET2 GENERATES THE ARRAYS REPRESENTING THE CENTRAL DIFFERENCE
C      APPROXIMATION OF THE VORTICITY TRANSPORT EQUATION IN TERMS OF
C      THE VELOCITY VECTOR POTENTIAL.
C
C      USAGE
  
```

MST2 10
 MST2 20
 MST2 30
 MST2 40
 MST2 50
 MST2 60
 MST2 70
 MST2 80
 MST2 90


```

100 CALL MSET2(X,N,MDIM,CFMAT,KSET)
110
120 DESCRIPTION OF PARAMETERS
130
140 X - THE NAME OF THE ARRAY BEING GENERATED. MUST BE DIMENSIONED
150 IN THE CALLING PROGRAM
160
170 N- THE ROW DIMENSION OF THE MATRIX X. MUST BE .GE. N.
180
190 MDIM - THE COLUMN DIMENSION OF THE MATRIX X. MUST BE .GE. N.
200
210 CFMAT - THE NAME OF A FUNCTION SUBPROGRAM WITH 4 PARAMETERS,
220 JSTA, K, CQM1 & CQM2. CFMAT MUST BE DECLARED
230 EXTERNAL IN THE CALLING PROGRAM.
240
250 THE FOLLOWING IS OUTPUT BY MSET2
260
270 X - THE N BY N MATRIX INTO WHICH THE COEFFICIENTS OF THE CENTRAL
280 DIFFERENCING ARE PUT.
290
300 OTHER ROUTINES NEEDED
310
320 FUNCTION SUBPROGRAM NAME PASSED IN THE CALLING PARAMETER 'CFMAT'
330 AND MSET1.
340
350 .....
360 SUBROUTINE MSET2 (X,N,MDIM,CFMAT,KSET)
370 REAL *8REY,R,DEL,DFLOAT,AMDA,ETA
380 COMPLEX *16X(MDIM,MDIM),CQM1(5),CQM2(3)
390 COMPLEX *16A,G
400 COMPLEX *16CFMAT
410 COMMON /COEFNT/ A,G,REY,DEL,AMDA
420
430 DEFINE THE SPACING OF THE INTERIOR MESH POINTS.
440
450 DEL = 1D0/DFLOAT(N+1)
460
470 INITIALIZE ALL ELEMENTS IN THE ARRAY TO ZERO.
480
490
500
510 DO 1 I=1,N
520
530 DO 1 J=1,N
540
550
560
570

```


1 X(I,J) = (000,000)

C
C
C
C
C
C
C

ESTABLISH THE CENTRAL DIFFERENCE APPROXIMATION AT EACH POINT IN THE MESH.

ETA = DEL
CALL MSET1 (ETA, CQM1, CQM2, KSET)
X(1,1) = CFMAT(1,1,CQM1,CQM2)
X(1,2) = CFMAT(1,2,CQM1,CQM2)
X(1,3) = CFMAT(1,3,CQM1,CQM2)
X(1,4) = CFMAT(1,4,CQM1,CQM2)
X(1,5) = CFMAT(1,5,CQM1,CQM2)

C

ETA = 2D0*DEL
CALL MSET1 (ETA, CQM1, CQM2, KSET)
X(2,1) = CFMAT(2,1,CQM1,CQM2)
X(2,2) = CFMAT(2,2,CQM1,CQM2)
X(2,3) = CFMAT(2,3,CQM1,CQM2)
X(2,4) = CFMAT(2,4,CQM1,CQM2)
X(2,5) = CFMAT(2,5,CQM1,CQM2)

C
C

IL = N-2

DO 2 I=3,IL

K = I-3

ETA = DEL*DFLOAT(I)

CALL MSET1 (ETA, CQM1, CQM2, KSET)

C

DC 2 J=1, 5
2 X(I,K+J) = CFMAT(3,J,CQM1,CQM2)

C
C

ETA = 1D0-2D0*DEL
CALL MSET1 (ETA, CQM1, CQM2, KSET)
X(N-1,N-3) = CFMAT(4,1,CQM1,CQM2)
X(N-1,N-2) = CFMAT(4,2,CQM1,CQM2)
X(N-1,N-1) = CFMAT(4,3,CQM1,CQM2)
X(N-1,N) = CFMAT(4,4,CQM1,CQM2)

C

ETA = 1D0-DEL
CALL MSET1 (ETA, CQM1, CQM2, KSET)
X(N,N-3) = CFMAT(5,1,CQM1,CQM2)
X(N,N-2) = CFMAT(5,2,CQM1,CQM2)
X(N,N-1) = CFMAT(5,3,CQM1,CQM2)
X(N,N) = CFMAT(5,4,CQM1,CQM2)

MST2 580
MST2 590
MST2 600
MST2 610
MST2 620
MST2 630
MST2 640
MST2 650
MST2 660
MST2 670
MST2 680
MST2 690
MST2 700
MST2 710
MST2 720
MST2 730
MST2 740
MST2 750
MST2 760
MST2 770
MST2 780
MST2 790
MST2 800
MST2 810
MST2 820
MST2 830
MST2 840
MST2 850
MST2 860
MST2 870
MST2 880
MST2 890
MST2 900
MST2 910
MST2 920
MST2 930
MST2 940
MST2 950
MST2 960
MST2 970
MST2 980
MST2 990
MST2 1000
MST2 1010
MST2 1020
MST2 1030
MST2 1040
MST2 1050

RETURN
END

MST21060
MST21070

.....FUNCTION CQM1E1(JSTA,K,CQM1,CQM2).....
(POLAR COORDINATES)

CQM1 10
CQM1 20
CQM1 30
CQM1 40
CQM1 50
CQM1 60
CQM1 70
CQM1 80
CQM1 90
CQM1 100
CQM1 110
CQM1 120
CQM1 130
CQM1 140
CQM1 150
CQM1 160
CQM1 170
CQM1 180
CQM1 190
CQM1 200
CQM1 210
CQM1 220
CQM1 230
CQM1 240
CQM1 250
CQM1 260
CQM1 270
CQM1 280
CQM1 290
CQM1 300
CQM1 310
CQM1 320
CQM1 330
CQM1 340
CQM1 350
CQM1 360
CQM1 370
CQM1 380
CQM1 390
CQM1 400
CQM1 410
CQM1 420
CQM1 430
CQM1 440

PURPOSE

RETURNS THE VALUES FOR THE COEFFICIENTS IN THE ARRAYS
REPRESENTING THE CENTRAL DIFFERENCE APPROXIMATION OF THE
VORTICITY TRANSPORT EQUATION USING THE COEFFICIENTS COMPUTED
BY SUBROUTINE MSET1.

DESCRIPTION OF PARAMETERS

JSTA - INDICATES WHICH DIFFERENCE EQUATION SET WILL BE USED.

- JSTA=1 - CONSISTENT 3RD ORDER TRUNCATION ERROR FINITE
DIFFERENCE EQUATIONS FOR R=DEL WILL BE USED.
- JSTA=2 - SAME AS ABOVE BUT R=2D0*DEL
- JSTA=3 - CENTRAL DIFFERENCE EQUATIONS WITH CONSISTENT 2ND
ORDER TRUNCATION ERROR WILL BE USED.
- JSTA=4 - SAME AS JSTA=3 BUT FOR R=1C0-2D0*DEL.
- JSTA=5 - SAME AS ABOVE BUT FOR R=1D0-DEL.

K - INDICATES THE ABSOLUTE POSITION OF THE POINT IN EACH ROW
OF THE FINITE DIFFERENCE MESH. IF THE FIRST NON-ZERO ENTRY
IN ROW J IS ELEMENT(J,3), THEN K=1 DENOTES ELEMENT(J,3),
K=2 DENOTES ELEMENT(J,4), ETC.

CQM1,CQM2 - THE COEFFICIENT ARRAYS FOR THE FINITE DIFFERENCE
APPROXIMATION OF THE FUNCTION Q. CQM1 CONTAINS THE
COEFFICIENTS FOR THE NON-GAMMA TERMS, WHILE CQM2 CONTAINS
THE COEFFICIENTS OF THE GAMMA TERMS. BOTH ARRAYS MUST BE
DIMENSIONED COMPLEX*16.

EXAMPLE OF THE CALLING ARGUMENT:

CQ M(1,2) E1(JSTA,K,CQM1,CQM2)

CQ - Q COMPONENT OF THE VELOCITY VECTOR POTENTIAL.

M(1,2) - 1 REFERS TO TERMS NOT CONTAINING GAMMA AS A
FACTOR.

2 REFERS TO TERMS CONTAINING GAMMA AS A FACTOR.

CC

450 CQM1
 460 CQM1
 470 CQM1
 480 CQM1
 490 CQM1
 500 CQM1
 510 CQM1
 520 CQM1
 530 CQM1
 540 CQM1
 550 CQM1
 560 CQM1
 570 CQM1
 580 CQM1
 590 CQM1
 600 CQM1
 610 CQM1
 620 CQM1
 630 CQM1
 640 CQM1
 650 CQM1
 660 CQM1
 670 CQM1
 680 CQM1
 690 CQM1
 700 CQM1
 710 CQM1
 720 CQM1
 730 CQM1
 740 CQM1
 750 CQM1
 760 CQM1
 770 CQM1
 780 CQM1
 790 CQM1
 800 CQM1
 810 CQM1
 820 CQM1
 830 CQM1
 840 CQM1
 850 CQM1
 860 CQM1
 870 CQM1
 880 CQM1
 890 CQM1
 900 CQM1
 910 CQM1
 920 CQM1

E1 - REFERS TO THE LINEAR COMBINATION OF THE FIRST AND
 THIRD EQUATIONS RESULTING FROM EXPRESSION OF THE
 VORTICITY TRANSPORT EQUATION IN TERMS OF THE VELOCITY
 VECTOR POTENTIAL.

USAGE

CQM1 MUST BE DECLARED COMPLEX*16 IN THE CALLING PROGRAM.

OTHER ROUTINES REQUIRED

NONE

FUNCTION CQM1 (JSTA, K, CQM1, CQM2)
 IMPLICIT COMPLEX*16(A-H,O-Z)
 COMMON / COEFNT/ A, G, REY, DEL, AMDA
 COMPLEX *16 CQM1(5), CQM2(3)
 REAL *8 REY, R, DEL

GC TO (1,7,13,19,24), JSTA

FINITE DIFFERENCE EQUATIONS AT ETA=DEL (NON GAMMA).

1 GO TO (2,3,4,5,6), K
 2 CQM1 = -0.911564626D0*CQM1(5)/DEL**4+1.825165563D0*CQM1(4)/DEL**CQM1(4)
 13-35D0*CQM1(3)/(60D0*DEL**2)-2D0*CQM1(2)/(3D0*DEL)+CQM1(1)
 GO TO 29
 3 CQM1 = 2.750242955D0*CQM1(5)/DEL**4-3.250331126D0*CQM1(4)/DEL**3
 1+8D0*CQM1(3)/(15D0*DEL**2)+2D0*CQM1(2)/(3D0*DEL)
 GO TO 29
 4 CQM1 = -3.043731779D0*CQM1(5)/DEL**4+1.660927152D0*CQM1(4)/DEL**CQM1(4)
 13+CQM1(3)/(20D0*DEL**2)
 GO TO 29
 5 CQM1 = 1.42468416D0*CQM1(5)/DEL**4-.235761589D0*CQM1(4)/DEL**3
 GO TO 29
 6 CQM1 = -0.219630709D0*CQM1(5)/DEL**4
 GO TO 29

FINITE DIFFERENCE EQUATIONS AT ETA=2D0*DEL (NON GAMMA).

7 GO TO (8,9,10,11,12), K
 8 CQM1 = -3.10340136D0*CQM1(5)/DEL**4+.868874172D0*CQM1(4)/DEL**3+CQM1(4)
 11D0*CQM1(3)/(12D0*DEL**2)-4D0*CQM1(2)/(3D0*DEL)
 GO TO 29
 9 CQM1 = 6.903012634D0*CQM1(5)/DEL**4-.937748345D0*CQM1(4)/DEL**3-CQM1(4)

CCCCCCCCCCCC

CCCC

CC


```

123D0*CQM1(3)/(15D0*DEL**2)+4D0*CQM1(2)/(3D0*DEL)+CQM1(1)
GO TO 29
10 CQM1E1 = -5.342274053D0*CQM1(5)/DEL**4-.254304636D0*CQM1(4)/DEL**3
1+19D0*CQM1(3)/(20D0*DEL**2)
GO TO 29
11 CQM1E1 = 1.666083577D0*CQM1(5)/DEL**4+.3251788C8D C*CQM1(4)/DEL**3
GO TO 29
12 CQM1E1 = -0.123420797D0*CQM1(5)/DEL**4
GO TO 29
C
C
C
CENTRAL DIFFERENCE APPROXIMATION FOR COMPONENT Q (NON GAMMA).
13 GO TO (14,15,16,17,18), K
14 CQM1E1 = CQM1(5)/DEL**4-CQM1(4)/(2D0*DEL**3)
GO TO 29
15 CQM1E1 = -4D0*CQM1(5)/DEL**4+CQM1(4)/DEL**3+CQM1(3)/DEL**2-CQM1(2)
1/(2D0*DEL)
GO TO 29
16 CQM1E1 = 6D0*CQM1(5)/DEL**4-2D0*CQM1(3)/DEL**2+CQM1(1)
GO TO 29
17 CQM1E1 = -4D0*CQM1(5)/DEL**4-CQM1(4)/DEL**3+CQM1(3)/DEL**2+CQM1(2)
1/(2D0*DEL)
GO TO 29
18 CQM1E1 = CQM1(5)/DEL**4+CQM1(4)/(2D0*DEL**3)
GO TO 29
C
C
C
FINITE DIFFERENCE EQUATIONS AT ETA=1D0-2D0*DEL (NCN-GAMMA).
19 GO TO (20,21,22,23), K
20 CQM1E1 = CQM1(5)/DEL**4-0.5D0*CQM1(4)/DEL**3
GO TO 29
21 CQM1E1 = -4D0*CQM1(5)/DEL**4+CQM1(4)/DEL**3+CQM1(3)/DEL**2-CQM1(2)
1/(2D0*DEL)
GO TO 29
22 CQM1E1 = 6D0*CQM1(5)/DEL**4-2D0*CQM1(3)/DEL**2+CQM1(1)
GO TO 29
23 CQM1E1 = -4D0*CQM1(5)/DEL**4-CQM1(4)/DEL**3+CQM1(3)/DEL**2+CQM1(2)
1/(2D0*DEL)
GO TO 29
C
C
C
FINITE DIFFERENCE EQUATIONS AT ETA=1D0-DEL (NON GAMMA).
24 GO TO (25,26,27,28), K
25 CQM1E1 = -0.25D0*CQM1(5)/DEL**4
GO TO 29
26 CQM1E1 = 8D0*CQM1(5)/(3D0*DEL**4)-CQM1(4)/(3D0*DEL**3)
GO TO 29
27 CQM1E1 = -9D0*CQM1(5)/DEL**4+CQM1(3)/DEL**2-CQM1(2)/(2D0*DEL)
GO TO 29

```



```

GO TO 29
28 CQM2E1 = 1600*CQM1(5) / DEL**4 + 300*CQM1(4) / DEL**3 - 200*CQM1(3) / DEL**2
1+CQM1(1)
25 RETURN
C
ENTRY CQM2E1(JSTA,K,CQM1,CQM2)
C
GO TO (30,35,40,45,5C), JSTA
C
FINITE DIFFERENCE EQUATIONS AT ETA=DEL ( GAMMA ).
C
30 GO TO (31,32,33,34,34), K
31 CQM2E1 = -3500*CQM2(3) / (6000*DEL**2) - 200*CQM2(2) / (300*DEL) + CQM2(1)
GO TO 54
32 CQM2E1 = 800*CQM2(3) / (1500*DEL**2) + 200*CQM2(2) / (300*DEL)
GO TO 54
33 CQM2E1 = CQM2(3) / (2000*DEL**2)
GO TO 54
34 CQM2E1 = (000,000)
C
FINITE DIFFERENCE EQUATIONS AT ETA=200*DEL ( GAMMA )
C
35 GO TO (36,37,38,39,35), K
36 CQM2E1 = 1100*CQM2(3) / (1200*DEL**2) - 400*CQM2(2) / (300*DEL)
GO TO 54
37 CQM2E1 = -2800*CQM2(3) / (1500*DEL**2) + 400*CQM2(2) / (300*DEL) + CQM2(1)
GO TO 54
38 CQM2E1 = 1900*CQM2(3) / (2000*DEL**2)
GO TO 54
39 CQM2E1 = (000,000)
C
CENTRAL DIFFERENCE EQUATIONS FOR THE COMPONENT Q ( GAMMA ).
C
40 GO TO (41,42,43,44,41), K
41 CQM2E1 = (000,000)
GO TO 54
42 CQM2E1 = CQM2(3) / DEL**2 - CQM2(2) / (200*DEL)
GO TO 54
43 CQM2E1 = -200*CQM2(3) / DEL**2 + CQM2(1)
GO TO 54
44 CQM2E1 = CQM2(3) / DEL**2 + CQM2(2) / (200*DEL)
GO TO 54
C
FINITE DIFFERENCE EQUATIONS AT ETA=100-200*DEL ( GAMMA ).
C
45 GC TO (49,46,47,48), K
46 CQM2E1 = CQM2(3) / DEL**2 - CQM2(2) / (200*DEL)
CQM11410
CQM11420
CQM11430
CQM11440
CQM11450
CQM11460
CQM11470
CQM11480
CQM11490
CQM11500
CQM11510
CQM11520
CQM11530
CQM11540
CQM11550
CQM11560
CQM11570
CQM11580
CQM11590
CQM11600
CQM11610
CQM11620
CQM11630
CQM11640
CQM11650
CQM11660
CQM11670
CQM11680
CQM11690
CQM11700
CQM11710
CQM11720
CQM11730
CQM11740
CQM11750
CQM11760
CQM11770
CQM11780
CQM11790
CQM11800
CQM11810
CQM11820
CQM11820
CQM11840
CQM11850
CQM11860
CQM11870
CQM11880

```



```

47  GC TO 54      -2D0*CQM2(3)/DEL**2+CQM2(1)
    CQM2E1 =
48  GC TO 54      CQM2(3)/DEL**2+CQM2(2)/(2D0*DEL)
    CQM2E1 =
49  GC TO 54      (0D0,0D0)
    CQM2E1 =
    GO TO 54
C
C
C      FINITE DIFFERENCE EQUATIONS AT ETA=1D0-DEL ( GAMMA ).
50  GC TO (53,53,51,52) K
51  CQM2E1 = CQM2(3)/DEL**2-CQM2(2)/(2D0*DEL)
    GO TO 54
52  CQM2E1 = -2D0*CQM2(3)/DEL**2+CQM2(1)
    GO TO 54
53  CQM2E1 = (0D0,0D0)
54  RETURN
    END

```

```

CQMI1890
CQMI1900
CQMI1910
CQMI1920
CQMI1930
CQMI1940
CQMI1950
CQMI1960
CQMI1970
CQMI1980
CQMI1990
CQMI2000
CQMI2010
CQMI2020
CQMI2030
CQMI2040
CQMI2050
CQMI2060

```

```

..... SUBROUTINE CDMTIN(N,A,NDIM,IERR).....
PURPOSE
  INVERT A COMPLEX*16 MATRIX
USAGE
  CALL CDMTIN(N,A,NCIM,DETERM)
DESCRIPTION OF PARAMETERS
  N      - ORDER OF COMPLEX*16 MATRIX TO BE INVERTED
          (INTEGER) MAXIMUM 'N' IS 100
  A      - COMPLEX*16 INPUT MATRIX (DESTROYED). THE
          INVERSE OF 'A' IS RETURNED IN ITS PLACE
  NDIM   - THE SIZE TO WHICH 'A' IS DIMENSIONED
          (ROW DIMENSION OF 'A' ACTUALLY APPEARING
          IN THE DIMENSION STATEMENT OF USER'S
          CALLING PROGRAM)
  IERR   - ERROR PARAMETER RETURNED BY CDMTIN. IERR = 0 INDICATES
          NORMAL INVERSION. IERR = 9999 INDICATES SINGULAR MATRIX.
REMARKS

```

```

CDMT 10
CDMT 20
CDMT 30
CDMT 40
CDMT 50
CDMT 60
CDMT 70
CDMT 80
CDMT 90
CDMT 100
CDMT 110
CDMT 120
CDMT 130
CDMT 140
CDMT 150
CDMT 160
CDMT 170
CDMT 180
CDMT 190
CDMT 200
CDMT 210
CDMT 220
CDMT 230
CDMT 240
CDMT 250
CDMT 260
CDMT 270
CDMT 280

```



```

C      MATRIX 'A' MUST BE A COMPLEX*16 GENERAL MATRIX          CDMT 290
C      IF MATRIX 'A' IS SINGULAR THAT MESSAGE IS PRINTED      CDMT 300
C      'N' MUST BE .LE. NDIM                                  CDMT 310
C      SUBROUTINES AND FUNCTIONS REQUIRED                       CDMT 320
C      ONLY BUILT-IN FORTRAN FUNCTIONS                       CDMT 330
C      METHOD                                                 CDMT 340
C      GAUSSIAN ELIMINATION WITH COLUMN PIVOTING IS USED.   CDMT 350
C      .....                                                CDMT 360
C      SUBROUTINE CDMTIN (N,A,NDIM,IERR)                     CDMT 370
C      IMPLICIT REAL*8(A-H,O-Z)                              CDMT 380
C      INTEGER *4 IPIVOT(100),INDEX(100,2),IERR            CDMT 390
C      REAL *8 TEMP,ALPHA(100)                               CDMT 400
C      COMPLEX *16 A(NDIM,NDIM),PIVOT(100),AMAX,T,SWAP,U    CDMT 410
C      INITIALIZATION                                        CDMT 420
C      DO 2 J=1,N                                           CDMT 430
C      IERR = 0                                             CDMT 440
C      ALPHA(J) = 0D0                                       CDMT 450
C      DO 1 I=1,N                                           CDMT 460
C      ALPHA(J) = ALPHA(J)+A(J,I)*DCONJG(A(J,I))           CDMT 470
C      ALPHA(J) = DSQRT(ALPHA(J))                          CDMT 480
C      IPIVOT(J) = 0                                         CDMT 490
C      DO 16 I=1,N                                           CDMT 500
C      SEARCH FOR PIVOT ELEMENT                             CDMT 510
C      AMAX = (0D0,0D0)                                       CDMT 520
C      DO 7 J=1,N                                           CDMT 530
C      IF ( IPIVOT(J)-1) 3,7,3                               CDMT 540
C      DO 6 K=1,N                                           CDMT 550
C      IF ( IPIVOT(K)-1) 4,6,21                             CDMT 560
C      TEMP = AMAX*DCONJG(AMAX)-A(J,K)*DCONJG(A(J,K))       CDMT 570
C      IF (TEMP) 5,5,6                                       CDMT 580
C      IROW = J                                             CDMT 590
C      I=IPIVOT(J)                                          CDMT 600
C      I=IPIVOT(I)                                          CDMT 610
C      I=IPIVOT(I)                                          CDMT 620
C      I=IPIVOT(I)                                          CDMT 630
C      I=IPIVOT(I)                                          CDMT 640
C      I=IPIVOT(I)                                          CDMT 650
C      I=IPIVOT(I)                                          CDMT 660
C      I=IPIVOT(I)                                          CDMT 670
C      I=IPIVOT(I)                                          CDMT 680
C      I=IPIVOT(I)                                          CDMT 690
C      I=IPIVOT(I)                                          CDMT 700
C      I=IPIVOT(I)                                          CDMT 710
C      I=IPIVOT(I)                                          CDMT 720
C      I=IPIVOT(I)                                          CDMT 730
C      I=IPIVOT(I)                                          CDMT 740
C      I=IPIVOT(I)                                          CDMT 750
C      I=IPIVOT(I)                                          CDMT 760

```



```

ICOLUM = K
AMAX = A(J,K)
6 CONTINUE
C
7 CCNTINUE
C
IPIVOT(ICOLUM) = IPIVOT(ICOLUM)+1
INTERCHANGE ROWS TO PUT PIVOT ELEMENT ON DIAGONAL
C
IF (IROW-ICCLUM) 8,10,8
8 CONTINUE
C
DC 9 L=1,N
SWAP = A(IROW,L)
A(IROW,L) = A(ICOLUM,L)
9 A(ICOLUM,L) = SWAP
C
SWAP = ALPHA(IROW)
ALPHA(IROW) = ALPHA(ICOLUM)
ALPHA(ICOLUM) = SWAP
10 INDEX(I,1) = IROW
INDEX(I,2) = ICOLUM
PIVOT(I) = A(ICOLUM, ICOLUM)
U = PIVOT(I)
TEMP = PIVOT(I)*DCONJG(PIVOT(I))
IF (TEMP) 11,20,11
C
DIVIDE PIVOT ROW BY PIVOT ELEMENT
C
11 A(ICOLUM,ICOLUM) = (10,000)
C
DC 12 L=1,N
U = PIVOT(I)
12 A(ICOLUM,L) = A(ICOLUM,L)/U
C
REDUCE NON-PIVOT ROWS
C
13 IF (L1-ICOLUM) 13,15,13
T = A(L1,ICOLUM)
A(L1,ICOLUM) = (000,000)
C
14 DC 14 L=1,N
U = A(ICOLUM,L)
A(L1,L) = A(L1,L)-U*T
C

```

```

CDMT 770
CDMT 780
CDMT 790
CDMT 800
CDMT 810
CDMT 820
CDMT 830
CDMT 840
CDMT 850
CDMT 860
CDMT 870
CDMT 880
CDMT 890
CDMT 900
CDMT 910
CDMT 920
CDMT 930
CDMT 940
CDMT 950
CDMT 960
CDMT 970
CDMT 980
CDMT 990
CDMT 1000
CDMT 1010
CDMT 1020
CDMT 1030
CDMT 1040
CDMT 1050
CDMT 1060
CDMT 1070
CDMT 1080
CDMT 1090
CDMT 1100
CDMT 1110
CDMT 1120
CDMT 1130
CDMT 1140
CDMT 1150
CDMT 1160
CDMT 1170
CDMT 1180
CDMT 1190
CDMT 1200
CDMT 1210
CDMT 1220
CDMT 1230
CDMT 1240

```


C	15	CONTINUE	CDMT11250
C			CDMT11260
C	16	CONTINUE	CDMT11270
C			CDMT11280
C		INTERCHANGE COLUMNS	CDMT11290
C			CDMT11300
C			CDMT11310
C			CDMT11320
C			CDMT11330
	19	I=1,N	CDMT11340
		L = N+1-1	CDMT11350
		IF (INDEX(L,1)-INDEX(L,2))	CDMT11360
	17	JROW = INDEX(L,1)	CDMT11370
		JCOLUMN = INDEX(L,2)	CDMT11380
C			CDMT11390
	18	DO 18 K=1,N	CDMT11400
		SWAP = A(K,JROW)	CDMT11410
		A(K,JROW) = A(K,JCOLUMN)	CDMT11420
		A(K,JCOLUMN) = SWAP	CDMT11430
C	18	CONTINUE	CDMT11440
C	19	CONTINUE	CDMT11450
			CDMT11460
			CDMT11470
		RETURN	CDMT11480
	20	WRITE (6,22)	CDMT11490
		IERR = 9999	CDMT11500
	21	RETURN	CDMT11510
C	22	FORMAT (20H MATRIX IS SINGULAR)	CDMT11520
		END	CDMT11530
C			CDMT11540
C	 SUBROUTINE MULM(X1,X2,N,MDIM,TEMPV).....	MULM 10
C		PURPOSE	MULM 20
C		PERFORMS THE MATRIX MULTIPLICATION OF A SQUARE MATRIX BY A	MULM 30
C		SQUARE MATRIX. THE RESULT IS RETURNED IN MATRIX X1.	MULM 40
C			MULM 50
C			MULM 60
C			MULM 70
C		USAGE	MULM 80
C		CALL MULM(X1,X2,N,MDIM,TEMPV)	MULM 90
C		DESCRIPTION OF PARAMETERS	MULM 100
C			MULM 110
C			MULM 120
C			MULM 130
C		X1 - THE MULTIPLYING MATRIX ON INPUT AND THE RESULTANT PRODUCT	MULM 140
C		ON OUTPUT.	MULM 150
C			MULM 160


```

C      X2 -- THE MULTIPLIED MATRIX.
C      N -- THE ORDER OF X1 AND X2.
C      MDIM -- THE DIMENSION OF X1 AND X2 FROM THE CALLING PROGRAM.
C      TEMPV -- A WORKING VECTOR. MUST BE DIMENSIONED MDIM.
C      OTHER ROUTINES REQUIRED
C      NONE
C      .....
C      SUBROUTINE MULM (X1,X2,N,MDIM,TEMPV)
C      COMPLEX *16X1(MDIM,MDIM),X2(MDIM,MDIM),TEMPV(MDIM),TEMP
C      .....
C      STORE ROW I OF X1 IN TEMPV.
C      DO 4 I=1,N
C      DO 1 J=1,N
C      1  TEMPV(J) = X1(I,J)
C      .....
C      MULTIPLY COLUMN J OF X2 BY ROW I OF X1 AND STORE IN X1(I,J).
C      DO 3 J=1,N
C      TEMP = (0D0,0D0)
C      DO 2 K=1,N
C      2  TEMP = TEMP+TEMPV(K)*X2(K,J)
C      3  X1(I,J) = TEMP
C      4  CONTINUE
C      RETURN
C      END
C      .....
MULM 170
MULM 180
MULM 190
MULM 200
MULM 210
MULM 220
MULM 230
MULM 240
MULM 250
MULM 260
MULM 270
MULM 280
MULM 290
MULM 300
MULM 310
MULM 320
MULM 330
MULM 340
MULM 350
MULM 360
MULM 370
MULM 380
MULM 390
MULM 400
MULM 410
MULM 420
MULM 430
MULM 440
MULM 450
MULM 460
MULM 470
MULM 480
MULM 490
MULM 500
MULM 510
MULM 520
MULM 530
MULM 540
MULM 550
MULM 560
MULM 570
MULM 580
MULM 590
MULM 600
MULM 610
MULM 620

```


10	DSPL SUBROUTINE DSPLIT(N,MDIM,A,AR,AI).....
20	DSPL	
30	DSPL	
40	DSPL	
50	DSPL	PURPOSE
60	DSPL	DSPLIT TAKES A MATRIX OF COMPLEX*16 NUMBERS AND
70	DSPL	SPLITS IT INTO TWO MATRICES, ONE CONTAINING THE REAL
80	DSPL	PART OF THE ORIGINAL MATRIX, AND ONE CONTAINING THE
90	DSPL	IMAGINARY PART.
100	DSPL	USAGE
110	DSPL	CALL DSPLIT(N,MDIM,A,AREAL,AIMAG)
120	DSPL	DESCRIPTION OF PARAMETERS
130	DSPL	N - THE SIZE OF THE MATRIX A, AN N BY N SQUARE
140	DSPL	MATRIX.
150	DSPL	MDIM - THE COLUMN DIMENSION OF MATRIX A
160	DSPL	A - THE INPUT MATRIX. MUST BE DIMENSIONED MDIM BY
170	DSPL	AT LEAST N IN THE CALLING PROGRAM (COMPLEX*16)
180	DSPL	AREAL,AIMAG - THE OUTPUT MATRICES CONTAINING THE
190	DSPL	REAL AND IMAGINARY PARTS, RESPECTIVELY, OF
200	DSPL	MATRIX A. MUST BE DIMENSIONED (MDIM,MDIM) IN THE
210	DSPL	CALLING PROGRAM.
220	DSPL	NOTES....
230	DSPL	MATRIX A AND MATRIX AREAL MAY OVERLAP IF THEY ARE
240	DSPL	DIMENSIONED IN THE CALLING PROGRAM AS FOLLOWS....
250	DSPL	COMPLEX*16 A(MDIM,MDIM)
260	DSPL	REAL*8 AREAL(MDIM,MDIM),AIMAG(MDIM,MDIM)
270	DSPL	EQUIVALENCE(A(1,1),AREAL(1,1))
280	DSPL	OTHER ROUTINES NEEDED
290	DSPL	NONE
300	DSPL
310	DSPL	SUBROUTINE DSPLIT (N,MDIM,A,AR,AI)
320	DSPL	REAL *8A(2,MDIM,MDIM),AR(MDIM,MDIM),AI(MDIM,MDIM)
330	DSPL
340	DSPL	
350	DSPL	
360	DSPL	
370	DSPL	
380	DSPL	
390	DSPL	
400	DSPL	
410	DSPL	
420	DSPL	
430	DSPL	
440	DSPL	
450	DSPL	
460	DSPL	
470	DSPL	
480	DSPL	

DSPL 490
 DSPL 500
 DSPL 510
 DSPL 520
 DSPL 530
 DSPL 540
 DSPL 550
 DSPL 560
 DSPL 570
 DSPL 580

```

C DO 1 J=1,N
C DO 1 I=1,N
  AR(I,J) = A(1,I,J)
  1 AI(I,J) = A(2,I,J)
C
C RETURN
C END
  
```

EBAC0010
 EBAC0020
 EBAC0030
 EBAC0040
 EBAC0050
 EBAC0060
 EBAC0070
 EBAC0080
 EBAC0090
 EBAC0100
 EBAC0110
 EBAC0120
 EBAC0130
 EBAC0140
 EBAC0150
 EBAC0160
 EBAC0170
 EBAC0180
 EBAC0190
 EBAC0200
 EBAC0210
 EBAC0220
 EBAC0230
 EBAC0240
 EBAC0250
 EBAC0260
 EBAC0270
 EBAC0280
 EBAC0290
 EBAC0300
 EBAC0310
 EBAC0320
 EBAC0330
 EBAC0340
 EBAC0350
 EBAC0360

```

SUBROUTINE EBALAC (AR, AI, N, IA, K, L, D)
-----LIBRARY 1-----
C
C FUNCTION
C USAGE
C PARAMETERS      AR
                  AI
C
C N
C IA
C K
C L
C
C D
C
C PRECISION
C LANGUAGE
C
C LATEST REVISION      MARCH 9, 1977
C
C SUBROUTINE EBALAC (AR, AI, N, IA, K, L, D)
C DIMENSION      AR(IA, 1), AI(IA, 1), D(N)
  
```

```

- BALANCES A COMPLEX GENERAL MATRIX AND ISOLATES
  EIGENVALUES WHENEVER POSSIBLE.
- CALL EBALAC (AR, AI, N, IA, K, L, D)
- INPUT/OUTPUT MATRICES OF DIMENSION N BY N.
  ON INPUT, AR AND AI CONTAIN THE REAL
  AND IMAGINARY PARTS, RESPECTIVELY, OF
  THE COMPLEX MATRIX OF ORDER N TO BE
  BALANCED. ON OUTPUT, AR AND AI CONTAIN THE
  REAL AND IMAGINARY PARTS OF THE
  TRANSFORMED MATRIX.
- INPUT VARIABLE CONTAINING THE ORDER
  OF THE MATRIX A = (AR, AI) TO BE BALANCED.
- INPUT VARIABLE CONTAINING THE ROW DIMENSION OF
  AR AND AI IN THE CALLING PROGRAM.
- OUTPUT INTEGERS CONTAINING THE BOUNDARY
  INDICES FOR THE BALANCED MATRIX A = (AR, AI)
  SUCH THAT
  AR(I, J) = 0. AND AI(I, J) = 0. IF
  (1) J IS GREATER THAN J AND
  (2) J = 1, ..., K-1 OR
  I = L+1, ..., N
- OUTPUT VECTOR OF LENGTH N CONTAINING
  INFORMATION DETERMINING THE PERMUTATIONS
  USED AND THE SCALING FACTORS.
- SINGLE/DOUBLE
- FORTRAN
  
```


EBAC0370
 EBAC0380
 EBAC0390
 EBAC0400
 EBAC0410
 EBAC0420
 EBAC0430
 EBAC0440
 EBAC0450
 EBAC0480
 EBAC0490
 EBAC0500
 EBAC0510
 EBAC0520
 EBAC0530
 EBAC0540
 EBAC0550
 EBAC0560
 EBAC0570
 EBAC0580
 EBAC0590
 EBAC0600
 EBAC0610
 EBAC0620
 EBAC0630
 EBAC0640
 EBAC0650
 EBAC0660
 EBAC0670
 EBAC0680
 EBAC0690
 EBAC0700
 EBAC0710
 EBAC0720
 EBAC0730
 EBAC0740
 EBAC0750
 EBAC0760
 EBAC0770
 EBAC0780
 EBAC0790
 EBAC0800
 EBAC0810
 EBAC0820
 EBAC0830
 EBAC0840
 EBAC0850
 EBAC0860

LOGICAL NOCONV
 RADIX IS A MACHINE DEPENDENT
 PARAMETER SPECIFYING THE BASE OF
 THE MACHINE FLOATING POINT REPRESENTATION

DOUBLE PRECISION AR, AI, D, RADIX, ZERO, ONE, PT95, B2, F, C, G, R, S
 DOUBLE PRECISION RRADIX, RB2
 DATA RADIX/16.0D0/
 DATA ZERO, ONE, PT95/0.0D0, 1.0D0, 0.95D0/

B2 = RADIX * RADIX
 RRADIX = ONE / RADIX
 RB2 = RRADIX * RRADIX
 K = 1
 L = N
 GO TO 30

5 D(M) = J
 IF (J.EQ.M) GO TO 20
 DO 10 I = 1, L
 F = AR(I, J)
 AR(I, J) = AR(I, M)
 AR(I, M) = F
 F = AI(I, J)
 AI(I, J) = AI(I, M)
 AI(I, M) = F

10 CONTINUE
 DO 15 I = K, N
 F = AR(J, I)
 AR(J, I) = AR(M, I)
 AR(M, I) = F
 F = AI(J, I)
 AI(J, I) = AI(M, I)
 AI(M, I) = F

15 CONTINUE
 20 GO TO (25, 45), IEXC

25 IF (L.EQ. 1) GO TO 115
 L = L - 1

30 L1 = L + 1
 DO 40 JJ = 1, L
 J = L1 - JJ
 DO 35 I = 1, L
 IF (I.EQ. J) GC TO 35
 IF (AR(J, I) .NE. ZERO .OR. AI(J, I) .NE. ZERC) GO TO 40

35 CONTINUE

IN-LINE PROCEDURE FOR ROW AND COLUMN EXCHANGE

SEARCH FOR ROWS ISOLATING AN EIGENVALUE AND PUSH THEM DOWN
 DO J=L, 1, -1

EBAC0870
 EBAC0880
 EBAC0890
 EBAC0900
 EBAC0910
 EBAC0920
 EBAC0930
 EBAC0940
 EBAC0950
 EBAC0960
 EBAC0970
 EBAC0980
 EBAC0990
 EBAC1000
 EBAC1010
 EBAC1020
 EBAC1030
 EBAC1040
 EBAC1050
 EBAC1060
 EBAC1070
 EBAC1080
 EBAC1090
 EBAC1100
 EBAC1110
 EBAC1120
 EBAC1130
 EBAC1140
 EBAC1150
 EBAC1160
 EBAC1180
 EBAC1190
 EBAC1210
 EBAC1220
 EBAC1230
 EBAC1240
 EBAC1250
 EBAC1260
 EBAC1270
 EBAC1280
 EBAC1290
 EBAC1300
 EBAC1310
 EBAC1320
 EBAC1330
 EBAC1340
 EBAC1350
 EBAC1360

```

M = L
IEXC = 1
GO TO 5
40 CONTINUE
GO TO 50

45 K = K+1
50 DO 60 J = K,L
   DO 55 I = K,L
     IF (I.EQ.J) GC TO 55
     IF (AR(I,J).NE.ZERO) GO TO 60
   CONTINUE
   M = K
   IEXC = 2
   GO TO 5

60 CONTINUE

DO 65 I = K,L
D(I) = ONE
65 CONTINUE

70 NOCONV = .FALSE.
DO 110 I = K,L
  C = ZERO
  R = ZERO
  DO 75 J = K,L
    IF (J.EQ.I) GO TO 75
    C = C+DABS(AR(J,I))+DABS(AI(J,I))
    R = R+DABS(AR(I,J))+DABS(AI(I,J))
  CONTINUE
  G = R*RRADIX
  F = ONE
  S = C+R
  IF (C.GE.G) GO TO 85
  F = F*RADIX
  C = C*B2
  GO TO 80
  G = R*RRADIX
  IF (C.LT.G) GO TO 95
  F = F*RRADIX
  C = C*B2
  GO TO 90

95 IF ((C+R)/F.GE.PT95*S) GO TO 110
   G = ONE/F
   D(I) = D(I)*F

      SEARCH FOR COLUMNS ISOLATING AN
      EIGENVALUE AND PUSH THEM LEFT

      BALANCE THE SUBMATRIX IN ROWS
      K TO L

      ITERATIVE LOOP FOR NORM REDUCTION

      BALANCE
      GO TO 110
  
```



```

NOCNV = .TRUE.
DO 100 J = K, N
  AR(I, J) = AR(I, J)*G
  AI(I, J) = AI(I, J)*G
CONTINUE
DO 105 J = 1, L
  AR(J, I) = AR(J, I)*F
  AI(J, I) = AI(J, I)*F
CONTINUE
100 CONTINUE
105 CONTINUE
110 IF (NOCNV) GO TO 70
115 RETURN
END

```

-----D-----LIBRARY I-----

```

EBAC1370
EBAC1380
EBAC1390
EBAC1400
EBAC1410
EBAC1420
EBAC1430
EBAC1440
EBAC1450
EBAC1460
EBAC1470
EBAC1480
EBAC1490
EHEC0010
EHEC0020
EHEC0030
EHEC0040
EHEC0050
EHEC0060
EHEC0070
EHEC0080
EHEC0090
EHEC0100
EHEC0110
EHEC0120
EHEC0130
EHEC0140
EHEC0150
EHEC0160
EHEC0170
EHEC0180
EHEC0190
EHEC0200
EHEC0210
EHEC0220
EHEC0230
EHEC0240
EHEC0250
EHEC0260
EHEC0270
EHEC0280
EHEC0290
EHEC0300
EHEC0310
EHEC0320
EHEC0330
EHEC0340

```

```

SUBROUTINE EHESSC (AR, AI, K, L, N, IA, ID)
-----D-----LIBRARY I-----
FUNCTION
USAGE
PARAMETERS AR
AI
K
L
N
IA
ID
PRECISION
-- REDUCTION OF A COMPLEX MATRIX TO COMPLEX
UPPER HESSENBERG FORM.
-- CALL EHESSC(AR, AI, K, L, N, IA, ID)
-- INPUT/OUTPUT MATRIX OF DIMENSION N BY N. OF
THE MATRIX TO BE REDUCED.
ON OUTPUT CCNTAINS THE REAL COMPONENTS
OF THE REDUCED HESSENBERG FORM IN THE
UPPER TRIANGULAR PORTION (INCLUDING MAIN
AND SUB-DIAGONAL) AND THE DETAILS OF
THE REDUCTION IN THE LOWER TRIANGULAR
PORTION.
-- INPUT/OUTPUT MATRIX OF DIMENSION N BY N
CONTAINING THE IMAGINARY COUNTERPARTS
TO AR, ABOVE.
-- INPUT SCALAR CONTAINING THE ROW AND COLUMN
INDEX OF THE STARTING ELEMENT TO BE
REDUCED BY ROW SCALING. FOR UNBALANCED
MATRICES SET K = 1.
-- INPUT SCALAR CONTAINING THE ROW AND
COLUMN INDEX OF THE LAST ELEMENT TO BE
REDUCED BY ROW SCALING. FOR UNBALANCED
MATRICES SET L = N.
-- INPUT SCALAR CONTAINING THE ORDER OF
THE MATRIX TO BE REDUCED.
-- INPUT SCALAR CCNTAINING ROW DIMENSION
OF AR AND AI IN THE CALLING PROGRAM.
-- OUTPUT VECTOR OF LENGTH L CONTAINING
DETAILS OF THE TRANSFORMATIONS.
-- SINGLE/DOUBLE

```



```

C LANGUAGE - FORTRAN
C -----
C LATEST REVISION - FEBRUARY 7, 1973
C
C SUBROUTINE EHESSC (AR, AI, K, L, N, IA, ID)
C
C DIMENSION
C DOUBLE PRECISION
C COMPLEX *16
C EQUIVALENCE
C
C DATA
C LA=L-1
C KPI=K+1
C IF (LA .LT. KPI) GO TO 45
C DC 40 M=KPI, LA
C I=M
C XR=ZERO
C XI=ZERO
C DO 5 J=M, L
C IF (DABS (AR(J, M-1))+DABS (AI(J, M-1)) .LE. DABS (XR)+DABS (XI))
C GO TO 5
C XR=AR(J, M-1)
C XI=AI(J, M-1)
C I=J
C CONTINUE
C ID(M)=I
C IF (I .EQ. M) GO TO 20
C
C MM1=M-1
C DO 10 J=MM1, N
C YR=AR(I, J)
C AR(I, J)=AR(M, J)
C AR(M, J)=YR
C YI=AI(I, J)
C AI(I, J)=AI(M, J)
C AI(M, J)=YI
C CONTINUE
C DO 15 J=1, L
C YR=AR(J, I)
C AR(J, I)=AR(J, M)
C AR(J, M)=YR
C YI=AI(J, I)
C AI(J, I)=AI(J, M)
C AI(J, M)=YI
C CONTINUE
C
C INTERCHANGE ROWS AND COLUMNS OF
C ARRAYS AR AND AI
C
C END INTERCHANGE

```

```

EHEC0350
EHEC0360
EHEC0370
EHEC0380
EHEC0390
EHEC0400
EHEC0410
EHEC0420
EHEC0430
EHEC0450
EHEC0460
EHEC0470
EHEC0490
EHEC0500
EHEC0510
EHEC0520
EHEC0530
EHEC0540
EHEC0550
EHEC0560
EHEC0570
EHEC0580
EHEC0610
EHEC0620
EHEC0630
EHEC0640
EHEC0650
EHEC0660
EHEC0670
EHEC0680
EHEC0690
EHEC0700
EHEC0710
EHEC0720
EHEC0730
EHEC0740
EHEC0750
EHEC0760
EHEC0770
EHEC0780
EHEC0790
EHEC0800
EHEC0810
EHEC0820
EHEC0830
EHEC0840
EHEC0850
EHEC0860

```



```

20 IF (XR .EQ. ZERO .AND. XI .EQ. ZERO) GO TO 40
   MPI=M+1
   DO 35 I=MPI,L
     YR=AR(I,M-1)
     YI=AI(I,M-1)
     IF (YR .EQ. ZERO .AND. YI .EQ. ZERO) GO TO 35
     Y=Y/X
     AR(I,M-1)=YR
     AI(I,M-1)=YI
     AR(I,J)=AR(I,J)-YR*AR(M,J)+YI*AI(M,J)
     AI(I,J)=AI(I,J)-YR*AI(M,J)-YI*AR(M,J)
     DO 25 J=M,N
       CONTINUE
     DO 30 J=1,L
       AR(J,M)+YR*AR(J,I)-YI*AI(J,I)
       AI(J,M)=AI(J,I)+YI*AR(J,I)
       CONTINUE
     CONTINUE
30 CONTINUE
35 CONTINUE
40 RETURN
45 END

```

```

EHEC0870
EHEC0880
EHEC0890
EHEC0900
EHEC0910
EHEC0920
EHEC0930
EHEC0940
EHEC0950
EHEC0960
EHEC0970
EHEC0980
EHEC0990
EHEC1000
EHEC1010
EHEC1020
EHEC1030
EHEC1040
EHEC1050
EHEC1060
EHEC1070

```

```

C-----D-----LIBRARY I-----
SUBROUTINE ELRH2C (HR,HI,K,L,N,IH,WR,WI,ZR,ZI,ID,INFER,IER)
FUNCTION
USAGE
PARAMETERS FR
HI
K
L
N
-- COMPUTE THE EIGENVALUES AND EIGENVECTORS OF
A COMPLEX UPPER HESSENBERG MATRIX AND
BACK TRANSFORM THE EIGENVECTORS.
-- CALL ELRH2C (FR,HI,K,L,N,IH,WR,WI,ZR,ZI,ID,
INFER,IER)
-- INPUT MATRIX OF DIMENSION N BY N CONTAINING
THE REAL COMPONENTS OF THE COMPLEX
HESSENBERG MATRIX. HR IS DESTROYED ON
OUTPUT.
-- INPUT MATRIX OF DIMENSION N BY N CONTAINING
THE IMAGINARY COUNTERPARTS TO HR, ABOVE.
HI IS DESTROYED ON OUTPUT.
-- INPUT SCALAR CONTAINING THE LOWER BOUNDARY
INDEX FOR THE INPUT MATRIX.
FOR UNBALANCED MATRICES SET K = 1.
-- INPUT SCALAR CONTAINING THE UPPER BOUNDARY
INDEX FOR THE INPUT MATRIX.
FOR UNBALANCED MATRICES SET L = N.
-- INPUT SCALAR CONTAINING THE ORDER OF THE
HESSENBERG MATRIX AND THE EIGENVECTOR
MATRIX.

```

```

ELR20010
ELR20020
ELR20030
ELR20040
ELR20050
ELR20060
ELR20070
ELR20080
ELR20090
ELR20100
ELR20110
ELR20120
ELR20130
ELR20140
ELR20150
ELR20160
ELR20170
ELR20180
ELR20190
ELR20200
ELR20210
ELR20220
ELR20230
ELR20240
ELR20250

```


ELR20750
 ELR20770
 ELR20780
 ELR20790
 ELR20800
 ELR20810
 ELR20820
 ELR20830
 ELR20840
 ELR20850
 ELR20860
 ELR20870
 ELR20880
 ELR20890
 ELR20900
 ELR20910
 ELR20920
 ELR20930
 ELR20940
 ELR20950
 ELR20960
 ELR20970
 ELR20980
 ELR20990
 ELR21000
 ELR21010
 ELR21020
 ELR21030
 ELR21040
 ELR21050
 ELR21060
 ELR21070
 ELR21080
 ELR21090
 ELR21100
 ELR21110
 ELR21120
 ELR21130
 ELR21140
 ELR21150
 ELR21160
 ELR21170
 ELR21180
 ELR21190
 ELR21200
 ELR21210
 ELR21220
 ELR21230

```

C      DATA      EPS/23410000000000000000/
      INITIALIZE IER
      IER=0
      INFER=0
      TR=ZERO
      TJ=ZERO
      DC 5 I=1,N
      DO 3 J=1,N
      ZR(I,J)=ZERO
      ZI(I,J)=ZERO
      CONTINUE
      ZR(I,I) = ONE
      3 CONTINUE
      5 CONTINUE
      IEND=L-K-1
      IF (IEND .LE. 0) GO TO 25
      DO 20 II=1,IEND
      I=L-II
      IP1=I+1
      IM1=I-1
      DO 10 M=IP1,L
      ZR(M,I)=HR(M,IM1)
      ZI(M,I)=HI(M,IM1)
      10 CONTINUE
      J=ID(I)
      IF (I .EQ. J) GO TO 20
      DO 15 M=I,L
      ZR(I,M)=ZR(J,M)
      ZI(I,M)=ZI(J,M)
      ZR(J,M)=ZERO
      ZI(J,M)=ZERO
      15 CONTINUE
      ZR(J,I)=ONE
      20 CONTINUE
      25 DO 30 I=1,N
      IF (I .GE. K .AND. I .LE. L) GO TO 30
      WR(I)=HR(I,I)
      WI(I)=HI(I,I)
      30 CCNTINUE
      NN=L
      C 35 IF (NN .LT. K) GO TO 150
      ITS=0
      NNM1=NN-1
      NNM2=NN-2
      SEARCH FOR NEXT EIGENVALUE
      FORM THE MATRIX OF ACCUMULATED
      TRANSFORMATIONS FROM THE INFOR-
      MATION LEFT BY ROUTINE 'EHESSC'
      DO I=L-1,K+1,-1
  
```



```

C      IF (NN .EQ. K) GO TO 50
C      LOOK FOR SINGLE SMALL SUB-DIAGONAL
C      ELEMENT
C      DO M=NN,K+1,-1
C      DO 45 KK=K,NNM1
C      M=NPL-KK
C      MMI=M-1
C      IF (DABS(HR(M,MM1))+DABS(HI(M,MM1)) .LE. EPS*(DABS(HR(MM1,MM1)) GO TO 55
C      +DABS(HI(MM1,MM1))+DABS(HR(M,M)))+DABS(HI(M,M))) GO TO 55
C      1
C      45 CONTINUE
C      50 M=K
C      55 IF (M .EQ. NN) GO TO 145
C      IF (ITS .EQ. 30) GO TO 205
C      IF (ITS .EQ. 10 .OR. ITS .EQ. 20) GO TO 60
C      SR=HR(NN,NN)
C      SI=HI(NN,NN)
C      XR=HR(NN,NN)*HR(NN,NNM1)-HI(NN,NN)*HI(NN,NNM1)
C      XI=HR(NN,NNM1,NN)*HI(NN,NNM1)+HI(NN,NNM1,NN)*HR(NN,NNM1)
C      IF (XR .EQ. ZERO .AND. XI .EQ. ZERO) GO TO 65
C      YR=(HR(NN,NNM1,NNM1)-SR)/TWO
C      YI=(HI(NN,NNM1,NNM1)-SI)/TWO
C      Z=CDSQR T(DCMPLX(YR**2-YI**2+XR,TWC*YR*YI+XI))
C      IF (YR*ZZR+YI*ZZI .LT. ZERO) Z=-Z
C      X=X/(Y+Z)
C      SR=SR-XR
C      SI=SI-XI
C      GO TO 65
C      60 SR=DABS(HR(NN,NNM1))+DABS(HR(NN,NNM2))
C      SI=DABS(HI(NN,NNM1))+DABS(HI(NN,NNM2))
C      65 DO 70 I=K,NN
C      HR(I,I)=HR(I,I)-SR
C      HI(I,I)=HI(I,I)-SI
C      70 CONTINUE
C      TR=TR+SR
C      TI=TI+SI
C      ITS=ITS+1
C      XR=DABS(HR(NN,NNM1))+DABS(HI(NN,NNM1))
C      YR=DABS(HR(NN,NNM1))+DABS(HI(NN,NNM1))
C      ZZR=DABS(HR(NN,NNM1))+DABS(HI(NN,NNM1))
C      NNMJ=NNM1-M
C      IF (NNMJ .EQ. 0) GO TO 80
C      DO 75 NM=1,NNMJ
C      MM=NN-NM
C      DO MM=NN-1,M+1,-1
C      LOOK FOR TWO CONSECUTIVE SMALL
C      SUB-DIAGONAL ELEMENTS
C      LOOK FOR SINGLE SMALL SUB-DIAGONAL
C      ELEMENT
C      DO M=NN,K+1,-1
C      DO 45 KK=K,NNM1
C      M=NPL-KK
C      MMI=M-1
C      IF (DABS(HR(M,MM1))+DABS(HI(M,MM1)) .LE. EPS*(DABS(HR(MM1,MM1)) GO TO 55
C      +DABS(HI(MM1,MM1))+DABS(HR(M,M)))+DABS(HI(M,M))) GO TO 55
C      1
C      45 CONTINUE
C      50 M=K
C      55 IF (M .EQ. NN) GO TO 145
C      IF (ITS .EQ. 30) GO TO 205
C      IF (ITS .EQ. 10 .OR. ITS .EQ. 20) GO TO 60
C      SR=HR(NN,NN)
C      SI=HI(NN,NN)
C      XR=HR(NN,NN)*HR(NN,NNM1)-HI(NN,NN)*HI(NN,NNM1)
C      XI=HR(NN,NNM1,NN)*HI(NN,NNM1)+HI(NN,NNM1,NN)*HR(NN,NNM1)
C      IF (XR .EQ. ZERO .AND. XI .EQ. ZERO) GO TO 65
C      YR=(HR(NN,NNM1,NNM1)-SR)/TWO
C      YI=(HI(NN,NNM1,NNM1)-SI)/TWO
C      Z=CDSQR T(DCMPLX(YR**2-YI**2+XR,TWC*YR*YI+XI))
C      IF (YR*ZZR+YI*ZZI .LT. ZERO) Z=-Z
C      X=X/(Y+Z)
C      SR=SR-XR
C      SI=SI-XI
C      GO TO 65
C      60 SR=DABS(HR(NN,NNM1))+DABS(HR(NN,NNM2))
C      SI=DABS(HI(NN,NNM1))+DABS(HI(NN,NNM2))
C      65 DO 70 I=K,NN
C      HR(I,I)=HR(I,I)-SR
C      HI(I,I)=HI(I,I)-SI
C      70 CONTINUE
C      TR=TR+SR
C      TI=TI+SI
C      ITS=ITS+1
C      XR=DABS(HR(NN,NNM1))+DABS(HI(NN,NNM1))
C      YR=DABS(HR(NN,NNM1))+DABS(HI(NN,NNM1))
C      ZZR=DABS(HR(NN,NNM1))+DABS(HI(NN,NNM1))
C      NNMJ=NNM1-M
C      IF (NNMJ .EQ. 0) GO TO 80
C      DO 75 NM=1,NNMJ
C      MM=NN-NM
C      DO MM=NN-1,M+1,-1

```

```

ELR21240
ELR21250
ELR21260
ELR21270
ELR21280
ELR21290
ELR21300
ELR21310
ELR21320
ELR21330
ELR21360
ELR21370
ELR21380
ELR21390
ELR21400
ELR21410
ELR21420
ELR21430
ELR21440
ELR21450
ELR21460
ELR21470
ELR21480
ELR21490
ELR21510
ELR21520
ELR21530
ELR21540
ELR21550
ELR21560
ELR21570
ELR21600
ELR21610
ELR21620
ELR21630
ELR21640
ELR21650
ELR21660
ELR21670
ELR21680
ELR21690
ELR21700
ELR21710
ELR21750
ELR21760
ELR21770
ELR21780
ELR21790

```


ELR211800
 ELR211810
 ELR211820
 ELR211840
 ELR211850
 ELR211860
 ELR211880
 ELR211890
 ELR211900
 ELR211910
 ELR211920
 ELR211930
 ELR211940
 ELR211950
 ELR211960
 ELR211970
 ELR211980
 ELR211990
 ELR222010
 ELR222020
 ELR222030
 ELR222040
 ELR222050
 ELR222060
 ELR222070
 ELR222080
 ELR222090
 ELR222100
 ELR222110
 ELR222120
 ELR222130
 ELR222140
 ELR222150
 ELR222160
 ELR222170
 ELR222180
 ELR222190
 ELR222200
 ELR222210
 ELR222220
 ELR222230
 ELR222240
 ELR222250
 ELR222260
 ELR222270
 ELR222280
 ELR222290
 ELR222300

```

MMM1=MM-1
YI=YR
YR=DABS(HR(MM,MMM1))+DABS(HI(MM,MMM1))
XI=ZZR
ZZR=XR
XR=DABS(HR(MMM1,MMM1))+DABS(HI(MMM1,MMM1))
IF(YR.LE.EPS*ZZR/YI*(ZZR+XR+XI)) GO TO 85
75 CONTINUE
80 MM=M
C
85 MP1=MM+1
DO 110 I=MP1,NN
  IM1=I-1
  XR=HR(IM1,IM1)
  XI=HI(IM1,IM1)
  YR=HR(I,IM1)
  YI=HI(I,IM1)
  IF(DABS(XR)+DABS(XI)).GE.DABS(YR)+DABS(YI)) GO TO 95
  INTERCHANGE ROWS OF HR AND HI
  DO 90 J=IM1,N
    ZZR=HR(IM1,J)
    HR(IM1,J)=HR(I,J)
    HR(I,J)=ZZR
    ZZI=HI(IM1,J)
    HI(IM1,J)=HI(I,J)
    HI(I,J)=ZZI
  CONTINUE
  Z=X/Y
  WR(I)=ONE
  GO TO 100
  Z=Y/X
  WR(I)=-ONE
  HR(I,IM1)=ZZR
  HI(I,IM1)=ZZI
  DO 105 J=I,N
    HR(I,J)=HR(I,J)-ZZR*HR(IM1,J)+ZZI*HI(IM1,J)
    HI(I,J)=HI(I,J)-ZZR*HI(IM1,J)+ZZI*HR(IM1,J)
  CONTINUE
  COMPOSITION R*L=H
110 CONTINUE
C
DO 140 J=MP1,NN
  JM1=J-1
  XR=HR(J,JM1)
  XI=HI(J,JM1)
  HR(J,JM1)=ZERO
  HI(J,JM1)=ZERO
C
C
INTERCHANGE COLUMNS OF HR, HI,
ZR, AND ZI IF NECESSARY

```


ELR22310
 ELR22320
 ELR22330
 ELR22340
 ELR22350
 ELR22360
 ELR22370
 ELR22380
 ELR22390
 ELR22400
 ELR22410
 ELR22420
 ELR22430
 ELR22440
 ELR22450
 ELR22460
 ELR22470
 ELR22480
 ELR22490
 ELR22500
 ELR22510
 ELR22520
 ELR22530
 ELR22540
 ELR22550
 ELR22560
 ELR22570
 ELR22580
 ELR22590
 ELR22600
 ELR22610
 ELR22620
 ELR22630
 ELR22640
 ELR22650
 ELR22660
 ELR22670
 ELR22680
 ELR22690
 ELR22700
 ELR22710
 ELR22720
 ELR22730
 ELR22740
 ELR22750
 ELR22760
 ELR22780
 ELR22790
 ELR22800

```

IF (WR(J) .LE. ZERC) GO TO 125
DO 115 I=1, J
  ZR=HR(I, J, JMI)
  HR(I, JMI)=HR(I, J)
  HR(I, J)=ZR
  ZZI=HI(I, JMI)
  HI(I, JMI)=HI(I, J)
  HI(I, J)=ZZI
  CONTINUE
DO 120 I=K, L
  ZR=ZR(I, JMI)
  ZR(I, JMI)=ZR(I, J)
  ZR(I, J)=ZR
  ZZI=ZI(I, JMI)
  ZI(I, JMI)=ZI(I, J)
  ZI(I, J)=ZZI
  CONTINUE
C
120 CONTINUE
END INTERCHANGE COLUMNS
DO 125 I=1, J
  HR(I, JMI)=HR(I, JMI)+XR*HR(I, J)-XI*HI(I, J)
  HI(I, JMI)=HI(I, JMI)+XR*HI(I, J)+XI*HR(I, J)
  CONTINUE
DO 135 I=K, L
  ZR(I, JMI)=ZR(I, JMI)+XR*ZR(I, J)-XI*ZI(I, J)
  ZI(I, JMI)=ZI(I, JMI)+XR*ZI(I, J)+XI*ZR(I, J)
  CONTINUE
C
135 CONTINUE
END ACCUMULATE TRANSFORMATIONS
GC TO 40
C
140 CONTINUE
A ROOT FOJND
WR(NN)=HR(NN, NN)+TR
WI(NN)=HI(NN, NN)+TI
NN=NNMI
GO TO 35
C
145 CONTINUE
ALL ROOTS FOUND. BACKSUBSTITUTE TO
FIND VECTORS OF UPPER TRIANGULAR
FORM
IF (N .EQ. 1) GO TO 9005
FNORM=ZERO
DC 160 I=1, N
  FNORM=FNORM+DABS(WR(I))+DABS(WI(I))
  IF (I .EQ. N) GO TO 160
  IPI=I+1
  DO 155 J=IPI, N
    FNORM=FNORM+DABS(HR(I, J))+DABS(HI(I, J))
  CONTINUE
C
155 CONTINUE
IF (FNORM .EQ. ZERO) GO TO 9005

```



```

C      NP2 = N+2
      DO 180 NM=2,N
        NN=NP2-NM
        XR=WR(NN)
        XI=WI(NN)
        NNM1=NN-1
C      DO 175 II=1, NNM1
        I=NN-II
        ZZR=HR(I, NN)
        ZZI=HI(I, NN)
        IF (I .EQ. NNM1) GO TO 170
        IP1=I+1
        DO 165 J=IP1, NNM1
          ZZR=ZZR+HR(I, J)*HR(J, NN)-HI(I, J)*HI(J, NN)
          ZZI=ZZI+HR(I, J)*HI(J, NN)+HI(I, J)*HR(J, NN)
        C CONTINUE
        YR=XR-WR(I)
        YI=XI-WI(I)
        IF (YR .EQ. ZERO .AND. YI .EQ. ZERO) YR=EPS*FNORM
        Z=Z/Y
        HR(I, NN)=T3(1)
        HI(I, NN)=T3(2)
      C CONTINUE
165      CONTINUE
170      CONTINUE
      NMI=N-1
C      DO 190 I=1, NMI
        IF (I .GE. K .AND. I .LE. L) GO TO 190
        IP1=I+1
        DO 185 J=IP1, N
          ZR(I, J)=HR(I, J)
          ZI(I, J)=HI(I, J)
        C CONTINUE
185      CONTINUE
190      CONTINUE
      IF (L .EQ. 0) GO TO 9005
C      NPL=N+K
      DO 200 JJ=K, NMI
        J=APL-JJ
        JM1=J-1
        DO 200 I=K, L
          ZZR=ZR(I, J)

```

```

ELR222810
ELR222820
ELR222830
ELR222840
ELR222850
ELR222860
ELR222870
ELR222880
ELR222890
ELR222900
ELR222910
ELR222920
ELR222930
ELR222940
ELR222950
ELR222960
ELR222970
ELR222980
ELR222990
ELR23000
ELR23010
ELR23020
ELR23030
ELR23040
ELR23050
ELR23060
ELR23070
ELR23080
ELR23090
ELR23100
ELR23110
ELR23120
ELR23130
ELR23140
ELR23150
ELR23160
ELR23170
ELR23180
ELR23190
ELR23200
ELR23210
ELR23220
ELR23230
ELR23240
ELR23250
ELR23260
ELR23270
ELR23280

```

```

MULTIPLY BY TRANSFORMATION MATRIX
TO GIVE VECTORS OF ORIGINAL FULL
MATRIX

```

```

DO J=N, K+1, -1

```


ELR233290
 ELR233300
 ELR233310
 ELR233320
 ELR233330
 ELR233340
 ELR233350
 ELR233360
 ELR233370
 ELR233380
 ELR233390
 ELR233400
 ELR233410
 ELR233420
 ELR233430
 ELR233440
 ELR233450
 ELR233460
 ELR233470

```

ZZI=ZJ(I,J)
MM=JMI
IF (L.LT. J) MM=L
DO 195 M=K,MM
  ZZR=ZZR+ZR(I,M)*HR(M,J)-ZI(I,M)*HI(M,J)
  ZZI=ZZI+ZR(I,M)*HI(M,J)+ZI(I,M)*HR(M,J)
CONTINUE
ZR(I,J)=ZZR
ZI(I,J)=ZZI
195 CONTINUE
200 CONTINUE
GO TO 9005

C
C
205 IER=129
INFER=NN
CONTINUE
9000 CALL UERTST (IER,6HEL RH2C)
9005 RETURN
END

```

SET ERROR - NO CONVERGENCE TO AN
 EIGENVALUE AFTER 30 ITERATIONS

EBBC0010
 EBBC0020
 EBBC0030
 EBBC0040
 EBBC0050
 EBBC0060
 EBBC0070
 EBBC0080
 EBBC0090
 EBBC0100
 EBBC0110
 EBBC0120
 EBBC0130
 EBBC0140
 EBBC0150
 EBBC0160
 EBBC0170
 EBBC0180
 EBBC0190
 EBBC0200
 EBBC0210
 EBBC0220
 EBBC0230
 EBBC0240
 EBBC0250
 EBBC0260
 EBBC0270

```

SUBROUTINE EBBCKC (ZR,ZI,N,IZ,K,L,M,D)
C-----D-----LIBRARY 1-----
C
C FUNCTION
C
C USAGE
C PARAMETERS
C
C ZR
C ZI
C
C N
C
C IZ
C
C K
C L
C
C M
C D

```

- BACKTRANSFORM THE EIGENVECTORS OF A BALANCED
 COMPLEX GENERAL MATRIX.
 - CALL EBBCKC (ZR,ZI,N,IZ,K,L,M,D)
 - INPUT/OUTPUT MATRICES OF DIMENSION N BY M.
 ON INPUT, THE FIRST M COLUMNS OF ZR AND
 ZI CONTAIN THE REAL AND IMAGINARY PARTS,
 RESPECTIVELY, OF THE EIGENVECTORS TO BE
 BACK TRANSFORMED. ON OUTPUT, THESE M
 COLUMNS CONTAIN THE REAL AND IMAGINARY
 PARTS OF THE TRANSFORMED EIGENVECTORS.
 - INPUT SCALAR CONTAINING THE NUMBER OF
 ROWS IN THE MATRIX Z = (ZR,ZI). N MUST
 NOT BE GREATER THAN IZ.
 - INPUT SCALAR CONTAINING THE ROW DIMENSION
 OF MATRICES ZR AND ZI IN THE CALLING
 PROGRAM.
 - INPUT SCALARS CONTAINING THE BOUNDARY
 INDICES FOR THE BALANCED MATRIX. K AND L
 ARE TWO OUTPUT PARAMETERS FROM IMSL ROUTINE
 EBALAC.
 - INPUT SCALAR CONTAINING THE NUMBER OF COLUMNS
 OF Z = (ZR,ZI) TO BE BACK TRANSFORMED.
 - INPUT VECTOR OF LENGTH N CONTAINING THE

DETAILS OF THE TRANSFORMATIONS PRODUCED
BY IMSL ROUTINE EBALAC.

```

C      PRECISION          EBBC0280
C      LANGUAGE          EBBC0290
C      - SINGLE/DOUBLE  EBBC0300
C      - FORTRAN        EBBC0310
C-----
C      LATEST REVISION   - MARCH 9, 1977  EBBC0320
C
C      SUBROUTINE EBCKC (ZR,ZI,N,IZ,K,L,M,D) EBBC0330
C
C      DIMENSION          ZR(IZ,1),ZI(IZ,1),D(1) EBBC0340
C      DOUBLE PRECISION  ZR,ZI,D,S      EBBC0350
C      IF (L.EC.K) GO TO 15 EBBC0360
C      DO 10 I = K,L      EBBC0370
C      S = D(I)          EBBC0380
C
C      DO 5 J = 1,M      EBBC0390
C      ZR(I,J) = ZR(I,J)*S EBBC0400
C      ZI(I,J) = ZI(I,J)*S EBBC0410
C
C      5 CONTINUE       EBBC0420
C
C      10 CONTINUE      EBBC0430
C
C      15 DO 25 I = 1,N EBBC0440
C      I = II           EBBC0450
C      IF (I.GE.K.AND.I.LE.L) GO TO 25 EBBC0460
C      IF (I.LT.K) I = K-II EBBC0470
C      KK = D(I)        EBBC0480
C      IF (KK.EQ.I) GO TO 25 EBBC0490
C      DO 20 J = 1,M    EBBC0500
C      S = ZR(I,J)      EBBC0510
C      ZR(I,J) = ZR(KK,J) EBBC0520
C      ZR(KK,J) = S     EBBC0530
C      S = ZI(I,J)      EBBC0540
C      ZI(I,J) = ZI(KK,J) EBBC0550
C      ZI(KK,J) = S     EBBC0560
C
C      20 CONTINUE     EBBC0570
C      25 CONTINUE     EBBC0580
C      RETURN          EBBC0590
C      END             EBBC0600
C
C      SUBROUTINE UERTST ( IER,NAME )      EBBC0610
C-----
C      UERTST-----LIBRARY 1----- EBBC0620
C      FUNCTION          - ERROR MESSAGE GENERATION EBBC0630
C
C      UERT0010          EBBC0640
C      UERT0020          EBBC0650
C      UERT0030          EBBC0660
C      UERT0040          EBBC0670
C      UERT0050          EBBC0680

```

LEFT HAND EIGENVECTORS ARE BACK
TRANSFORMED IF THE ABOVE
STATEMENT IS REPLACED BY S=1.0/D(I)

DO 25 I=K-1,1,-1 AND
DO I=L+1,N,1


```

//EIG$FCN JOB (1719,0947,AX74), 'SMC 1882', TIME=2
//EXEC FORTCLGW
//FORT.SYSIN DD *
.....
PROGRAM EIGFCN
PERTURBATION VELOCITY PLOT PROGRAM
  NI = 0
PURPOSE
  TO PLOT THE NONDIMENSIONALIZED PERTURBATION VELOCITY U AGAINST
  NONDIMENSIONALIZED RADIUS UTILIZING THE DATA GENERATED
  BY PROGRAM PIPEO (MODE=1). PLOTTING IS PERFORMED ON
  THE NPS VERSATEC PLOTTER USING SUBROUTINE PLOTG.
.....
IMPLICIT REAL*8(A-H,O-Z)
COMPLEX *16QPRIM(85),DQPRIM(85),ALPHA,UP(85),UPPRIM(85),CONST,GAMMA
IA
REAL *8ETA(85),U(85),UR(85),UI(85)
REAL *4URI(85),UII(85),RADI(85),SREY,SLAMDA,SGAMMA,AR,AI
READ N,REY,ALPHA,LAMBDA,GAMMA & QPRIME'S
READ (5,6) N,REY,ALPHA
NO = N+1
N1 = N+2
READ (5,7) AMDA,GAMMA
READ (5,8) KSET
SLAMDA = AMDA
SREY = REY
AR = ALPHA
AI = AIMAG(ALPHA)
SGAMMA = GAMMA
SGAMMA = SGAMMA + AR
DC 1 I=2,NO
READ (5,9) ETA(I),QPRIM(I)
1 CONTINUE

```

```

EIGF 10
EIGF 20
EIGF 30
EIGF 40
EIGF 50
EIGF 60
EIGF 70
EIGF 80
EIGF 90
EIGF 100
EIGF 110
EIGF 120
EIGF 130
EIGF 140
EIGF 150
EIGF 160
EIGF 170
EIGF 180
EIGF 190
EIGF 200
EIGF 210
EIGF 220
EIGF 230
EIGF 240
EIGF 250
EIGF 260
EIGF 270
EIGF 280
EIGF 290
EIGF 300
EIGF 310
EIGF 320
EIGF 330
EIGF 340
EIGF 350
EIGF 360
EIGF 370
EIGF 380
EIGF 390
EIGF 400
EIGF 410
EIGF 420
EIGF 430
EIGF 440
EIGF 450

```



```

C      COMPUTE UPPRIME'S
C      DEL = 1D0/DFLOAT(N+1)
C      ETA(I) = 0.0D0
C      QPRIM(1) = 1.795918367D0*QPRIM(2) - 1.24781341D0*QPRIM(3) + 0.60641399
14D0*QPRIM(4) - 0.177842566D0*QPRIM(5) + 0.023323615D0*QPRIM(6)
C      UPPRIM(1) = 2D0*QPRIM(1)
C      CALL COEFNT (ETA(2), AMDA, COEF, KSET)
C      DQPRIM(2) = 2D0*(-QPRIM(2)+QPRIM(3))/(3D0*DEL)
C      DQPRIM(2) = COEF*DQPRIM(2)
C      UPPRIM(2) = 2D0*QPRIM(2)+ETA(2)*DQPRIM(2)
C      CALL COEFNT (ETA(3), AMDA, COEF, KSET)
C      DQPRIM(3) = 4D0*(-QPRIM(2)+QPRIM(3))/(3D0*DEL)
C      DQPRIM(3) = COEF*DQPRIM(3)
C      UPPRIM(3) = 2D0*QPRIM(3)+ETA(3)*DQPRIM(3)
C
C      DO 2 I=4,N
C      CALL COEFNT (ETA(I), AMDA, COEF, KSET)
C      DQPRIM(I) = (QPRIM(I+1)-QPRIM(I-1))/(2D0*DEL)
C      DQPRIM(I) = COEF*DQPRIM(I)
C      UPPRIM(I) = 2D0*QPRIM(I)+ETA(I)*DQPRIM(I)
C      2 CONTINUE
C
C      CALL COEFNT (ETA(N0), AMDA, COEF, KSET)
C      DQPRIM(N0) = -QPRIM(N)/(2D0*DEL)
C      DQPRIM(N0) = COEF*DQPRIM(N0)
C      UPPRIM(N0) = 2D0*QPRIM(N0)+ETA(N0)*DQPRIM(N0)
C
C      ETA(N1) = 1.0D0
C      UPPRIM(N1) = (0D0,0D0)
C      WRITE (6,10)
C      DETERMINE U VECTOR OF LARGEST MAGNITUDE
C      C = 0.0D0
C
C      DO 3 I=1,N1
C      IF (CDABS(UPPRIM(I)).GT.C) INDEX=I
C      IF (CDABS(UPPRIM(I)).GT.C) C=CDABS(UPPRIM(I))
C      3 CONTINUE
C
C      CONST = DCONJG(UPPRIM(INDEX))/C**2

```

```

EIGF 460
EIGF 470
EIGF 480
EIGF 490
EIGF 500
EIGF 510
EIGF 520
EIGF 530
EIGF 540
EIGF 550
EIGF 560
EIGF 570
EIGF 580
EIGF 590
EIGF 600
EIGF 610
EIGF 620
EIGF 630
EIGF 640
EIGF 650
EIGF 660
EIGF 670
EIGF 680
EIGF 690
EIGF 700
EIGF 710
EIGF 720
EIGF 730
EIGF 740
EIGF 750
EIGF 760
EIGF 770
EIGF 780
EIGF 790
EIGF 800
EIGF 810
EIGF 820
EIGF 830
EIGF 840
EIGF 850
EIGF 860
EIGF 870
EIGF 880
EIGF 890
EIGF 900
EIGF 910
EIGF 920
EIGF 930

```



```

C          NORMALIZE UPRIMES'S AND SPLIT INTO REAL & IMAGINARY VECTORS
C
C          DO 4 I=1,N1
C          UP(I) = CONST*UPPRIM(I)
C          UR(I) = UP(I)
C          UI(I) = (ODO,-1DO)*UP(I)
C          4 CONTINUE
C
C          CONVERT U'S AND ETA'S TO SINGLE PRECISION FCR PLOTG
C
C          DO 5 I=1,N1
C          RAD1(I) = ETA(I)
C          UR1(I) = UR(I)
C          UI1(I) = UI(I)
C          WRITE (6,11) RAD1(I),UR1(I),UI1(I)
C          5 CONTINUE
C
C          PLOT RESULTS
C          CALL PLOTG(RAD1,UR1,N1,1,1,1,'RADIUS',6,'PERTURBATION VELOCITY',
C          $ 21,0,1,1,1,1,7,7)
C          CALL PLOTG(RAD1,UI1,N1,2,1,5,'RADIUS',6,'PERTURBATION VELOCITY',
C          $ 21,0,1,1,1,1,7,7)
C          CALL CHART (N,SREY,AR,AI,SGAMMA,SLAMDA)
C          CALL PLOT (0.0,0.0,999)
C          STOP
C
C          6 FORMAT (I2,3D20.10)
C          7 FORMAT (F15.7,2(1PD20.10))
C          8 FORMAT (I2)
C          9 FORMAT (F15.7,2(1PD20.10))
C          10 FORMAT (.1,;)
C          11 FORMAT (.1,3F15.7)
C          END
C
C          .....SUBROUTINE COEFFNT(ETA,AMDA,COEF,KSET).....
C          COEF 10
C          COEF 20
C          COEF 30
C          COEF 40
C          COEF 50
C          COEF 60
C          COEF 70
C
C          PURPOSE--WHEN AN OFFSET MESH IS USED, THIS SUBROUTINE GENERATES
C          THE COEFFICIENT REQUIRED TO CONVERT DQ/DELTA TO DQ/DR AND
C          CONVERTS THE UNIFORM ETA VALUE INTO THE NONUNIFORM R VALUE

```



```

3 COEF = 1D0
  RETURN
  END
.....SUBROUTINE CHART(N,SREY,AR,AI,SGAMMA,SLAMDA).....
PURPOSE
  TO LABEL THE GRAPH WITH INFORMATION PERTAINING TO THE PLOT
EXAMPLE OF THE CALLING ARGUMENT
  CALL CHART(N,SREY,AR,AI,SGAMMA,SLAMDA)
DESCRIPTION OF PARAMETERS
  THE PARAMETERS ARE SELF-EXPLANATORY AND MUST BE IN SINGLE
  PRECISION FOR PLOTTING.
OTHER SUBROUTINES NEEDED
  ONLY BUILT-IN VERSATEC PLOTTING FUNCTIONS NEWPEN,SYMBOL &
  NUMBER. NOTE THAT THESE ROUTINES MAY ONLY BE
  ACCESSED WHEN RUNNING UNDER 'FORTCLGW'.
.....
SUBROUTINE CHART (N,SREY,AR,AI,SGAMMA,SLAMDA)
  X0 = 2.5
  Y0 = 6.5
  HT1 = 0.15
  HT1 = 0.7*HT
  DELY1 = .08+HT
  DELY2 = .065+HT1
  DELX = .1
  GRAPH TITLE
  CALL NEWPEN (2)
  CALL SYMBOL(X0,Y0,HT,'NORMALIZED PERTURBATION VELOCITY',0.,32)

```

```

COEF 560
COEF 570
COEF 580
CHAR 10
CHAR 20
CHAR 30
CHAR 40
CHAR 50
CHAR 60
CHAR 70
CHAR 80
CHAR 90
CHAR 100
CHAR 110
CHAR 120
CHAR 130
CHAR 140
CHAR 150
CHAR 160
CHAR 170
CHAR 180
CHAR 190
CHAR 200
CHAR 210
CHAR 220
CHAR 230
CHAR 240
CHAR 250
CHAR 260
CHAR 270
CHAR 280
CHAR 290
CHAR 300
CHAR 310
CHAR 320
CHAR 330
CHAR 340
CHAR 350
CHAR 360
CHAR 370
CHAR 380
CHAR 390
CHAR 400
CHAR 410
CHAR 420
CHAR 430

```



```

X0 = X0+7.*DELX
Y0 = Y0-DELY1
CALL SYMBOL(X0,Y0,HT,'FOR THE CASE N = 0',0.,18)
MESH VALUE
CALL NEWPEN(1)
X0 = X0+4.*DELX
Y0 = Y0-DELY1
SN = FLOAT(N)
CALL SYMBOL(X0,Y0,HT1,'NMESH = ',0.,9)
CALL NUMBER(999.,999.,HT1,SN,0.,-1)
REY VALUE
YC = Y0-DELY2
CALL SYMBOL(X0,Y0,HT1,'REY = ',0.,9)
CALL NUMBER(999.,999.,HT1,SREY,0.,-1)
ALPHA VALUE
X1 = X0+11.*DELY2
Y0 = Y0-DELY2
CALL SYMBOL(X0,Y0,HT1,'ALPHA = ',0.,9)
CALL NUMBER(999.,999.,HT1,AR,0.,1)
CALL NUMBER(X1,Y0,HT1,AI,0.,1)
GAMMA RL* VALUE
Y0 = Y0-DELY2
CALL SYMBOL(X0,Y0,HT1,'GAMMA* = ',0.,9)
CALL NUMBER(999.,999.,HT1,SGAMMA,0.,4)
LAMBDA VALUE
Y0 = Y0-DELY2
CALL SYMBOL(X0,Y0,HT1,'LAMBDA = ',0.,9)
CALL NUMBER(999.,999.,HT1,SLAMDA,0.,1)
SYMBOL LEGEND
Y0 = 1.75
CALL SYMBOL(X0,Y0,HT1,'OCTAGON = U(REAL)',0.,17)
Y0 = Y0-DELY2
CALL SYMBOL(X0,Y0,HT1,'DIAMOND = U(IMAG)',0.,17)
RETURN
END

```

```

CHAR 440
CHAR 450
CHAR 460
CHAR 470
CHAR 480
CHAR 490
CHAR 500
CHAR 510
CHAR 520
CHAR 530
CHAR 540
CHAR 550
CHAR 560
CHAR 570
CHAR 580
CHAR 590
CHAR 600
CHAR 610
CHAR 620
CHAR 630
CHAR 640
CHAR 650
CHAR 660
CHAR 670
CHAR 680
CHAR 690
CHAR 700
CHAR 710
CHAR 720
CHAR 730
CHAR 740
CHAR 750
CHAR 760
CHAR 770
CHAR 780
CHAR 790
CHAR 800
CHAR 810
CHAR 820
CHAR 830
CHAR 840
CHAR 850
CHAR 860
CHAR 870
CHAR 880
CHAR 890
CHAR 900

```


.....
THE FOLLGWIING CARDS COMPRISE THE DATA DECK FOR PROGRAM EIGFCN.
/*
//GO.SYSIN DD *
.
DATA DECK FROM ONE RUN OF PROGRAM PIPEO (MODENO = 1)
.
.
/*
.....


```

C      CALL SEARCH (-1,X1,Y1,NPLT1,NDIM)
C      CALL SEARCH (0,X2,Y2,NPLT2,NDIM)
C      CALL SEARCH (1,X3,Y3,NPLT3,NDIM)
C      JUMP TO PLOT LABEL ROUTINE IF NO INCIPIENT POINTS
C      IF (NPLT1) 8,8,2
C      IF POINTS COMPUTED, NEW PAGE AND WRITE THEM OUT
C      2 WRITE (6,11)
C      DC 3 I=1,NPLT1
C      WRITE (6,12) X1(I),Y1(I)
C      3 CONTINUE
C      PLOT INCIPIENT INSTABILITY POINTS
C      CALL PLOTG(X1,Y1,NPLT1,1,0,1,'ALPHA REAL',10,'ALPHA IMAGINARY',15,
C      $ XMIN,XMAX,YMIN,YMAX,7.,7.)
C      LEGEND FOR INCIPIENT SYMBOL
C      CALL NEWPEN (1)
C      CALL SYMBOL(1.3,0.7,.1,'OCTAGON = INCIPIENT INSTABILITY',0.,32)
C      JUMP TO PLOT LABEL ROUTINE IF NO CRITICAL POINTS
C      IF (NPLT2) 8,8,4
C      IF POINTS COMPUTED, NEW PAGE AND PRINT THEM OUT
C      4 WRITE (6,11)
C      DO 5 I=1,NPLT2
C      WRITE (6,12) X2(I),Y2(I)
C      5 CONTINUE
C      PLOT CRITICAL POINTS
C      CALL PLOTG(X2,Y2,NPLT2,2,0,2,'ALPHA REAL',10,'ALPHA IMAGINARY',15,
C      $ XMIN,XMAX,YMIN,YMAX,7.,7.)
C      LEGEND FOR CRITICAL SYMBOL

```

```

STBC 460
STBC 470
STBC 480
STBC 490
STBC 500
STBC 510
STBC 520
STBC 530
STBC 540
STBC 550
STBC 560
STBC 570
STBC 580
STBC 590
STBC 600
STBC 610
STBC 620
STBC 630
STBC 640
STBC 650
STBC 660
STBC 670
STBC 680
STBC 690
STBC 700
STBC 710
STBC 720
STBC 730
STBC 740
STBC 750
STBC 760
STBC 770
STBC 780
STBC 790
STBC 800
STBC 810
STBC 820
STBC 830
STBC 840
STBC 850
STBC 860
STBC 870
STBC 880
STBC 890
STBC 900
STBC 910
STBC 920
STBC 930

```



```

C      CALL NEWPEN (1)
C      CALL SYMBOL(1.3,.53,.1,'TRIANGLE = CRITICAL INSTABILITY',0.,31)
C
C      JUMP TO PLOT LABEL ROUTINE IF NO FULLY DEVELOPED POINTS
C      IF (NPLT3) 8,8,6
C
C      IF POINTS COMPUTED, NEW PAGE AND PRINT THEM OUT
C      WRITE (6,11)
C
C      DC 7 I=1,NPLT3
C      WRITE (6,12) X3(1),Y3(1)
C      7 CONTINUE
C
C      PLOT FULLY DEVELOPED POINTS
C      CALL PLOTG(X3,Y3,NPLT3,3,0,5,'ALPHA REAL',10,'ALPHA IMAGINARY',15,
C      $ XMIN,XMAX,YMIN,YMAX,7.,7.)
C      LEGEND FOR FULLY DEVELOPED SYMBOL
C
C      CALL NEWPEN (1)
C      CALL SYMBOL(1.3,.36,.1,'DIAMOND = FULLY DEVELOPED INSTABILITY',
C      $ C.,38)
C
C      LABEL THE PLOT
C      8 CALL CHART (SN,SREY,SLAMDA)
C      CALL PLOT (0.,0.,999)
C      STOP
C
C      9 FORMAT (I2)
C      10 FORMAT (3E20.10)
C      11 FORMAT ('!')
C      12 FORMAT ('.',2E20.10)
C      END
C
C.....SUBROUTINE SEARCH(NCASE,X,Y,NDIM).....
C      PURPOSE
C
C      TO SCAN THE STABILITY MAP FOR CHANGES OF SIGN WITH RESPECT
C      TO A SPECIFIED STABILITY VALUE AND GENERATE AN ARRAY OF X,Y
C      POINTS DEFINING A CONTOUR OF THE SPECIFIED STABILITY.
C
C      10
C      20
C      30
C      40
C      50
C      60
C      70
C      80

```

STBC 940
STBC 950
STBC 960
STBC 970
STBC 980
STBC 990
STBC 1000
STBC 1010
STBC 1020
STBC 1030
STBC 1040
STBC 1050
STBC 1060
STBC 1070
STBC 1080
STBC 1090
STBC 1100
STBC 1110
STBC 1120
STBC 1130
STBC 1140
STBC 1150
STBC 1160
STBC 1170
STBC 1180
STBC 1190
STBC 1200
STBC 1210
STBC 1220
STBC 1230
STBC 1240
STBC 1250
STBC 1260
STBC 1270
STBC 1280
STBC 1290
STBC 1300
STBC 1310

SEAR 10
SEAR 20
SEAR 30
SEAR 40
SEAR 50
SEAR 60
SEAR 70
SEAR 80


```

90 SEAR
100 SEAR
110 SEAR
120 SEAR
130 SEAR
140 SEAR
150 SEAR
160 SEAR
170 SEAR
180 SEAR
190 SEAR
200 SEAR
210 SEAR
220 SEAR
230 SEAR
240 SEAR
250 SEAR
260 SEAR
270 SEAR
280 SEAR
290 SEAR
300 SEAR
310 SEAR
320 SEAR
330 SEAR
340 SEAR
350 SEAR
360 SEAR
370 SEAR
380 SEAR
390 SEAR
400 SEAR
410 SEAR
420 SEAR
430 SEAR
440 SEAR
450 SEAR
460 SEAR
470 SEAR
480 SEAR
490 SEAR
500 SEAR
510 SEAR
520 SEAR
530 SEAR
540 SEAR
550 SEAR
560 SEAR

SAMPLE OF THE CALLING ARGUMENT
CALL SEARCH(NCASE,X,Y,NDIM)
DESCRIPTION OF PARAMETERS
NCASE - DEFINES THE INSTABILITY CASE (INCIPIENT,CRITICAL OR
FULLY DEVELOPED) TO BE USED WHEN GENERATING THE X,Y
ARRAYS.
NCASE = -1 INCIPIENT INSTABILITY CRITERION
NCASE = 0 CRITICAL INSTABILITY CRITERION
NCASE = 1 FULLY DEVELOPED INSTABILITY CRITERION
X - THE ARRAY OF ALPHA REAL COORDINATES DEFINING THE
LOCATION OF THE POINTS OF SPECIFIED INSTABILITY.
Y - THE ARRAY OF ALPHA IMAGINARY COORDINATES DEFINING THE
LOCATION OF THE POINTS OF THE SPECIFIED INSTABILITY.
NDIM - THE ORDER OF THE MAP ARRAY.
OTHER ROUTINES NEEDED
STATEMENT FUNCTION CRIT AND SUBROUTINE INTERP.
.....
SUBROUTINE SEARCH (NCASE,X,Y,K,NDIM)
DIMENSION X(500), Y(500)
COMMON /ARRAY/ G(41,41),AR(41),AI(41)
DEFINE THE STATEMENT FUNCTION CRIT(NCASE,ALPHA)
CRIT(NCASE,ALPHA) = FLOAT(NCASE)*ABS(ALPHA)
K = 0
NDIM = NDIM-1
SEARCH FOR SIGN CHANGES BY COLUMN & INTERPOLATE FOR
ALPHA IMAGINARY AT WHICH SIGN CHANGE OCCURS
DC 5 I=1,NDIM
DO 5 J=1,MDIM
IF (G(I,J)-CRIT(NCASE,AR(I))) 2,4,1
1 IF (G(I,J+1)-CRIT(NCASE,AR(I))) 3,3,5
2 IF (G(I,J+1)-CRIT(NCASE,AR(I))) 5,3,3

```



```

100 INTE
110 INTE
120 INTE
130 INTE
140 INTE
150 INTE
160 INTE
170 INTE
180 INTE
190 INTE
200 INTE
210 INTE
220 INTE
230 INTE
240 INTE
250 INTE
260 INTE
270 INTE
280 INTE
290 INTE
300 INTE
310 INTE

```

CALL INTERP(X1,X2,Y1,Y2,X3)
DESCRIPTION OF PARAMETERS
X1 & X2 - X-COORDINATES OF POINTS Y1 & Y2 RESPECTIVELY.
Y1 & Y2 - TWO POINTS OF OPPOSITE SIGN FOR WHICH THE POINT
OF ACTUAL SIGN CHANGE (Y = 0) IS TO BE INTERPOLATED.
X3 - THE VALUE OF X FOR WHICH Y = 0.
OTHER ROUTINES NEEDED
NONE
.....
SUBROUTINE INTERP (X1,X2,Y1,Y2,X3)
X3 = (X2*Y1-X1*Y2)/(Y1-Y2)
RETURN
END

```

10 BLKD
20 BLKD
30 BLKD
40 BLKD
50 BLKD
60 BLKD
70 BLKD
80 BLKD
90 BLKD
100 BLKD
110 BLKD
120 BLKD
130 BLKD
140 BLKD
150 BLKD
160 BLKD
170 BLKD
180 BLKD
190 BLKD
200 BLKD
210 BLKD
220 BLKD
230 BLKD
240 BLKD

```

.....BLOCK DATA.....
PURPOSE
TO INITIALIZE COMMON ARRAYS G1,ARI & AII TO ZERC.
SAMPLE OF CALLING ARGUMENT
NONE
DESCRIPTION OF PARAMETERS
G1 - THE MAP OF STABILITY VALUES GENERATED BY PIPEO.
EACH ELEMENT OF G1 IS A VALUE OF GAMMA* CORRESPONDING
TO A SPECIFIC VALUE OF THE REAL AND IMAGINARY PARTS OF
THE WAVE NUMBER, ALPHA.
ARI - THE LINEAR ARRAY OF X-COORDINATES OF THE STABILITY MAP
(THE REAL PART OF THE WAVE NUMBER, ALPHA).
AII - THE LINEAR ARRAY OF Y-COORDINATES OF THE STABILITY MAP
(THE IMAGINARY PART OF THE WAVE NUMBER, ALPHA).
OTHER ROUTINES NEEDED


```
C C C C C
NONE
.....
BLOCK DATA
COMMON /ARRAY/ G1(41,41),ARI(41),ARI(41),A11(41)
DATA G1,ARI,A11/1681*0.0,82*0.0/
END
BLKD 250
BLKD 260
BLKD 270
BLKD 280
BLKD 290
BLKD 300
BLKD 310
BLKD 320
BLKD 330
```

```
C C C C C C C C C C C C C C C C C C C C C C C C C C C C C C C C C C C C C C C C C C
.....
SUBROUTINE CHART(SN, SREY, SLAMDA) .....
PURPOSE
TO LABEL THE CONTOUR PLOT
SAMPLE OF THE CALLING ARGUMENT
CALL CHART(SN, SREY, SLAMDA)
DESCRIPTION OF PARAMETERS
SN      - THE NUMBER OF INTERIOR MESH POINTS USED FOR THE
          STABILITY CONTOUR MAP BEING PLOTTED.
SREY    - REYNOLDS NUMBER
SLAMDA  - THE NONUNIFORM MESH PARAMETER APPLICABLE TO THE
          DATA BEING PLOTTED.
OTHER ROUTINES NEEDED
ONLY BUILT-IN VERSATEC PLOTTING FUNCTIONS NEWPEN, SYMBOL &
NUMBER. NOTE THAT THESE ROUTINES MAY ONLY BE ACCESSED WHEN
RUNNING UNDER 'FORTCLGW'.
.....
SUBROUTINE CHART (SN, SREY, S, SLAMDA )
C C C
X0 = 2.5
Y0 = 6.5
HT = 0.15
HT1 = 0.7*HT
DELY1 = .08+HT
DELY2 = .065+HT1
CHAR 10
CHAR 20
CHAR 30
CHAR 40
CHAR 50
CHAR 60
CHAR 70
CHAR 80
CHAR 90
CHAR 100
CHAR 110
CHAR 120
CHAR 130
CHAR 140
CHAR 150
CHAR 160
CHAR 170
CHAR 180
CHAR 190
CHAR 200
CHAR 210
CHAR 220
CHAR 230
CHAR 240
CHAR 250
CHAR 260
CHAR 270
CHAR 280
CHAR 290
CHAR 300
CHAR 310
CHAR 320
CHAR 330
CHAR 340
CHAR 350
CHAR 360
CHAR 370
```



```

DELX = .1
GRAPH TITLE
CALL NEWPEN (2)
CALL SYMBOL(X0, Y0, HT, 'STABILITY CONTOUR PLOT', 0., 22)
X0 = X0 + 3. * DELX
YC = Y0 - DELY1
CALL SYMBOL(X0, Y0, HT, 'FOR THE CASE N = 0', 0., 18)

MESH VALUE

CALL NEWPEN (1)
X0 = X0 + 4. * DELX
Y0 = Y0 - DELY1
CALL SYMBOL(X0, Y0, HT1, 'NMESH = ', 0., 9)
CALL NUMBER (999., 999., HT1, SN, 0., -1)

REY VALUE
Y0 = Y0 - DELY2
CALL SYMBOL(X0, Y0, HT1, 'REY = ', 0., 9)
CALL NUMBER (999., 999., HT1, SREY, 0., -1)

LAMBDA VALUE
YC = Y0 - DELY2
CALL SYMBOL(X0, Y0, HT1, 'LAMBDA = ', 0., 9)
CALL NUMBER (999., 999., HT1, SLAMDA, 0., 1)

STABILITY AREA LABELS
NOTE - SINCE THE SHAPE OF THE CURVE VARIES WITH
      EACH SET OF INPUT DATA, THE COORDINATES OF THE FOLLOWING
      LABELS MUST BE ADJUSTED FOR EACH SPECIFIC PLOT.

CALL NEWPEN (2)
CALL SYMBOL (4.0, 4.5, HT1, 'SUPERCRITICAL', 0., 13)
CALL SYMBOL (5.6, 4.5, HT1, 'SUBCRITICAL', 0., 11)
CALL SYMBOL (6.9, 4.5, HT1, 'STABLE', 0., 6)
YC = 4.5 - DELY2
CALL SYMBOL (4.0, Y0, HT1, 'INSTABILITY', 0., 12)
CALL SYMBOL (5.6, Y0, HT1, 'INSTABILITY', 0., 11)

RETURN
END

```

```

CHAR 380
CHAR 390
CHAR 400
CHAR 410
CHAR 420
CHAR 430
CHAR 440
CHAR 450
CHAR 460
CHAR 470
CHAR 480
CHAR 490
CHAR 500
CHAR 510
CHAR 520
CHAR 530
CHAR 540
CHAR 550
CHAR 560
CHAR 570
CHAR 580
CHAR 590
CHAR 600
CHAR 610
CHAR 620
CHAR 630
CHAR 640
CHAR 650
CHAR 660
CHAR 670
CHAR 680
CHAR 690
CHAR 700
CHAR 710
CHAR 720
CHAR 730
CHAR 740
CHAR 750
CHAR 760
CHAR 770
CHAR 780
CHAR 790
CHAR 800
CHAR 810
CHAR 820
CHAR 830

```

C
C
C

C
C
C

C
C
C

C
C
C
C

C
C
C
C
C
C

C
C
C


```
.....  
C THE FOLLOWING CARDS CCMPRISE THE DATA DECK FOR PROGRAM STBCONT.  
C  
C /*  
C //GD.SYSIN DD *  
C  
C DATA DECK FROM ONE RUN OF PROGRAM PIPEO (MODENO = 2)  
C  
C /*  
C  
C.....
```


LIST OF REFERENCES

1. Davey, A., and Drazin, P.G., "The Stability of Poiseuille Flow in a Pipe," Journal of Fluid Mechanics, v. 36, part 2, p. 209, 22 August 1968.
2. Garg, V.K., and Rouleau, W.T., "Linear Spatial Stability of Pipe Poiseuille Flow," Journal of Fluid Mechanics, v. 54, part 1, p. 113, 6 January 1969.
3. Gill, A.E., "The Least-Damped Disturbance to Poiseuille Flow in a Circular Pipe," Journal of Fluid Mechanics, v. 61, part 1, p. 765, 3 December 1973.
4. Harrison, W.F., On the Stability of Poiseuille Flow, Ae. E. Thesis, Naval Postgraduate School, Monterey, California, 1975.
5. Huang, L.M. and Chen, T.S., "Stability of Developing Flow Subject to Non-axisymmetric Disturbances," Journal of Fluid Mechanics, v. 63, part 1, p. 183, 16 April 1973.
6. Johnston, R.H. III, A Program for the Stability Analysis of Pipe Poiseuille Flow, M.S. Thesis, Naval Postgraduate School, Monterey, California, 1976.
7. Leite, R.J., An Experimental Investigation of Axially Symmetric Poiseuille Flow, Report No. OSR-TR-56-2, Air Force Contract AF18(600)-350, November, 1956.
8. Naval Postgraduate School Report NPS-67Gn77051, Improved Finite Difference Formulas for Boundary Value Problems, by T.H. Gawain and R.E. Ball, 1 May 1977.
9. Naval Postgraduate School Report NPS67-78-006, A Basic Reformulation of the Pipe Flow Stability Problem and Some Preliminary Numerical Results, by T.H. Gawain, 1 September 1978.
10. Reynolds, O., "An Experimental Investigation of the Circumstances which Determine whether the Motion of Water Shall be Direct or Sinuous, and the Law of Resistance in Parallel Channels," Phil. Trans. Royal Soc., 174, p. 935-982, 1883.
11. Salwen, H., and Grosch, C.E., "The Stability of Poiseuille Flow in a Pipe of Circular Cross-section," Journal of Fluid Mechanics, v. 54, part 1, p. 93, 6 March 1972.

INITIAL DISTRIBUTION LIST

	No. Copies
1. Defense Documentation Center Cameron Station Alexandria, Virginia 22314	2
2. Library, Code 0142 Naval Postgraduate School Monterey, California 93940	2
3. Department Chairman, Code 67 Department of Aeronautics Naval Postgraduate School Monterey, California 93940	1
4. Prof. T.H. Gawain, Code 67Gn Department of Aeronautics Naval Postgraduate School Monterey, California 93940	5
5. LT Michael James Arnold, USN 10825 Single Tree Lane Spring Valley, California 92077	1

Thesis
A7259
c.1

Arnold

Investigation of
pipe flow instability
and results for wave
number zero.

179850

T
A
c

Thesis
A7259
c.1

Arnold

Investigation of
pipe flow instability
and results for wave
number zero.

179850

Investigation of pipe flow instability a



3 2768 002 01257 7
DUDLEY KNOX LIBRARY